The three solicited reviews were extremely helpful and constructive, and I believe that all of their editorial suggestions and requests for clarification have been taken. The unsolicited comments from Shakhova, Tumskoy, and Romanovskii seemed to show a strong preference for the hypothesis that a climatically-significant pulse of methane could erupt from the Siberian seafloor to the atmosphere in a time span of just a few years. I am unable to think of a way to get a model to do this, but addressing their concerns has also made the paper clearer.

Specific issues are discussed first in this document, updating text from the Author's comment posted to refer to the revised manuscript. A lineby-line discussion of changes in the paper follows that.

Reviewer Nicolsky requested the source code to be included in the supplemental material, which I have done.

## Methane hydrate stability in the permafrost zone

One of the strongest and least speculative conclusions of this paper is that methane hydrate is thermodynamically unstable in most of the permafrost zone in the upper sediment column, and will therefore not accumulate or be found there. Shakhova cited soil at temperatures of –17 °C, which is below the freezing temperature for hydrate if salinity is 35 PSU (as assumed by [Romanovskii et al., 2005]), claiming that these sediments are clearly within the hydrate stability zone. I believe this claim is incorrect.

The thermodynamics are illustrated in a phase diagram for ice and hydrate, using salt as a master variable. This figure will be the new Figure 1 in the revised paper. When the system consists only of ice and fluid phases, the equilibrium salinity  $S_{\rm eq}$  increases with decreasing temperature below freezing (Figure 1a, left). Above the melting temperature, ice is unstable, as indicated by the nonzero values of the disequilibrium temperature,  $\Delta T_{\rm eq,\,ice} = T - T_{\rm eq,\,ice}$ , in contours, even in zero-salinity water (right). For a system consisting of only the hydrate and fluid phases (assuming that ice formation is disallowed, and also assuming gas saturation for methane) (Figure 1b), the behavior is similar but with an added pressure dependence due to the compressibility of the gas phase. When both solid phases are allowed, the overall equilibrium salinity will whichever is higher between  $S_{\rm eq,\,ice}$  and  $S_{\rm eq\,hydrate}$ . Whichever phase can seize water at its lowest activity (from fluid of the highest salinity) will be

the stable phase. The salinity of the brine excluded from that phase will be too high to permit the existence of the other solid phase at that temperature. The contours show  $\Delta T_{eq}$  for hydrate (solid) and ice (dashed).

Figure 1d shows  $\Delta T_{eq, \, hydrate}$  in colors and contours of the excess salinity relative to hydrate equilibrium,  $S_{max}$  -  $S_{eq, \, hydrate}$ . Hydrate is only stable when  $\Delta T_{eq, \, hydrate}$  is zero (purple color). Under permafrost conditions of low pressure and low temperature (upper left corner),  $\Delta T_{eq, \, hydrate}$  is greater than zero, indicating that hydrate is unstable, coinciding with the salinity forcing from the ice, in overlain contours. A similar exclusion of ice in part of the hydrate stability zone is seen Figure 1e, but this would only happen in nature in conditions of unlimited methane. The resulting phase diagram for ice and methane hydrate is shown in Figure 1f. Hydrate stability is suppressed in the permafrost zone by this thermodynamic mechanism.

I have done a sensitivity study in which pure- $H_2O$  ice is forbidden to form, which has the effect of allowing the hydrate stability zone to outcrop near the sediment surface, briefly, during the coldest times in the glacial cycle (Figure 12). This altered-physics simulation, even though indefensible, does not generate a large transient methane emission spike in response to a global warming forcing. This is because the sediment column has already been subjected to the extreme temperature change of inundation by rising sea level. The stability zone boundary quickly retreated downward when that happened. A temperature anomaly from global warming has to diffuse hundreds of meters into the sediment column, to catch up to and accelerate that retreating hydrate stability boundary.

## **Sensitivity studies**

Many of the reviewers' comments can be addressed by doing model sensitivity studies, which I have now done. A new section has been added to the text, "3.5 Summary of Model Sensitivity Studies". Figures 15 and 17-19 have been expanded to show the new results, and a table has been added. A web page server for movies from the model has been updated to include 68 new sensitivity study movies, at

http://geosci.uchicago.edu/~archer/spongebob. Reviewer Nicolsky called for these most directly, but questions about uncertainties in initial salinity, geothermal temperature gradient, rates of hydrological flow, and permafrost impact on gas mobility, can also be addressed by showing the

model sensitivity to the parameter in question. Simulations of the glacial cycles and global warming forcings have been done for the parameters summarized below:

Biogenic and thermogenic methane production rates. The model brings no real constraint to the rates of methane production, since the sediment accumulation history etc. are not well-constrained. See the discussion on the initial condition to the glacial simulations, next. New sensitivity runs were done scaling the biogenic and thermogenic methane production rates by arbitrary values, to see the impact on the methane cycle in the model. Increasing production increased the release rate to the atmosphere, but in all cases the release was modulated in time through the glacial cycles by inhibition of gas flow by permafrost during low sea-level, and by inundation in the ocean during high sea level. Deposition of high-POC soils on land (Yedoma) increased the methane flux to the atmosphere but had only a small impact on the abundance of methane hydrate.

**Permafrost inhibition of gas migration**. It turns out that changing the critical ice fraction at which to shut down methane gas mobility has a large impact on the glacial cycles of methane fluxes to the atmosphere, indicating a need for more attention to be paid to this question.

Geothermal temperature gradient. Changes in the temperature gradient with depth in the sediment column alter the depth of the base of the hydrate stability zone. The top of the zone remains unchanged, because surface sediment temperatures are less affected by the change in heat flux, anchored as they to ocean temperatures.

Flow. Runs were done including the vertical permeable channels (at the request of review #1), and excluding horizontal flow. Neither of these changes affects the model behavior or main conclusions in a fundamental or significant way.

# Glacial cycle simulation initial condition

Tumskoy was offended at the over-simplicity of the spinup simulation used to generate the initial condition for the glacial simulations. The more usual approach in modeling hydrates is to start with an ad-hoc initial condition (for example, [Reagan, 2008; Reagan and Moridis, 2009; Reagan et al., 2011]). For SpongeBOB the model state at any time is the result of the time history of sedimentation, which is driven by the time-evolving depth of the sea floor, and interacting with isostatic adjustment of the crust. The simplest way to generate an initial condition in the model without a startup transient is to spin the model up at low

resolution. Because of the over-simplicity of the tectonic, sea level, and sedimentation forcing of the spinup phase, its POC concentrations and methane production rates do not constrain those of the real shelf. The sensitivity of the glacial methane cycles to these uncertain methane production rates is evaluated by scaling the model methanogenesis rates from the spinup result. I would clarify this in the text, and I hope that the sensitivity studies will make this point more concretely as well.

# Prior literature on hydrology and the fresh / salt water boundary

Reviewer #1 wondered if the ratcheting effect of the fresh-water pump described in my results had been modeled before. I found some hydrological models of the salt / fresh boundary as it changes with sea level change [Kooi et al., 2000; Lu and Werner, 2013; Watson et al., 2010]. These all share the result that freshening is much faster than getting saltier, although these were applied to systems of much smaller space (~10 meters depth) and time (~century) scales than I addressed. These models included mechanisms for salt fingering, which reviewer Nicolsky pointed out is lacking in my code. Since salt fingering is a diffusive mechanism while hydrological flow is advective, salt fingering should get even less important as the size and time scale for the domain increases.

# Ocean temperature boundary condition

Tumskoy questioned the validity of the temperature changes in the ocean through the glacial cycles. The model assumed uniform water column temperatures ranging from +1 to -2 °C. For the region of interest, the continental shelf in shallow water, this forcing was incorrect, given that overlying waters are still near freezing today, during an interglacial. All of the new model glacial cycle runs are done with a more realistic thermal boundary condition, with temperature changes at depth in the water column (below about 200 meters) as before, but waters in the top 100 meters pegged at -1.5 °C until the global warming scenario begins.

# Glacial / interglacial atmospheric methane fluxes.

Shakhova claims that the higher methane concentration in the interglacial atmosphere is proof that emission from the Arctic could not have been greater during glacial time. However, the Arctic is, then as now, a small

part of the global atmospheric methane budget, so one cannot constrain Arctic emissions from the global concentrations recorded in ice cores.

# Sub-grid scale flows.

Tumskoy called the response time of the model to a sea level change "absurd", but I don't know of modeling results or measurements of the time scale for adjustment of a salty sediment column on initial exposure to fresh water forcing, on which to base that claim. Permeabilities of the deep sediment column are poorly constrained, and not very important to the major conclusions of the paper. He also claims that I ignore the literature on hydrological flow in permafrost. Flow in my model is attenuated by ice, although it was not stated as clearly as it could have been, so perhaps it was missed. It is not completely impeded, so that flow through taliks and faults can still in principal be accounted for as subgrid scale process, the way clouds are done in climate models. Flow in the model is not well constrained or accurate, but it turns out not to matter to the main conclusions of the paper, according to a new sensitivity study in which horizontal flow is disabled. Although offshore flow carrying methane appears to be an important flux in the methane budget, this sensitivity run only had subtle impacts on the other components of the budget (Figure 17).

# Application to a large-scale abrupt methane release due to global warming.

The proposal from Shakhova is that methane has been building up as bubbles, sealed by permafrost, which is now unsealing, or as methane hydrate in shallow sediments, which is now melting, to deliver, on a time scale of a few years, about 50 Gton of methane to the atmosphere. For a transient gas in the atmosphere like methane, it makes a huge difference if it is released in a few years, or a few hundred years. I don't doubt the possibility of mobilizing 50 Gton of methane eventually, only the fast time constant (which is required to get a strong methane-driven climate impact). An expanded discussion of this topic begins on 1043.

