Supplementary material for Gourdji et al., "North American CO<sub>2</sub> exchange: inter-comparison of modeled estimates with results from a fine-scale atmospheric inversion"

## Supplement A – Inversion setup and inputs

## A.1 Atmospheric CO<sub>2</sub> data and boundary conditions

### A.1.1 CO<sub>2</sub> mixing ratio data processing and filtering

After preliminary data filtering to exclude low-quality flags and other obvious errors or anomalous spikes in the data, continuous measurements from all towers were averaged to a three-hourly timescale for use in the inversions. Aircraft data were included for measurements below 4000m in altitude, where the strongest influence from surface CO<sub>2</sub> fluxes on the continent is expected (Gerbig et al., 2003a).

With a perfect transport model, it would be possible to use all available well-calibrated CO<sub>2</sub> measurement data in an inversion to help improve the atmospheric constraint on flux estimates. In fact, inversions using synthetic data, and assuming perfect transport, yield better flux estimates with more defined spatial features as measurements are added throughout the day and night (Mueller, 2011). However, including measurements during times of the day that are subject to systematic transport model errors can lead to biased flux estimates, and potentially misleading scientific conclusions (Lin and Gerbig, 2005; Prather et al., 2008).

Following previous inversion studies (e.g. Peters et al., 2007; Schuh et al., 2010; Göckede et al., 2010b), we therefore rely primarily on afternoon CO<sub>2</sub> measurements when vertical convective mixing is strongest, and the height of the PBL is generally well-represented by WRF-STILT (Zhao et al., 2009). Specifically, 3-hourly averages centered at 1 and 4pm were included year-round for all towers, as well as some morning and evening data (3-hourly averages centered at 10am and 7pm) during the height of the growing season when the air should also be well-mixed due to longer day lengths. Also, following data selection criteria for the CarbonTracker system (Peters et al., 2007), some night-time data (3-hourly averages centered on 1 and 4am) were included for the two tall towers (LEF and WKT, Table 1), where towers sample the residual mixed layer representative of  $CO_2$  fluxes from the previous day.

Finally, some of the flask and aircraft data collected on the Pacific coast were excluded, given that a substantial misfit was found at these sites between measured CO<sub>2</sub> mixing ratios and bottom-up fluxes transported forward to the measurement locations. This may be due to difficulties in modeling coastal transport (Riley et al., 2005), or the coarse resolution of WRF-STILT in these locations (i.e. 40km) relative to the scale of meteorological phenomena (Ahmadov et al., 2009).

In addition to systematic transport model errors, using observed continuous CO<sub>2</sub> measurements in a regional inversion is subject to other challenges. First, instruments are subject to failure, creating gaps in the concentration record for some towers. For example, at Harvard Forest, 45% of the potential measurements are missing for the year, particularly in the early part of the year, while at Moody, Texas, observations are missing for January, August, and most of September. Other shorter data gaps occur throughout the year for all towers.

Second, some of the variability in the measurements is due to very local influences (Gerbig et al., 2009) that cannot be resolved by the transport model, the driving meteorological data, or the flux estimation grid, leading to representation errors in the inversion. These data should ideally be excluded, although there is no perfect method for identifying purely local influence. Therefore, with the exception of filtering out extreme events that appear to be plumes from nearby point sources (i.e. short-term increases >30 ppm over background air), this study did not attempt to filter the data for local variability before averaging to the 3-hourly resolution. This decision also leverages the fact that the WRF-STILT model, by simulating transport at relatively high spatiotemporal resolution in a Lagrangian framework, should be better able to resolve local influences relative to coarse-grid Eulerian models (Rastigejev et al., 2010).