# Line-by-line description of manuscript changes

(line numbers from the file with changes highlighted).

Abstract has been completely rewritten to reflect the new sensitivity studies, and to clarify the thermodynamic underpinning of the conclusion.

Section 1.1 clarified as to initial condition and sensitivity studies.

Beginning line 312, discussion deleted of model architecture on the water table depth, which is not really applicable to the current simulations.

Line 365. Added comment about canyon spacings chosen, based on request from reviewer #1.

Section 2.3 Permafrost. Rewritten to explain more clearly the goal of the model (thermodynamic equilibrium), rather than the actual formulation, which was difficult to evaluate (reviewer Nicolsky).

Line 421. Discussion added of thermodynamic competition between ice and hydrate.

Line 503. Corrected my embarrassing confusion of horsts vs. grabens, and clarified the connection to the model results. A deep sediment column is required in order to get thermogenic methane production, a central part of the simulation. However, since TG methanogenesis is now varied in sensitivity studies, it is less essential that the detailed history, and therefore TG methanogenesis rate, of the model sediment column match those of the real deep sediment column.

Initial Spinup section. This was newly added to explain why the initial condition was spun up as it was (rather than chosing an ad-hoc initial condition), and how the sensitivity studies are used to account for the over-simplicity and inaccuracy of the spinup phase of the simulation.

Line 608. New discussion was added on the time asymmetry between freshwater invasion into a sediment column driven by pressure head vs. reinvasion of salt after re-immersion. The model lacks salt fingering, as pointed out by Nicolsky, but the time asymmetry is also found in smaller-scale models that include this effect (cited in the text). As the spatial scale gets larger, advective processes should become even more dominant over diffusive ones. This freshening of the sediment column has not been documented in detail in the Arctic (as pointed out by Romanovskii), although lower-than-marine salinity pore waters have

documented there. The freshened vs. marine simulations are used as a sensitivity study to a reasonable range of initial salinities, so an exact knowledge of the initial salinity is not claimed or required (because the sensitivity to this parameter was actually rather subtle).

Line 749. New discussion on the competition between ice and hydrate, in the form of an altered-physics simulation in which ice is not allowed to form. The depth range of the hydrate stability zone changes in this run, but there is still no massive  $CH_4$  degassing upon global warming, as looked for by Shakhova and her colleagues.

Line 842. Discussion rewritten on Figures 15, 17, and 18, to reflect the new sensitivity runs. Figure 15 also now shows the inventory of methane as bubbles through the simulations, in response to the hypothesis of Shakhova.

Section 3.5 This is entirely new, a summary of sensitivity studies.

Line 1091. Rewritten discussion as requested by Nicolsky on the Shakhova hypothesis.

1 2	A model of the methane cycle, permafrost, and hydrology of the Siberian continental margin
3	David Archer, University of Chicago
4	d-archer@uchicago.edu
5	
6	Abstract
7 8 9 10 11 12	A two-dimensional model of a passive continental margin was adapted to the simulation of the methane cycle on Siberian continental shelf and slope, attempting to account for the impacts of glacial / interglacial cycles in sea level, alternately exposing the continental shelf to freezing conditions with deep permafrost formation during glacial times, and immersion in the ocean in interglacial times. The model is then subjected to a potential future climate warming scenario.
14 15 16 17 18 19 20 21 22 23 24	Pore fluid salinity plays a central role in the model geochemical dynamics. In the permafrost zone, pure water ice tolerates a higher fluid salinity than methane hydrate can, eliminating hydrate as an equilibrium phase. An analogous region in the ice – hydrate – brine phase diagram excludes ice in favor of hydrate, but the two phases can coexist at a sub-saturated methane concentration. In the permafrost zone (cold and low pressure), in contrast, the dissolved methane concentration cannot be higher than equilibrium with gas, so the hydrate exclusion from this zone is inescapable. This thermodynamic constraint restricts methane hydrate to at least 300 meters depth below the sediment surface, precluding a fast hydrate dissolution response to sea-floor warming.
25 26 27 28 29 30 31 32	The initial salinity of the sediment column may have been affected by previous hydrological forcing, because freshwater invasion driven by a pressure head is probably much faster than salinity invasion due to convective-diffusive processes. This has a ratcheting effect, leaving relict fresh water lenses below sea level in many parts of the world. The pore fluid salinity determines the relative volumes of the ice, brine, and hydrate phases in the sediment column, and therefore the timing of ice formation and melting, but the chemical composition, in particular the
33	salinity of the brine phase, is fixed, in equilibrium, by the local

35 sensitive to the initial salinity of the sediment column. 36 Through the glacial / interglacial cycles, Permafrost formation inhibits bubble transport through the sediment column, by construction in the 37 38 model. The impact of permafrost on the methane budget is to replace 39 the bubble flux by offshore groundwater flow containing dissolved methane, rather than accumulating methane for catastrophic release 40 41 when the permafrost seal fails during warming. By far the atmospheric 42 methane flux is affected most strongly by largest impact of the glacial / interplacial cycles on the atmospheric methane flux is attenuation 43 44 bychanges in sea level, because dissolution of bubbles dissolve in the 45 ocean when sea level is high. Methane emissions to the atmosphere are 46 highest during the regression-sea-level fall part of the cycle (as soil is freezinging) part of the cycle, rather than during during transgression 47 48 (the warming deglaciations thawing). Timings of the atmospheric methane 49 flux changes are sensitive to assumptions made about bubble transport inhibition by permafrost. The atmospheric flux is sensitive to biogenic 50 and thermogenic methane production rates, but the hydrate inventory is 51 52 only sensitive to thermogenic methane production. The geothermal heat 53 flux affects the thickness of the hydrate stability zone (primarily the 54 depth of its base), but not the inventory of hydrate in the model until a 55 low-gradient threshold is passed. The model produces methane inventory 56 changes of 50 Gton C as bubbles, and as much as hundreds of Gton C as 57 hydrate, but these reservoir changes interact mostly with pore water 58 dissolved methane rather than driving immediate methane loss from the 59 sediment column. 60 The model-predicted methane flux to the atmosphere in response to a 61 warming climate is small, relative to the global methane production rate, 62 because of the ongoing flooding of the continental shelf. The 63 atmospheric methane flux response to sudden warming takes thousands 64 of years, because of the slow thermal diffusion time to the hydrate 65 stability zone, and because a warming perturbation beginning now would 66 follow a much larger warming perturbation that started thousands of

temperature. The model hydrate inventory on the shelf is however

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thousands of years in the future, the increased methane flux increase due

to warming could be completely counteracted by sea level rise, which decreases the efficiency of bubble transit through the water column.

years ago, when the sediment surface flooded. On time scales of

71 The model is used to gauge the impact of the glacial cycles, and potential anthropogenic warming in the deep future, on the atmospheric methane 72 emission flux, and the sensitivities of that flux to processes such as 73 permafrost formation and terrestrial organic carbon (Yedoma) 74 75 deposition. A slight increase due to warming could be completely counteracted by sea level rise on geologic time scales, decreasing the 76 77 efficiency of bubble transit through the water column. The methane cycle on the shelf responds to climate change on a long time constant of 78 79 thousands of years, because 80 Hydrological forcing drives a freshening and ventilation of pore waters in areas exposed to the atmosphere, which is not quickly reversed by 81 82 invasion of seawater upon submergence, since there is no analogous saltwater pump. This hydrological pump changes the salinity enough to 83 84 affect the stability of permafrost and methane hydrates on the shelf. 85 hydrate is excluded thermodynamically from the permafrost zone by 86 water limitation, leaving the hydrate stability zone at least 300 meters 87 below the sediment surface. 88

#### 1. Introduction

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### 1.1 The Siberian Continental Shelf System

The Siberian Arctic continental shelf has been the focus of attention from scientists and the public at large for its potential to release methane, a greenhouse gas, in response to climate warming, a potential amplifying positive feedback to climate change [Shakhova, 2010; Westbrook, 2009][Shakhova, 2010; Westbrook, 2009]. The goal of this paper is to simulate the geophysical and carbon cycle dynamics of the Siberian continental margin within the context of a basin- and geologic time-scale mechanistic model of the coastal margin carbon cycle called SpongeBOB [Archer et al., 2012][Archer et al., 2012]. An initial condition for the glacial cycle simulations was generated by configuration of sediment structure and composition was spunspinning the up at low resolution over 62 million simulated years. Then the model at higher resolution is driven by cyclic changes in sea level and air temperature resulting from glacial cycles, to simulate the impact of the hydrological pressure head and permafrost formation on the fluid flow and methane cycle on the shelf. Finally, an 100,000-year interglacial interval in the simulation is subjected

107 to anthropogenic warming of the overlying water and potential 60-meter 108 changes sea level. Sensitivity studies are presented for the biogenic and 109 thermogenic methane production rates, initial salinity, geothermal 110 temperature gradient, rates of hydrological flow, and permafrost impact 111 on gas mobility. 112 1.1.1 Permafrost 113 One component of the problem simulation is a wedge of frozen sediment 114 (permafrost) submerged beneath the ocean on the continental shelf of 115 Siberia, left behind from glacial time when the shelves were exposed to 116 the frigid atmosphere by lowered sea level [Romanovskii and Hubberten, 117 2001] Romanovskii and Hubberten, 2001]. The ice is thought to provide 118 a seal to upward migration of methane gas [Shakhova et al., 119 2009] Shakhova et al., 2009, especially where ancient fresh 120 groundwater flow produced a layer of very high saturation ice infill, a 121 formation called the Ice Complex in Siberia [Romanovskii et al., 122 2000] Romanovskii et al., 2000], although there are high ice saturations found in the Alaskan Arctic as well [Zimov et al., 2006] [Zimov et al., 123 124 <del>2006]</del>. 125 With inundation by the natural sea level rise over the last 10+ thousand 126 years, the permafrost is transiently melting, although the time constant 127 for this is generally long enough that significant frozen volume remains, 128 especially in shallower waters which were flooded more recently [Khvorostyanov et al., 2008a; Nicolsky and Shakhova, 2010; Romanovskii 129 130 and Hubberten, 2001; Romanovskii et al., 2004; Shakhova et al., 2009; Taylor et al., 1996][Khvorostyanov et al., 2008a; Nicolsky and Shakhova, 131 2010; Romanovskii and Hubberten, 2001; Romanovskii et al., 2004; 132 Shakhova et al., 2009; Taylor et al., 1996]. Even overlying water at the 133 134 freezing temperature can provoke subsurface melting by providing a 135 warmer boundary condition against which geothermal heat establishes the 136 subsurface temperature profile, but with climate warming, the waters 137 could surpass the freezing temperature, allowing heat to flow from above 138 as well as below [Khvorostyanov et al., 2008b] [Khvorostyanov et al., 139 2008b] Elevated methane concentrations have been measured in the water 140 141 column over the Siberian shelf, even in areas of shallow water where the permafrost should still be strongly intact [Shakhova, 2010; Shakhova et 142 al., 2005][Shakhova, 2010; Shakhova et al., 2005]. Chemical and 143

144 isotopic signatures of hydrocarbons adsorbed onto surface sediments 145 indicate a thermal origin [Cramer and Franke, 2005] Cramer and Franke, 146 2005], suggesting that the methane is produced many kilometers deep in 147 the sediment column. The apparent ability for this methane to transverse 148 the barrier of the Ice Complex has been attributed to hypothesized 149 openings in the ice (called "taliks"), resulting from lakes or rivers on the 150 exposed shelf, or geologic faults [Nicolsky and Shakhova, 2010; 151 Romanovskii et al., 2004; Shakhova et al., 2009][Nicolsky and Shakhova, 2010; Romanovskii et al., 2004; Shakhova et al., 2009] 152 153 1.1.2 Salt 154 Dissolved salt in the pore waters can have a strong impact on the timing 155 of thawing permafrost [Nicolsky and Shakhova, 2010; Shakhova et al., 156 2009][Nicolsky and Shakhova, 2010; Shakhova et al., 2009]. When sea 157 level drops and exposes the top of the sediment column to the 158 atmosphere and fresh water, t\(\pi\)he salinity of the subsurface pore waters 159 of a sediment column exposed to the atmosphere can be flushed out by hydrological groundwater flow, driven by the pressure head from the 160 161 elevated terrestrial water table above sea level. Ground waters tend to 162 be fresh, a product of hydrological cycle, rather than saline as in marine 163 sediments. The boundary between fresh and salty pore water tends to 164 intersect the beach sediment surface at the water's edge [Moore et al., 2011][Moore et al., 2011]. From there, the boundary tends to dip 165 166 landward, to a depth of approximately 40 meters below sea level for 167 every 1 meter of elevation of the table water. The ratio of water table 168 elevation to freshwater lens depth is driven by the relative densities of 169 fresh and salt water, as the fluid seeks an isostatic balance in which the 170 fresh water displaces an equal mass of salt water [Verriuit, 1968][Verrjuit, 1968] 171 172 The SpongeBOB model has been modified to simulate the processes 173 responsible for these observations. We do not attempt to simulate a 174 detailed outcropping history over 62 million-year spinup time of the 175 sediment column, but rather demonstrate the general process by 176 subjecting the nearly complete sediment column to a one-time sea level 177 lowering, exposing the continental shelf to groundwater forcing. After a 178 few million years, the sediment column subsides, due to compaction and 179 absence of sediment deposition, resulting in a sediment column that has 180 been considerably freshened by the atmospheric exposure. We find that 181 **t**This freshening persists in the model for millions of years, because there

182 183 184 185 186	is no corresponding "salt-water pump" during high sea-level stands. This behavior is consistent with the discovery of vast nearly fresh aquifers in currently submerged continental shelf regions around the world [Post et al., 2013][Post et al., 2013], left over from groundwater forcing during glacial time.
187	1.1.3 Carbon
188 189 190 191 192 193 194 195 196 197	The Another component of the Arctic methane storysimulation is the Yedoma, deposits of wind-blown dust and organic carbon that accumulated on the coastal plains of exposed continental shelves during glacial times [Zimov et al., 2006][Zimov et al., 2006]. The deposits contain a substantial fraction of organic carbon, consisting of grass roots and remains, preserved by the freezing conditions. When they thaw, they begin to release CO <sub>2</sub> and methane to the atmosphere [Dutta et al., 2006; Schuur et al., 2008; Zimov et al., 2006][Dutta et al., 2006; Schuur et al., 2008; Zimov et al., 2006]. Oxidation of the carbon can give off enough heat to accelerate the melting driven by primary climate forcing [Khvorostyanov et al., 2008b].
199	2. Model Description
200	2.1 Previously Published Model Formulation
201 202 203 204 205 206 207 208 209 210 211 212 213 214	SpongeBOB is a two-dimensional basin spatial-scale and geological time-scale model for the methane cycle in continental margin sediments. The model, configured for a passive margin basin, was described by Archer et al [2012][Archer et al., 2012], as applied to the Atlantic coast of the United States. The model attempts to "grow" a sediment column based on first principles or parameterizations of sediment and pore water physical and chemical dynamics. The approach integrates our understandingprocesses of the carbon and methane cycles within the evolving sediment column matrix, providing constraints to the rates and processes that may inform the response of the system to future changes in climate. Where model parameterizations or parameters are poorly constrained, The prognostic model can also be applied in sensitivity studies are used to assess which of the uncertainties are the most significant.
215 216 217	the impact of large-scale geophysical drivers such as ocean temperature and carbon cycle dynamics that might have shaped the evolution of the sedimentary methane cycle in the geologic past.

218 Sediment is delivered from the coast of the model as riverine material. 219 and it settles according to a parameterization of grain size, with finer 220 material advecting further offshore before deposition. The organic 221 carbon concentration of the depositing material is determined in the 222 model as a function of water depth at the time of sedimentation. Rather 223 than attempt to simulate the complex biogeochemical dynamics of the 224 ocean and surficial sediments (early diagenesis), the POC fraction, and 225 the H/C ratio of the organic matter, are specified by a parameterization 226 based on water\_-depth to reproduce the observed patterns of sediment 227 surface POC deposition, as a driver to the subsurface model. 228 The H/C ratio of the depositing organic matter limits the potential extent 229 of methane production from the organic matter. The degradation rate of 230 organic carbon is estimated based on its age, a relationship that captures 231 many orders of magnitude of variability in the natural world [Middelburg 232 et al., 1997 [Middelburg et al., 1997]. The reaction pathways presume a 233 reactive intermediate H<sub>2</sub>, which either reduces SO<sub>4</sub><sup>2-</sup> if it is available or it reacts with DIC to produce methane. Isotopic fractionation of CO<sub>2</sub>, CH<sub>4</sub>, 234 235 and radioiodine are simulated by maintaining parallel concentration fields 236 of different isotopologs, and applying fractionation factors to the 237 chemical kinetic rate constants or equilibrium conditions. Dissolved 238 methane in the pore water has the potential to freeze into methane 239 hydrate or degas into bubbles, depending on the T, Ptemperature. 240 pressure, salinity, and CH<sub>4</sub> concentrations. 241 Sediment compaction drives pore fluid advection through the sediment 242 column, but the fluid flow is also focused in some simulations by ad hoc 243 vertical channels of enhanced permeability, to simulate in at least a 244 qualitative way the impact of heterogeneity in the fluid flow on the 245 characteristics of the tracer field. Methane hydrate is concentrated in 246 these channels by focused upward flow, and the pore-water tracers in the 247 channels resembles that of hydrate-bearing regions (in SO<sub>4</sub><sup>2</sup>-248 concentration and 129-lodine ages). 249 Most of the model configuration and formulation was described by Archer

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et al. [2012][2012]. Here we describe in detail tThe new modifications

required to simulate groundwater hydrological flow and permafrost

formation are described in detail below.

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## 2.2 Groundwater Hydrology

254 **2.2.1 Pressure Head** 

- 255 When the sediment column is exposed to the atmosphere, the pressure
- 256 field from the variable elevation of the water table (the pressure head)
- 257 begins to affect the fluid flow. The pressure head for a fluid particle at
- 258 the depth of the water table varies as
- 259  $P_{\text{head}}(z) = g \int_{z}^{z_{\text{wt}}} \rho_{\text{seawater}} dz$

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- 260 where  $z_{wt}$  is the elevation of the water table. The pressure head at each
- depth in the domain is a function of the physical water table height above
- it and the density anomalies integrated from the water table to the depth
- of the point in question. The pressure head resulting from a varying
- 264 water table can therefore be altered at depth by variations in pore fluid
- 265 density driven by salinity or temperature.
- **2.2.2 Fluid Flow**
- 267 The pressure head acts in concert with the excess pressure P<sub>excess</sub> to drive
- 268 horizontal Darcy flow through the sediment, as

$$269 \qquad u_{\text{Darcy}, i \to i+1} = \frac{k_{\text{h}, i} + k_{\text{h}, i+1}}{2\mu} \frac{\left(P_{\text{excess}, i} - P_{\text{excess}, i+1}\right) + \left(P_{\text{head}, i} - P_{\text{head}, i+1}\right)}{(\Delta x_i + \Delta x_{i+1})/2}$$

270 while the vertical flow in the model is driven only by compaction pressure

$$271 \qquad w_{\text{Darcy}, j \rightarrow j+1} = \frac{k_{v,j}}{\mu} \frac{P_{\text{excess}, j} - P_{\text{excess}, j+1}}{(\Delta z_i + \Delta z_{i+1})/2}$$

- 272 The value of P<sub>excess</sub> is determined from the porosity and sediment load of
- the sediment in each grid box, as described in Archer et al [2012][Archer
- 274 et al., 2012]. -An assumed sediment rheology is used to calculate the
- load-bearing capacity of the solid matrix within a given grid cell. Pexcess is
- 276 calculated by assuming that the load of the solid phase overlying the grid
- 277 cell that is notis carried by the solid matrix must be carried by the sum of
- 278 the load-bearing capacity of the sediment (a one-to-one function of
- 279 porosity), and P<sub>excess</sub> carried by in the fluid phase, which is calculated by
- 280 difference. When ice forms (described below), it leaves Percess unchanged.