Third, measurements are influenced by both land and oceanic fluxes, although oceanic flux locations were not explicitly estimated in this study. This is an issue particularly for coastal locations, where mixing with ocean air can dilute the influence of land fluxes, thereby potentially yielding misleading land flux estimates if this is not taken into account in the inversion framework. Therefore, filtering was applied to exclude measurements with primary sensitivity to ocean fluxes (defined as greater than 90% of the total sensitivity coming from ocean regions, or an integrated land footprint less than 0.5 ppm/( $\mu$ mol m<sup>-2</sup>s<sup>-1</sup>)).

After all data filters were applied, the number of data points included in the inversion per tower, accounting for data gaps, is shown with other tower information in Table A1. In addition, Figure A1 shows a 2-week moving average of 1pm measurements at all nine continuous measurement locations throughout the year, after subtracting the influence of the boundary conditions and fossil fuel emissions.

Table A1: Measurement locations, along with other identifying characteristics of the sites and data included in the inversion. The first 9 locations in the table have continuous data, while the last two represent flask and aircraft measurements from multiple locations across the continent.

Measurement site code	Site name	Site latitude/ longitude	Altitude above ground level (m)	Night-time data included	Number of observations for year
LEF	Park Falls, Wisconsin	45.93N, 90.27W	396	1 and 4 am year-round	1500
WKT	Moody, Texas	31.32N, 97.33W	457	1 and 4 am year-round	959
SBL	Sable Island, Nova Scotia	43.93N, 60.02W	25	None	663
BRW	Barrow, Alaska	71.32N, 156.60W	10	None	248
ARM	Norman, Oklahoma	36.62N, 97.50W	60	None	879
HFM	Harvard Forest, Massachusetts	42.54N, 72.17W	/ 30 None		558
AMT	Argyle, Maine	45.03N, 68.68W	107 None		795
FRD	Fraserdale, Ontario	49.84N, 81.52W	40	None	878
CDL	Candle Lake, Saskatchewan	53.99N, 105.12W	30 None		927
FLA	Flask samples from 6 sites	See Figure 1	0 to 4	None	153
AIR	Aircraft vertical profiles from 15 sites	See Figure 1	139 to 3999	None	943

Figure A1: 2-week moving averages of CO<sub>2</sub> mixing ratio observations (3-hourly averages centered on 1pm local time) for 9 continuous measurement towers in 2004. The combined influence of fossil fuel emissions and the continental boundary conditions on the observations was subtracted, such that the residuals shown here represent only the influence of biospheric fluxes on the observations. Minor gaps were smoothed with linear interpolation, whereas gaps longer than 5 days were removed from the time series. The 2-week moving averages are centered on the dates shown in the x-axis.



### A.1.2 Boundary conditions

The empirical (EMP) North American boundary condition dataset is similar to the boundary condition developed by Gerbig et al. (2003b), and nominally represents a timevarying curtain (latitude, longitude, altitude) of atmospheric CO<sub>2</sub> mixing ratios through the middle of the Pacific Ocean. This curtain was created by extrapolating un-polluted surface and aircraft measurements of CO<sub>2</sub> from the NOAA-ESRL Cooperative Air Sampling Network (Tans & Conway, 2005) in both space and time, in a manner similar to other GLOBALVIEW data products (Masarie & Tans, 1995; GLOBALVIEW-CO2). This boundary condition is strictly valid only for air masses reaching the continent from the West. Relatively little air enters the study area from the East, and given that the dominant outflow from North America is to the East, un-polluted aircraft measurements of CO<sub>2</sub>, which are representative of background air, are lacking over the Atlantic. Across the northern and southern boundaries, available data suggest that longitudinal gradients are likely to be small.

The CarbonTracker (CT) boundary conditions represent gridded 4-dimensional (i.e. latitude, longitude, altitude and time) modeled CO<sub>2</sub> concentrations from the CarbonTracker data assimilation system (Peters et al., 2007, 2010a). Therefore, in contrast to the EMP dataset, the CT boundary conditions contain both longitudinal and synoptic variability. However, these modeled CO<sub>2</sub> concentrations are subject to biases in the inferred global fluxes resulting from the data assimilation system. In fact, the EMP boundary condition dataset was developed in response to known seasonal biases in the CT CO<sub>2</sub> fields, particularly during the Northern hemisphere growing season, for measurement locations upwind of North America (Peters et al., 2010b).