281 282	but the flow is inhibited by scaling the permeability k by the decrease in fluid porosity.
283 284 285 286 287 288 289 290 291 292 293 294 295	In previous versions of the SpongeBOB model, the fluid flow was calculated explicitly, each time step, as a function of $P_{\rm excess}$ at the beginning of the time step. Numerical stability motivated a modification of the vertical flow to an implicit numerical scheme, which finds by iteration an internally consistent array of vertical flow velocities and resulting $P_{\rm excess}$ values from a time point at the end of the time step. Ocean and atmosphere models often use this methodology for vertical flow. A benefit to this change is stability in the vertical flow field, as a reductingon in the numerical noise in the flow, which that can causees trouble when it interacts dynamically with other aspects of the model such as ice formation. Implicit schemes can be more efficient computationally, but in this case the execution time is not improved by the implicit method, just the stability.
296 297 298 299 300 301	Note that the flow scheme in its formulation is entirely elastic, whereas in reality, pore fluid excluded by the pressure of a sediment column above sea level, for example, where it is uncompensated by buoyancy in seawater, should remain excluded when sea level rises again, like toothpaste from the tube. However, my attempts to embed this plastic behavior into an implicit solver failed to converge.
302	2.2.3 Water Table Depth
303 304 305 306 307 308 309	The model maintains $z_{wt}$ , the elevation of the water table within the sediment column, as a continuous variable that ranges through the discreet vertical grid of the model. The formulation allows boxes to be empty of water or partially "saturated" at the top of the fluid column. In these simulations, however, the water table remained very close to the sediment surface, as unsaturated soil produced by subsurface flow is quickly replenished by hydrological recharge.
310 311	It was necessary to treat $z_{\rm wt}$ as a continuous variable because of the impact of the pressure head on the fluid flow.
312 313 314 315 316	The evolution of the height of $z_{\rm wt}$ is determined by changes in the relative volumes of fluid and air phases. In grid cells at and above the depth of the water table, the air volume is computed such that the porosity of the cell relaxes toward the drained condition, in which $P_{\rm excess} = 0$ . In this way the outcropping grid cell provides an upper pressure boundary condition for

317 the model, as opposed to grid columns which are entirely submerged, which take the sea floor as a Posses = 0 boundary condition. 318 319 In outcropping grid cells, any decrease in fluid volume in the course of a time step is offset by a corresponding increase in the air volume. Any air 320 321 volume in a time step, in turn, is vulnerable to replacement by precipitation from the atmosphere (recharge). A maximum possible 322 323 groundwater recharge rate of 0.1 meters per year is imposed, but in 324 these simulations, given the very shallow land surface elevation gradient, 325 the capacity for subsurface flow to accommodate recharge is much lower than this, of order a few millimeters per year. The depth of the water 326 327 table in these simulations corresponds very closely with the depth of the 328 sediment surface, as any unsaturated soil produced by lateral flow is 329 replenished by recharge.

**2.2.4 Canyons** 

The model as described so far represents a laterally homogeneous slab, a poor approximation for hydrology above sea level because of the formation of canyons and river networks in a real drained plateau. The depth of the water table in a river canyon is depressed, relative to the surroundings, to the depth of the canyon. The water table is higher in between the canyons because of recharge, and the difference in head drives lateral flow, the canyons acting to drain the sediment column.

The model formulation has been altered to represent this mechanics in a simplified way. Rather than expand the model into the full third dimension, the 2-D field of the model is held to represent the sediment column at a hypothetical ridge crest, as altered by an adjacent canyon. The canyon elevation is represented by  $z_{\text{canyon}}$ , and its width by a scale  $\Delta y_{\text{canyon}}$ . A cross-column flow velocity  $v_{\text{Darcy,j}}$  is calculated as

$$344 \qquad v_{\text{Darcy},j} = \frac{k_{\text{h},j}}{\mu} \frac{\left(P_{\text{head},\text{canyon}} - P_{\text{head}}\right)}{\Delta y_{\text{canyon}}}$$

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where  $P_{\text{head,canyon}}$  is the pressure head as a function of depth in the hypothetical canyon, calculated assuming that the water table outcrops at  $z_{\text{canyon}}$ , and that the temperatures in the sediment column have adjusted to the formation of the canyon, such that the near-surface geothermal gradient is the same between the hypothetical canyon and

the bulk sediment column. The lateral "drainage" flow  $(v_{\text{Darcy},j})$  drives 350 351 vertical velocities by continuity. The horizontal distance scale  $\Delta y_{canvon}$  is somewhat arbitrary and difficult to 352 constrain, given that in the reality of river networks the distance to the 353 354 nearest canyon from any point in the domain is likely to be a function of 355 altitude, distance from the coast, and time. Another poorly resolved 356 factor is the depth of the canyon. In reality, canyons cut into a plateau 357 following a dynamic that erosion is proportional to slope, but stops at sea 358 level. As a simplification the model is set to hold the canyon depth at 359 current sea level. 360 The canyon mechanism accelerates the freshening of the sediment 361 column by providing a pathway for the escape of the salt water, although 362 it was found that the net effect in the model is not dramatic (results 363 shown below), in part because the canyon drainage mechanism only acts 364 on pore fluids above sea level, while the hydrological freshwater pumping 365 mechanism reaches much deeper than sea level. In the real fractal 366 geometry of canyons, the spacing between canyons across a plain is 367 similar to the width of the plain (length of the canyons), so the Base 368 simulation assumes a canyon width of 100 km, based on the 100+ km 369 width scale of the continental shelf. 370 2.3 Permafrost 371 The ice model is based on an assumption of thermodynamic equilibrium, in 372 which the heat content of the cell is distributed between the pure ice. 373 hydrate, and brine phases, and the salinity of the brine drives a freezing 374 point depression to match the local temperature. The ice content in a grid 375 cell relaxes toward equilibrium, quickly enough to approximate an 376 equilibrium state through the slow temperature evolution in the model 377 (which neglects a seasonal cycle at the surface), but slowly enough to 378 avoid instabilities with other components of the model such as fluid flow and methane hydrate formation. A limiter in the code prevents more 379 380 than 99% of the fluid in a grid cell from freezing, but the thermodynamic 381 equilibrium salinity is used to calculate, for example, the stability of 382 methane hydrate, to prevent the numerical limiter from affecting the 383 thermodynamic availability of water to drive chemical reactions. 384 numerical treatment of ice formation in the pore space in any given time 385 step is based on the energy constraint that the maximum extent of ice

freezing, for example, would release enough latent heat to raise the temperature of the grid cell to the freezing point.

388 Freeze<sub>max</sub> = 
$$\frac{\Delta T \cdot C_p \cdot V}{H_{\text{fusion}} \cdot MW_{\text{ice}}}$$

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where  $\Delta T$  is the temperature difference from equilibrium (which is a function of salinity and pressure), Co is the heat capacity of the grid cell overall (linear average of the material in the cell), H<sub>fusion</sub> is the heat of fusion of ice, and MW<sub>ice</sub> is the molecular weight. The actual extent of freezing imposed in a time step is taken to be 1% of Freeze, to prevent oscillations of the salinity, hence the freezing temperature, of the cell-Although the kinetics of freezing imposed by this formulation are much slower than would be realistic (given a time step of 0.02 years), the model forcing, due to changes in surface temperature, sea level, and subsurface diffusion of heat, is slow enough that the model relaxes nearly to an equilibrium condition, where the elevated salt concentration in the pore water drives a freezing point depression that matches the local temperature. A limiter in the code prevents more than 99% of the fluid in a grid cell from freezing, but the thermodynamic equilibrium salinity is used to calculate, for example, the stability of methane hydrate, to prevent the numerical limiter from affecting the thermodynamic availability of water to drive chemical reactions.

This model formulation implies that the salinity of pore fluid in subfreezing conditions (the permafrost zone) is independent of the original salinity of the bulk sediment column, but is rather determined only by the freezing-point depression implied by the temperature. If the original column is relatively fresh, there will be a smaller volume of pore fluid at a subfreezing temperature than if it is originally salty (see for example Figure 4 in [Nicolsky and Shakhova, 2010][Nicolsky and Shakhova, 2010]], but the activity of the water (a correlate of the salinity) is set by the temperature and the thermodynamics of pure ice, which are the same in the two cases. Layers of high-salinity unfrozen brines called cryopegs [Gilichinsky et al., 2005; Nicolsky et al., 2012][Gilichinsky et al., 2005; Nicolsky et al., 2012] are consistent with this formulation.

## 2.4 Thermodynamic competition between ice and hydrate

The high salinity (low activity of water) in the permafrost zone has the practical impact of excluding methane hydrate from permafrost soils that

421 are significantly colder than freezing. The thermodynamics are illustrated 422 in Figure 1. When the system consists only of ice and fluid phases, the 423 equilibrium salinity S<sub>eq</sub> increases with decreasing temperature below 424 freezing (Figure 1a, left). Above the melting temperature, ice is unstable, 425 as indicated by the nonzero values of the disequilibrium temperature. 426  $\Delta T_{eq.ice} = T - T_{eq.ice}$  in contours, even in zero-salinity water (right). For a 427 system consisting of only the hydrate and fluid phases (assuming that ice 428 formation is disallowed, and also gas saturation for methane) (Figure 1b), 429 the behavior is similar but with an added pressure dependence due to the 430 compressibility of the gas phase. When both solid phases are allowed, 431 the overall equilibrium salinity will whichever is higher between Segice and 432 Sea hydrate. Whichever phase can seize water at its lowest activity (highest 433 salinity) will be the stable phase. The salinity of the brine excluded from that phase will be too high to permit the existence of the other solid 434 435 phase at that temperature. The contours show ΔT<sub>eq</sub> for hydrate (solid) 436 and ice (dashed), which are also plotted in color in Figures 1d and e. This is illustrated in Figure 1d, in colors of  $\Delta T_{eq}$  hydrate and contours of the 437 excess salinity relative to hydrate equilibrium, Smax - Sen hydrate. Hydrate is 438 only stable when  $\Delta T_{eq}$  hydrate is zero (purple color). Under permafrost 439 conditions of low pressure and low temperature (upper left corner),  $\Delta T_{eq}$ 440 441 hydrate is greater than zero, indicating that hydrate is unstable, coinciding 442 with the salinity forcing from the ice, in overlain contours. A similar 443 exclusion of ice in part of the hydrate stability zone is seen Figure 1e, but 444 this would only happen in nature in conditions of unlimited methane. The 445 resulting phase diagram for ice and methane hydrate is shown in Figure 446 1f. Hydrate stability is suppressed in the permafrost zone by this 447 thermodynamic mechanism. 448

## (results shown below).

449 Permafrost formation has several impacts on the methane cycle in the 450 model. Biogenic methanogenesis is assumed stopped in the ice fraction 451 of a grid cell (which approaches unity but never reaches it in the model, 452 due to exclusion of salt into brine). Bubble transport in the model 453 balances bubble production, driven by a small and not very well 454 constrained standing bubble concentration within the pore space. It is 455 generally assumed [Shakhova et al., 2010b] [Shakhova et al., 2010b] that 456 permafrost inhibits gas transport through the sediment column, both 457 based on sediment column carbon and hydrogen budgets [Hunt, 458 1995][Hunt, 1995] and on the tight seal provided by the ice complex.

- 459 The seal provided to Arctic lakes, which can drain overnight if the seal is
- breached, also lends credence to this idea. In the model, this effect was
- simulated by stopping gas transport completely when a grid cell exceeds
- 462 50% ice fraction (with sensitivity runs assuming 10%, 30%, 70%, and
- 463 <u>90%)</u>.

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## 2.54 Atmospheric Methane Fluxes

- 465 Bubbles emerging from the sediment column into the water column of the
- ocean may dissolve in the water column, or they may reach the sea
- 467 surface, a direct methane flux to the atmosphere [Westbrook et al.,
- 468 2009][Westbrook et al., 2009]. In the model, bubble dissolution in the
- 469 water column is assumed to attenuate the bubble flux according to the
- 470 water depth with an e-folding attenuation scale of 30 meters [Gentz et
- 471 <u>al., 2014; Portnov et al., 2013; Westbrook et al., 2009][Gentz et al.,</u>
- 472 2014; Portnov et al., 2013; Westbrook et al., 2009]. In reality, a low-flux
- gas seep, producing small bubbles, will probably not reach as far into the
- 474 water column as a 30-meter scale height, while a faster seep can reach
- 475 further. Methane dissolved in the water column, in reality, may survive
- 476 oxidation (time constant of about a year), and degas to the atmosphere,
- 477 but this possibility is not included in the model. For land grid points
- 478 (exposed to the atmosphere by lowered sea level), any upward bubble
- 479 flux at the sediment surface is assumed 100% released to the
- 480 atmosphere. The model neglects methane oxidation in soils, as well as
- 481 many other terrestrial processes such as thaw bulbs beneath bodies of
- 482 water [Walter et al., 2006] [Walter et al., 2006], and the seasonal cycle
- 483 of melting and thawing in the surface active layer. In short, the methane
- 484 fluxes to the atmosphere computed from the model runs are crude, and
- 485 underlain by a sedimentary methane cycle with large uncertainties,
- 486 intended to capture the main sensitivities to various processes rather
- 487 than to provide strong quantitative constraint to the fluxes in the real
- 488 world.

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## 2.65 Comparison with Previous Models

- 490 The dynamics of the permafrost layer, and its present state, have been
- 491 extensively modeled within detailed maps of the crust and sediment
- 492 structure [Gavrilov et al., 2003; Nicolsky and Shakhova, 2010; Nicolsky et
- 493 al., 2012; Romanovskii and Hubberten, 2001; Romanovskii et al.,
- 494 2005] Gavrilov et al., 2003; Nicolsky and Shakhova, 2010; Nicolsky et
- 495 *al.*, 2012; *Romanovskii and Hubberten*, 2001; *Romanovskii et al.*, 2005].

496	The crust underlying the continental shelf area has been alternately rising
497	subsiding and subsiding rising in blocks called horsts and grabens
498	[Nicolsky et al., 2012][Nicolsky et al., 2012]. The sediment cover on the
499	horsts grabens is much thicker than it is in the grabenshorsts.
500	SpongeBOB, an idealized two-dimensional model, does not address this
501	complexity, but the thickness of the sediment cover on the shelf ranges
502	from 5 – 10 kilometers, reminiscent of the horsts grabens (subsiding
503	blocks). A thin sediment column would not reach the temperature
504	required for thermogenic methane production. The rates of thermogenic
505	methane production are not predicted or constrained by the model,
506	because of the different depositional histories of the sediment columns.
507	However, we can gauge the sensitivity of the methane cycle in the near-
508	surface sediments to thermogenic methane production by scaling the
509	model-predicted rate (by factors of 10 and 100).
510	Methane hydrate modeling has been done in the Arctic applied to the
511	Siberian continental slope [Reagan, 2008; Reagan and Moridis, 2009;
512	Reagan et al., 2011][Reagan, 2008; Reagan and Moridis, 2009; Reagan et
513	al., 2011, but only one calculation has been done in the context of
514	permafrost formation [ <i>Romanovskii et al.</i> , 2005][ <i>Romanovskii et al.</i> ,
515	2005], as found on the shelf. Romanovski [2005][2005] modeled the
516	extent of the methane hydrate stability zone through glacial cycles, but
517	apparently based the calculations on marine salinity values when
518	calculating the stability of hydrate, while I argue that in sub-freezing
519	conditions (in the permafrost zone) the only water available for hydrate
520	formation will be in a saline brine that would be in equilibrium with ice at
521	the local temperature. SpongeBOB therefore This formulation restricts
522	hydrate stability from the permafrost zone to greater depth below the
523	sea floor than predicted by Romanovski [2005][2005]. In the Mackenzie
524	Delta, hydrate was detected in a core drilled into onshore permafrost soils
525	[Dallimore and Collett, 1995][Dallimore and Collett, 1995], but only at
526	depths greater than 300 meters, near the base of the permafrost where
527	the temperature is close to freezingzone.
528	3. Results
529	3.1 Initial Spinup
530	The point of the spinup phase is to generate an initial condition for the
531	glacial cycle simulations. The more usual approach in modeling hydrates
532	is to start with an ad-hoc initial condition [Reagan, 2008; Reagan and

533	Moridis, 2009; Reagan et al., 2011]. For SpongeBOB the model state at
534	any time is the result of the time-history of sedimentation, which is driven
535	by the time-evolving depth of the sea floor, and interacting with isostatic
536	adjustment of the crust. The simplest way to generate an initial condition
537	in the model without a startup transient is to spin the model up from
538	bedrock at low resolution. Because of the over-simplicity of the tectonic,
539	sea level, and sedimentation forcing of the spinup phase, its POC
540	concentrations and methane production rates do not constrain those of
541	the real Siberian shelf. The sensitivity of the glacial methane cycles to
542	methane production rates will be evaluated by scaling the model
543	methanogenesis rates from the spinup result. The model setting was
544	grown for 62 million years of model time. The initial spinup used a
545	relatively coarse resolution as shown in Figure 2a.
546	The model setting was grown for 62 million years of model time. The
547	initial spinup used a relatively coarse resolution as shown in Figure 1a.
548	For the glacial / interglacial experiments, the initial condition was
549	interpolated to a higher resolution grid in the vertical, as shown in Figure
550	1b2b. Particulate organic carbon (POC) concentrations are highest just
551	off the shelf break (Figure 23), because this is where most of the
552	sediment is deposited, and because the sedimentary material is richest in
553	POC in shallow ocean water depths [Archer et al., 2012][Archer et al.,
554	2012]. The unchanging sea level in the spinup period kept the sediment
555	surface from outcropping, resulting in nearly uniform marine salinity
556	throughout the model domain (Figure 3a4a). Methane concentration
557	(Figure 54a) closely mirrors the solubility of dissolved methane, resulting
558	in near saturation concentrations through most of the model domain
559	(Figure <u>54b</u> ). As in the previous model simulations [ <i>Archer et al.</i> ,
560	2012][Archer et al., 2012], the imposition of permeable channels has a
561	strong effect on the chemistry of the permeable grid cells (Figure 54d),
562	although the impact on the integrated model behavior, such as the
563	methane flux to the atmosphere, was small in these simulations, so for
564	clarity model results neglecting these channels will be shown.
565	3.2 Impact of Freshwater Hydrology
566	When sea level drops such that the land surface of the sediment column
567	outcrops to the atmosphere, the pore fluid becomes subject to the
568	pressure head driving it seaward, and to fresh water recharge from
569	precipitation. The pressure head forcing and the buoyancy of the
570	sediment fluid column combine to create a mechanism to excavate

salinity from the upper sediment column. Initially after sea level fall, there is a pressure head gradient extending throughout the sediment column, provoking lateral flow at all depths. As the pore fluid at the surface is replaced by fresh runoff, the lighter density of that fluid tends to diminish the pressure head gradient in the deeper sediment column. The deeper pressure gradient and flow approach zero as the fresh water lens in the outcropping region approaches an isostatic equilibrium condition known as the Ghyben-Herzberg relation [Moore et al., 2011] [Moore et al., 2011], in which each meter elevation of the water table is compensated for by about 40 meters of fresh water below sea level, determined by the difference in densities of fresh and salt water.