Due to these biases in CT, there is a systematic offset between the two boundary condition datasets, with the influence at the observational sites from the EMP dataset on average about 0.5 ppm lower than the values from CT, with this offset somewhat higher during the growing season. Figure A2 shows a 2-week moving average of the offset by continuous measurement location throughout the year.

Figure A2: Two-week moving averages of the offset between the EMP and CT boundary condition  $CO_2$  time series at the 9 continuous measurement towers in 2004. The 2-week moving averages are centered on the dates shown in the x-axis.



## A.2 – Atmospheric transport model

#### A.2.1 STILT

The Lagrangian particle dispersion model (LPDM) employed in this work is the Stochastic Time Inverted Lagrangian Transport (STILT) model, run in the time-reversed (receptor-oriented) mode. The STILT model is based on the HYSPLIT model (Draxler and Hess, 1998) and has been developed specifically for inverse greenhouse gas flux estimates (Lin et al., 2003; Nehrkorn et al., 2010). Recent examples of its application at continental and regional scales are given by Kort et al. (2008, 2010), Zhao et al. (2009), Gourdji et al. (2010), and Göckede et al. (2010b). As in all LPDMs, atmospheric dispersion in STILT is simulated by tracking a large set of tracer particles, with each particle transported by mean winds obtained from a meteorological model as well as an unresolved, turbulent (subgrid) velocity component computed using the parameterizations of Hanna (1982). The inclusion of both the mean and stochastic wind components (whose interactions are the basic cause of dispersion in the atmosphere) sets the LPDM approach apart from conventional trajectory models that employ mean winds only, and thus cannot properly simulate dispersion or surface interactions (Stohl, 1998; Stohl et al., 2003; Uliasz and Pielke, 1990; Uliasz, 1994). The LPDM approach prevents particle tracks from intersecting the surface for numerical reasons (e.g., a strong descent forcing a particle to a level below the surface) as is common in trajectory models. Instead, the turbulent winds invoked by the LPDM ensure that particles travel through the planetary boundary layer (PBL) in a physically reasonable way.

In this application, the STILT model transports ensembles of 500 particles backwards in time 10 days from each of a set of receptor points (in this case, from a set of towers at 3-hourly resolution) using the input meteorological data. For each receptor, we calculate the response of the target gas concentration at the receptor point to surface sources ("footprint"), in units of ppmv/(µmol m<sup>-2</sup> s<sup>-1</sup>). The footprint, which represents the adjoint of the transport field, is calculated by counting the number of particles in a surface-influenced region (defined as a fraction of the estimated PBL height) for a given time period spent in the region (for details, see Lin et al., (2003)). When multiplied by an *a priori* field of surface flux, the footprint gives the associated contribution to the mixing ratio measured at the receptor.

The STILT model has undergone a number of upgrades, most recently a merger with parts of the latest HYSPLIT code, resulting in improved handling of nested fields. The model is being developed by a worldwide consortium, with the most up-to-date software and information available at <u>www.stilt-model.org</u>.

#### A.2.2 Meteorological Input

For the inversions described in this paper, the STILT model has been driven by meteorological fields from a version of the WRF model (Skamarock and Klemp, 2008) customized for STILT and other transport models (Nehrkorn et al., 2010). In particular, these WRF outputs include convective mass fluxes that are used directly in the STILT dispersion calculations, and time-averaged mass fluxes (rather than instantaneous advective velocities) are used to drive STILT, which results in good mass conservation (a critical consideration for inverse flux estimates). We have employed version 2.2 of WRF (<u>www.wrf-model.org</u>) over a domain covering North America, with two nested 10- and 40-km grids, as shown in Figure 1 in the full manuscript. NARR fields have been used for initial and lateral boundary conditions and for analysis nudging of WRF. To prevent drift of the WRF forecasts from the analyses, forecasts were reinitialized every 24 hours (at 00 UTC). Forecasts were run out to 30 hours, but only hours 7-30 from each forecast were used to avoid spin-up effects during the first 6 hours from each forecast. Model fields were output hourly and archived in the native WRF netcdf format. For use in the STILT model, a subset of model fields was archived in the (compressed) ARL format, at a substantial savings in storage.