In an attempt tTo create this condition within the model, two simulations are presented in which sea level was decreased by 30 and 120 meters, respectively, and held there for millions of years (Figure 65). The 30-meter drop experiment produced land outcrop in about 1/4 of the model domain, with the predicted equilibrium Ghyben-Herzberg halocline reaching about 1200 meters maximum depth. The model salinity relaxes into close agreement with the predicted halocline, lending support to the model formulation for density, pressure head, and fluid flow. As time progresses further, the outcropping land surface subsides (there is no land deposition in this scenario), until it drops below the new lowered sea level value after about 2.5 Myr. The sequence of events leaves behind a persistent fresh water lens below sea level.

Variants of this experiment were done with differing values of the lateral distance to drainage canyons in the model, which provide a pathway for fluid loss in sediments above sea level. When a hypothetical canyon is located 10 km from the SpongeBOB slab, the model salinity approaches equilibrium on an e-folding time scale of about 4300 kyr (Figure 76). When the canyon is 100 km distant or nonexistent, the equilibration time scale is about 6500 kyr. Based on the idea that canyons of order 100 km long should be about 100 km apart, The rest of thethe Base simulations in this paper assumes canyon spacing of were done using the relatively low-impact 100 -km-canyon spacing.

When sea level is lowered by 120 m, the sequence of events is similar, except that the pressure head is so high that to satisfy the Ghyben-Herzberg relation would require fresh pore waters at many kilometers depth, even deeper than bedrock on the "continental" side of the model domain. Because of the low permeability of the deepest sediment

609 column, the freshwater pumping groundwater mechanism is unable to 610 reach these deepest pore waters, which therefore remain salty. The time 611 scale for establishing a significant freshening of the upper kilometer of 612 the sediment column is still on the order of 100-500 kyr, and the 613 subsequent subsidence time of the sediment column in the model, until it 614 drops below the new lowered sea level, takes about 10 Myr. In both 615 cases, subsidence of the exposed sediment column prevents the 616 sediment surface in the model from remaining above sea level indefinitely 617 (without land deposition).

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The sequence of events leaves behind a fresh water lens below sea level that persists in the model for millions of years (Figure 6). Groundwater flow, driven by the pressure head, provides an advective means of pumping fresh water into the subsurface sediment column that has no counterpart for salty ocean water. The model lacks the mechanism of salt fingering, which can enhance the diffusion of salt from above into a fresh water agufer [Kooi et al., 2000]. However, higher-resolution models of smaller domains that accounted for salt fingering also show a time asymmetry, with faster fresh water invasion on sea level drop than salt invasion on sea level rise [Lu and Werner, 2013; Watson et al., 2010]. As the size of the domain increases with increasing sea level change, advective processes such as hydrological flow should become even more dominant over diffusive processes such as salt fingering. The recent discovery of vast freshwater aguifers on global continental shelves [Post et al., 2013], persisting since the time of lowered sea level 20,000 years ago, and the lower-than-marine salinities of the pore waters measured in submerged surface Arctic sediments (summarized by [Nicolsky et al., 20121) are also consistent with the existence of a fresh-water hydrological pump which has a significant impact on sediment column salinities. The hydrological pumping generates a low-methane plume that also persists for millions of years in the model (Figure 8). Two states, called "prefreshened" and "pure marine", serve as end-member initial conditions for glacial / interglacial simulations (Figure 4b), to evaluate the sensitivity of the model glacial cycles to the initial salinity of the sediment column.

The interesting and salient observation from both of these simulations is the persistence of the fresh pore water lens after its resubmergence in the salty ocean. Groundwater flow, driven by the pressure head, provides a means of pumping fresh water into the sediment column that has no

647	counterpart for salty ocean water. Submerged pore waters near the
648	sediment surface tend to pick up salt by diffusion in the model, but this is
649	much slower than the ground water pumping mechanism, and its depth
650	reach is much shallower. Any exposure of the continental shelf to the
651	atmosphere will tend to ratchet down the salinity of the pore waters on
652	the shelf, consistent with the recent discovery of vast freshwater
653	aquifers on global continental shelves [Post et al., 2013], and in particular
654	the lower-than-marine salinities of the pore waters measured in
655	submerged surface Arctic sediments (summarized by [Nicolsky et al.,
656	2012]). The impact of the groundwater pump on the methane cycle in
657	the upper sediment column is profound and long-lasting in the model
658	(Figure 7).
659	The final, resubmerged state of the 120 meter drop experiment is taken
660	as the initial condition for glacial / interglacial cycle forcing, described in
661	the next section, through which the behavior of this "prefreshened" initial
662	condition is compared with that from a "pure marine" (still salty
663	throughout) initial condition (Figure 3b).
664	3.3 Glacial Cycles
665	3.3.1 Setup and Forcing
666 667 668 669 670 671	Beginning from an entirely submerged initial condition, the model is subjected to 100-kyr sawtooth cycles of sea level ranging between –120 to +20 meters from the initial sea level (starting at –120 for prefreshened, 0 for pure marine) (Figure <u>98a</u> ). We show a suite of model simulations intended to isolate the impacts of various component processes.
672	The model scenarios and sensitivity studies are summarized in Table 1.
673	The simplest scenario (SL) varies the sea level while keeping the air and
674	water temperatures time-invariant. The sea-level air temperature is
675	maintained at 0 °C. This simulation is nearly permafrost-free, with a small
676	exception where the altitude of the sediment surface is much higher than
677	sea level (due to the lapse rate <u>in the atmosphere</u> ). There is no
678	deposition of sediment above sea level in this simulation. Permafrost
679	formation is added in simulation GL, in which the air temperature ramps
680	down to -16 °C at sea level, linearly with the glacial sea level fall (Figure
681	98b). In the ocean, shelf waters are always -1.8 °C, but an interglacial
682	subsurface temperature maximum of 1 °C at 200 meters decreases to

683 -1.8 °C during glacial times. Deposition of organic-rich sediments when 684 the surface is exposed to the atmosphere (Yedoma: represented as 685 accumulation of 10 meters in 100 kyr, with 30% POC) is added in 686 scenarios SL+LD and GL+LD (LD for land deposition). The atmospheric 687 temperature impact of a global warming scenario (GW) is also shown in 688 Figure 98b, beginning at 400 kyr, and compared with an extended-689 interplacial control forcing (Ctl). The potential impact of geologic-time scale sea level rise is added to the global warming scenario in simulation 690 691 GL+SL. 692 Other model sensitivity runs used varying values of the thermogenic and 693 biogenic methane production rates, the geothermal temperature gradient. 694 Several altered-physics runs were done, one adding vertical permeable 695 channels, one disabling horizontal flow, and several to evaluate the impact 696 of ice formation on methane hydrate stability. 697 3.3.2 Salinity and Ice 698 In the "prefreshened" initial condition (Fr), millions of years have elapsed 699 since the previous exposure of the sediment to hydrological forcing, but a 700 core of fresh water remains. Salinities near the sediment surface have 701 grown saltier due to diffusive contact with the salty oceanseawater 702 (Figure 109, left). A fully marine initial condition (Mar) (Figure 109, 703 right) was initialized from the unfreshened case, in which sea level was 704 held at a fixed value throughout the 65 Myr spinup of the sediment 705 column. The salinities are nearly uniform in this case. 706 When the sediment surface is re-exposed to the atmosphere during an 707 interval of sea level, in the absence of ice formation (simulation SL), the 708 surface layer tends to freshen relatively quickly due to the hydrological 709 forcing, but a subsurface salinity maximum persists (Figure 109c and d). 710 However, if the air temperatures are cold enough to form ice (simulation 711 <u>GL</u>), surface salinities in the model increase to up to nearly 190 psu, in 712 both prefreshened and pure marine cases (Figure 109e and f). By the 713 next interglacial time (Figure 109g and h), ice near the sediment surface has melted enough for near-surface pore waters to reach relatively low 714 715 salinities. 716 3.3.3 Pressure and Flow 717 The effect of the sea level and permafrost forcing on the pressures and 718 flow velocities are shown in Figure 1+10. On a spatial scale of the entire

model domain (Figure 110, left), the highest driving pressures are found at the base of the sediment column, underneath the region of maximum sediment accumulation (the depocenter just off the shelf break). Changes in sea level drive large fluctuations in the pressure head (contours) extending to bedrock. In the near-surface continental shelf (Figure 110, right), the driving pressure variations are dominated by the pressure head driven by sea level changes. The formation of permafrost (GL. Figure 110 e and f) seals the upper sediment column to fluid flow. When sea level rises again, in the model configuration including permafrost, there is a strong pulse of downward flow following partial melting of the permafrost (Figure 110 h). It is possible that this flow, which lasts a few thousand years, is an artifact of the elastic model configuration, in which the release of a load (by submergence of the upper sediment column into the ocean) provokes the expansion of pore spaces in the sediment. The anomalous flow, integrated over its duration, could displace the pore fluid by about 40 meters, which is less than one grid cell. However, tThe model configuration without the sealing effect of permafrost (SL) does not show this pulse of invasive flow on sea level rise.

## 3.3.4 Methane Fluxes Cycle

There are multiple ways in which the glacial cycles of sea level and air and water temperature might impact the flux of methane to the atmosphere. Submergence in the ocean is one modulating factor, because the emerging bubbles dissolve in the ocean rather than reaching the atmosphere. Another factor is the deposition of high-POC surface soils during low sea level stands, and its exposure to degradation later when the permafrost soils melt. A third factor is permafrost, impeding gas and fluid flow and excluding dissolved methane and salt from ice formation. The impacts of these processes are assessed by comparing the results from model configurations with and without each process in question.

The impact of phase competition between ice and hydrate is shown in

Figure 12. In the Base scenario (Figure 12a and c) hydrate stability is excluded from the permafrost zone as described above and in Figure 1. Preventing ice from forming in an altered-physics simulation (+ No Ice) decreases the fluid-phase salinity relative to the Base simulation, and allows the methane hydrate stability zone to nearly reach the sea floor (Figure 12b and d), during strongest glacial conditions. Another altered-physics simulation was done in which ice is allowed to form, but not

757 <u>affect the salinity as it drives methane hydrate stability (which was hard-</u>

758 wired to marine salinity). Methane hydrate is still unstable in the

759 permafrost zone through most of the simulation (see movie files in

760 <u>supplemental material</u>), indicating that thermal interaction must also have

761 <u>a strong impact on methane hydrate stability in the permafrost zone.</u>

762 The evolution of the dissolved methane disequilibrium condition (CH<sub>4</sub> /

763  $CH_{4 sat}$ ) is shown in Figure 134. At the initiation of the glacial cycles,

764 methane is undersaturated in near-surface sediments on the continental

shelf, by diffusive contact with the methane-free ocean upper boundary

condition. In the prefreshened sediment column scenario (Fr), methane

767 concentrations in the depth range of 100-1000 meters are lower than in

768 the marine case (Mar, Figure 134b), due to the ventilation by the

769 hydrological pump (Figure 134a). Further freshening of the pore waters

in the ice-free case (SL+LD) tends to deplete methane in the upper

771 sediment column (Figure 131c-e), while methane exclusion from the

772 permafrost ice leads to supersaturation in simulation GL+LD (Figure 134

773 f-h). The hydrate stability zone is somewhat expanded in the

prefreshened sediment column relative to the marine case (Figure 134 g

775 vs. h, heavy black contour).

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776 Figure 142 shows snapshots of various aspects of the shelf carbon cycle.

777 <u>beginning from a prefreshened initial condition</u>. Sections of POC

concentration in Figure 142, left show the accumulation of POC-rich

Yedoma deposits on land (Figure 142 g and j). The rate of methane

production in the model (Figure 142, right) depends on temperature and

781 organic carbon age, but it is also attenuated by permafrost formation in

782 the model, scaling to zero in the completely frozen case. Methanogenesis

783 rates are near zero in the permafrost zone during glacial time (Figure

784 142h), but partially recover during interglacial time (Figure 142k) even

785 though permafrost is still present.

A zone of methane hydrate stability exists below the permafrost zone

787 when permafrost is present, and some methane hydrate accumulates in

788 that zone. The highest pore-fraction values are found near the

789 continental slope, where the shelf stability field outcrops within the slope

790 depocenter. Dissolved methane concentrations exceed saturation within

791 the stability zone in the model (Figure 134), but the accumulation of

methane hydrate (Figure 142, right) is limited by the slow rate of

793 methane production.

794 TA-time series plots of the inventory of methane as hydrate on the shelf 795 <u>areis</u> shown in Figure 153. The integration cuts off at x=560 km to 796 exclude the sediment depocenter on the continental slope. Hydrate 797 inventories reach maximum values during deglaciations. There is more 798 hydrate when the pore water is fresher, and there would be more if ice 799 were excluded from forming (Figure 15a). The hydrate inventory is much 800 more sensitive to thermogenic methane production, deep in the sediment 801 column, than Yedoma deposition (Figure 15b). The impact of the 802 geothermal heat flux is to change the depth of the bottom of the hydrate 803 stability zone (Figure 12 e and f), but the impact is small on the hydrate 804 inventory, unless the temperature gradient is so low that hydrate persists 805 through the entire glacial cycle (Figure 15c). The hThe inventory of methane in hydrates rises to maxima of ~25 Gton C, right at the 806 807 deglaciation. However, this methane hydrate cycling is not as significant 808 to the larger carbon cycle as one might expect, because the hydrate is 809 formsed from the dissolved methane pool, (which exceeds 1000 Gton C 810 in shelf porewaters of the model), rather than interacting directly with the 811 atmosphere.

812 The impact of the glacial cycles on the methane pathway to the 813 atmosphere in the model is shown in Figure 164. When sea level is high, 814 the efficiency of bubble transport across the sediment-water interface 815 reaching the atmosphere ranges from about 75% near the coast to about 816 10% at the shelf break (Figure 164a). Most of the methane flux from the 817 sediment is located just off the shelf break (Figure 164e), where the 818 escape efficiency is low, so not much methane makes it to the 819 atmosphere during the interglacial. During glacial times, the sediment 820 column is exposed to the atmosphere, and the escape efficiency in the 821 model is 100% (Figure 164b). Permafrost inhibits the terrestrial methane 822 flux (Figure 164i) relative to the case without permafrost (Figure 164f). 823 During some the some of the deglaciations, the release of pent-up gas by 824 permafrost degradation leads to a deglaciation spike of excess methane flux to the atmosphere (Figure 164j-k relative to 164q-h). 825

Time series plots of the <u>major fluxes of the</u> methane cycle on the continental margin are shown in <u>Figure 175</u>. The methanogenesis rates in the model output are in units of moles per meter of coastline, since it is a 2-D model. We scale this up to the Siberian continental margin by assuming a width of 1,000 km. The area of the shelf is then 5 · 10<sup>11</sup> m<sup>2</sup>,

831 roughly comparable to the real shelf area of 460,000 km<sup>2</sup> [Stein and Fahl. 832 2000][Stein and Fahl, 2000] 833 834 The biological rate of methane production on the continental shelf evolves through time in Figure 175b. Yedoma deposition (case SL+LD) 835 836 tends to slowly increase the total shelf respiration rate in the model, 837 relative to a case with no land deposition (case SL). The formation of 838 permafrost, during glacial periods of case GL+LD, attenuates 839 methaneogenesis by inhibiting biological activity in the frozen soil. 840 841 The continental shelf methane cycle is summarized for four different model scenarios in Figure 15 c-f. The solid regions in Figure 17 c-h are 842 843 cumulative methane sinks for six different model scenarios, plotted 844 underneath red lines showing biogenic methane production. In time 845 average, where sinks balance sources, the colored areas should fill up the 846 region below the red line. 847 Trapping of methane by impermeable permafrost leads to a spike of 848 methane fluxes at the ends of deglaciations in simulations with permafrost (Figure 17 c and e). The spikes happen as sea level 849 approaches its highest extent, stifling the offshore groundwater flow by 850 851 decreasing the pressure head, but early in the interglacial time while 852 permafrost is the most intact. The spikes are stronger for the first glacial 853 cycles than the last, apparently due to long-term adjustment of the 854 methane cycle on the shelf (a growing together of the production rate 855 (red lines in Figure 17 c-f) and the various methane sinks (colored areas). 856 Permafrost formation blocks methane emission during times of low sea 857 level. This can be seen in the collapse of the blue regions in Figure 17 c 858 vs. d and e vs. f during times of low sea level. Blocking horizontal flow 859 disrupts offshore flow, the only significant methane sink on the shelf 860 during glacial periods (Figure 17h), resulting in somewhat higher deglacial 861 spikes of methane emission than predicted by the models including 862 transport. There is no direct link between ice fraction and methane 863 oxidation in the model, which is driven only by coexisting concentrations of sulfate and methane, but the rate of methane oxidation also drops to 864 865 negligible during glacial times in the simulations with permafrost (grey in

Figure 17 c and e). The absolute rates of methane loss differ between 866 867 the Prefreshened vs. Marine initial conditions, but this is in part due to differences in the width of the continental shelf between the two 868 869 simulations. The patterns of the methane cycle are very similar, however, 870 between the two cases, and also not much affected by the imposition of 871 permeable vertical channels (Figure 17g). 872 Results are shown with and without permafrost formation, and with and 873 without the hydrological pre-freshening in the initial condition. Solid regions are cumulative methane sinks. The red lines are the methane 874 875 source from biological processes. In steady state, where sinks balance 876 sources, the colored areas should fill up the region below the red line. 877 By model design, permafrost formation inhibits methane loss from the shelf sediments as bubbles. This can be seen in the collapse of the blue 878 region in Figure 15 c and e during times of low sea level. There is no 879 direct link between ice fraction and methane oxidation in the model, which 880 is driven only by coexisting concentrations of sulfate and methane, but 881 the rate of methane oxidation also drops to negligible during glacial times 882 883 in the simulations with permafrost (grey in Figure 15 c and e). 884 The only remaining sink for methane from the continental shelf sediment column is offshore subsurface ground water flow carrying dissolved 885 886 methane (brown in Figure 15 c and e). Offshore transport is also significant or dominant in the runs without permafrost formation (Figure 887 15 d and f). During interglacial intervals, there is a small onshore flow of 888 dissolved methane into the continental shelf region, for example at time 889 890 0, which is indicated in the plot as a negative starting point for the grey region, representing methane oxidation in the sediment column. 891 892 atmospheric fluxes 893 Trapping of methane by impermeable permafrost leads to a spike of methane fluxes at the ends of deglaciations in simulations with 894 permafrost (Figure 15 c and e). The spikes happen as sea level 895 approaches its highest extent, stifling the offshore groundwater flow by 896 decreasing the pressure head, but early in the interglacial time while 897 permafrost is the most intact. The spikes are stronger for the first glacial 898 cycles than the last, apparently due to long-term adjustment of the 899 methane cycle on the shelf (a growing together of the production rate 900 901 (red lines in Figure 15 c-f) and the various methane sinks (colored areas).