A summary of numerical and physics options used in these runs is given in Table A2.

Option	Description
Land-surface	Noah land-surface model with Monin-Obukhov surface layer (Ek et al., 2003)
PBL package	Yonsei University (YSU) scheme (Hong et al., 2006)
LW radiation	RRTM (Mlawer et al., 1997)
SW radiation	Goddard (Chou and Suarez, 1994)
Microphysics	Lin et al. (1983); Chen and Sun (2002)
Convection	Grell & Devenyi (2002)
Nesting	One-way
Nudging	u,v,T,q at all levels above PBL, every 3 hours, 1 hour relaxation time
Time stepping	3 <sup>rd</sup> order Runge-Kutta; 4 short time steps per long time step
Advection	5 <sup>th</sup> order horizontal, 3 <sup>rd</sup> order vertical
	positive definite advection for moisture and scalars
Diffusion	2 <sup>nd</sup> order horizontal diffusion using Smagorinsky first-order closure
Damping	No upper level or vertical velocity damping; default values for divergence and external model damping

Table A2: Numeric and physics options for WRF runs.

## A.3 – Branch and Bound variable selection method

The Bayes Information Criterion (BIC) is a criterion-based variable selection method, allowing for the comparison of all possible combinations of a superset of auxiliary variables (Ward, 2008). The original BIC equations from Schwarz (1978) were modified for a geostatistical setup with correlated residuals by Mueller et al., (2010), and were further updated here to be compatible with an inverse modeling formulation using the atmospheric measurements. After these modifications, the criterion that must be minimized for the BIC approach can be expressed as:

$$BIC = \ln|\Psi| + [\mathbf{z}^{\mathrm{T}}(\Psi^{-1} - \Psi^{-1}\mathbf{H}\mathbf{X}(\mathbf{X}^{\mathrm{T}}\mathbf{H}^{\mathrm{T}}\Psi^{-1}\mathbf{H}\mathbf{X})^{-1}\mathbf{X}^{\mathrm{T}}\mathbf{H}^{\mathrm{T}}\Psi^{-1})\mathbf{z}] + p\ln(n)$$
(A1)

where  $\Psi = HQH^T + R$ , *n* is the number of observations, *p* represents the number of covariates included within a given model (*X*), and all other variables are as described in Appendix A. Given that comparing all possible models (2<sup>p</sup>) quickly becomes computationally expensive and perhaps infeasible for large supersets of variables, the BIC was implemented with a Branch-and-Bound algorithm (Land and Doig, 1960) to help make the problem computationally tractable. This algorithm avoids unnecessary matrix multiplications by eliminating model "branches" as it runs that cannot possibly contain the "best" model (Yadav et al., in review).

# Supplement B – Supplementary tables and figures

Table B1: Correlation coefficients among $\beta$	uncertainties for variables included in the NARR					
inversion with the EMP boundary conditions.						

	Evapo-	Precipitation	Specific	Air Temperature (@
	transpiration	rate	humidity	2m)
Evapotranspiration	1.00			
Precipitation rate	0.18	1.00		
Specific humidity	-0.14	-0.28	1.00	
Air temperature (@2m)	-0.13	-0.07	-0.76	1.00

Figure B1: Contribution of auxiliary variables and stochastic component to the best estimates of flux in the GIM/ NARR inversion with EMP boundary conditions in April 2004 (see equation 5 in Gourdji et al., 2010).