902	<u>Fluxes of methane to the atmosphere the impacts of salinity, </u>
903	prefreshened versus marine initial conditions, on the methane flux to the
904	atmosphere, are shown in Figure 186. In the absence of permafrost
905	(Figure 18 a and b), or assuming that bubble migration is blocked only if
906	the ice fraction exceeds 90%, a condition rarely attained in the model
907	(Figure 18e), the highest methane fluxes to the atmosphere are found
908	during glacial (cold) times, rather than warm interglacials. This is due to
909	dissolution of methane gas into the ocean when the sediment column is
910	submerged. When permafrost blocks methane gas fluxes in the sediment
911	column, the highest atmospheric fluxes are generally found during the
912	time of early sea level fall, when unfrozen sediment is exposed to the
913	atmosphere before it has a chance to freeze. The timing of the variations
914	in atmospheric flux through the glacial cycles is very sensitive to the
915	critical ice fraction for blocking gas transport (Figure 18e). Due to the
916	differing deposition history, the continental shelf is narrower in the marine
917	than it is for the pre-freshened initial condition simulations.
918	The impacts of the pore water salt inventory are most apparent during
919	the time of sea level fall, with permafrost formation (red lines). The
920	saltier sediment column takes about 20 kyr to choke off the methane flux
921	to the atmosphere (Figure 186a), while the pre-freshened sediment
922	column stops the methane flux more abruptly, in just a few thousand
923	years (Figure 1 <u>8</u> 6b).
924	Atmospheric emissions also scale with methane production rates,
925	generally maintaining the temporal patterns of emission as set by
926	permafrost and submergence in the ocean.
320	permanost and submergence in the ocean.
927	The generally lower fluxes in the shelf methane cycle for the marine case
928	are to some extent due to the differing areas of integration (spanning
929	460 vs. 560 km of shelf width). In general, the dominant player of the
930	glacial / interglacial cycles on the atmospheric release rate of methane in
931	the model is dissolution of methane bubbles rising through the water
932	column.
933	3.4 Anthropogenic Global Warming
934	The global warming (GW) scenario begins from a high sea-level interglacial
935	state, and raising the temperature following the climate impact of the
936	"spike and long tail" time distribution of a slug of new CO <sub>2</sub> added to the
937	atmosphere [Archer et al., 2009][Archer et al., 2009] (Figure 8). There

is a stage of fast atmospheric drawdown as CO<sub>2</sub> invades the ocean, but once the ocean, atmosphere, and land surface reach equilibrium (after a few hundred years), the CO<sub>2</sub> content of the entire biosphere begins to slowly relax toward an initial "natural" value, on time scales of hundreds of thousands of years, by weathering reactions with carbonate and siliceous solid rocks. The net result is a CO2 drawdown that can be expressed as the sum of several exponential functions in time, with time scales ranging from 10<sup>2</sup> – 10<sup>6</sup> years.

Changes in water column temperature are assumed equal to those of the atmosphere, following paleoceanographic reconstructions [Martin et al., 2002][Martin et al., 2002] and long-term coupled ocean / atmosphere circulation model experiments [Stouffer and Manabe, 2003][Stouffer and Manabe, 2003]. The first global warmingGW scenario imposes this temperature change on the water column, relaxing toward equilibrium with the atmospheric CO<sub>2</sub> trajectory with a time constant of 100 years.

The effect of sea level rise is added to create a second global warming scenario GW+SL. On time scales of thousands of years the sea level response to changing global temperature is much stronger than the sea level response over the coming century, as prominently forecast by the IPCC. Reconstruction of sea level and global temperature covariation in the geologic past (glacial time to Eocene hothouse) reveals a covariation of 10-20 meters per °C [Archer and Brovkin, 2008][Archer and Brovkin, 2008]. The global warming with sea level scenario assumes an equilibrium sea level response of 15 meters / °C, which it relaxes toward with a time constant of 1000 years.

The atmospheric methane fluxes, shown in Figure 197, increase in the global warming (GW) model run, as they also do in the control (Ctl) simulation, which is essentially an extended but unwarmed interglacial period. The permafrost melts on a time scale of about 10,000 years for the GW simulation, and about 50,000 for the Ctl. The rates of methane production, and flux to the atmosphere, both increase with the loss of the permafrost, if there is no change in sea level. However, the new methane flux comes not as a sudden burst, but rather as a slow transition toward a new, higher, chronic release rate. When sea level is also changed (GW+SL), bubbles dissolve in the water column, which more than counteracts the increase in methane flux due to the extended interglacial (Ctl) or warming (GW) scenarios.

975	3.5 Summary of Model Sensitivity Studies
976 977 978 979 980 981 982	Sediment Porewater Salinity. Ice freezes until the salinity of the residual brine brings about a freezing point depression equal to the in situ temperature. A saltier initial sediment column will reach this condition with a lower ice fraction, its melting is accelerated, and its hydrate inventory is lower (Figure 18). The equilibrium salinity in the permafrost zone is not affected by the salt inventory of the column, only the relative volumes of the solid and fluid phases.
983 984 985 986 987 988 989 990	Methane Production Rates. The atmospheric flux increases with increases in either shallow, biogenic methane production, driven by deposition of Yedoma, and thermogenic methane production in the deep sediment column (Figure 19). Biogenic methane is produced too shallow in the sediment column to impact the inventory of methane hydrate (Figure 15). The timing through the glacial cycles of atmospheric methane emissions from these scenarios parallel each other, because they are controlled in common by the transport-blocking effects of permafrost and sediment submergence in the ocean.
992 993 994 995 996 997 998 999	Geothermal Temperature Gradient. When the heat flux is higher, the temperature gradient is steeper, pivoting about the sediment surface temperature, which is set by the ocean. The base of the methane hydrate stability boundary gets shallower, while the top remains at about the same depth, resulting in a thinning of the stability zone (Figure 12). The hydrate inventory through the glacial cycles however is not much affected, unless the heat flux gets small enough for hydrate to persist through the glaciations (Figure 15).
1000 1001 1002 1003 1004 1005 1006 1007 1008 1009	lce vs. hydrate thermodynamic competition. When ice is included as a competing phase, it excludes methane hydrate from the low-pressure, very cold permafrost zone. The hydrate stability zone thins (from above and below in the model: Figure 12), and the hydrate inventory decreases (Figure 15). When ice formation is disallowed, the hydrate stability zone approaches the sediment surface during coldest glacial time, but by the time of an interglacial-based global warming climate perturbation, the stability zone boundary has retreated to several hundred meters below the sea floor, precluding a sudden hydrate dissolution response to a suddenly warming ocean.

1010 1011 1012 1013 1014	Permafrost inhibition of gas migration. When the ice fraction of the model exceeds a critical threshold, gas migration is blocked.  Changing the value of this threshold has a strong impact on the rates of methane emission during glacial versus interglacial times. This process is therefore a high priority for future model refinement.
1015 1016 1017 1018 1019 1020 1021 1022	Vertical flow heterogeneity. The chemistry of continental margin sediments in this model [Archer et al., 2012] showed a strong sensitivity to flow heterogeneity, achieved by increasing the vertical permeability of every fifth grid cell. In the configuration presented here, the impact of the channels is much smaller. The dynamics of this simulation are thermally driven, rather than by sediment deposition driving fluid flow in the continental margin case. Atmospheric methane fluxes are spikier when the channels are included, but the mean rate is not much changed.
1023 1024 1025 1026 1027	Ground water flow. Groundwater flow carries enough methane to be a significant sink during times of low sea level. However, disabling that flow has only subtle impacts on the other aspects of the methane cycle on the shelf. Spikes of methane emission during late deglaciation get somewhat more intense.
1028	
1029 1030	4. Implications of the Model Results for the Real Siberian Continental Margin
1031 1032 1033 1034 1035 1036 1037 1038 1039 1040	This is the first simulation of the full methane cycle on the Siberian continental margin, or any other location with embedded permafrost soils, including hydrate formation and transient fluxes. It is internally consistent, linking processes from the ocean, the sea floor, and the deep Earth, within constraints of sediment accommodation and conservation of carbon, through geologic time. As such it has some lessons to teach us about the real Siberian continental margin. However, many of the model variables are not well known, such as the methaneogenesis rates or soil permeabilities, meaning that in some aspects the model results are not a strong constraint on reality.
1041 1042 1043 1044	The absolute values of the methane inventories in the system, as hydrate and bubbles, are not well constrained theoretically. The rate of methane production in shallow sediments is not well characterized. In reality there might be some flux of methane from the crust, but this is not included in

1046 mechanistically poorly understood, therefore not well represented in the 1047 code, which affects the inventories of bubbles in the sediment. 1048 Ultimately the bubble concentration in the model reaches a rough steady 1049 state where production of methane gas balances its escape through the 1050 sediment column, but the steady state value from the model could be 1051 wrong. The model lacks faults, permeable layers, or the ability to "blow 1052 out", producing the sedimentary wipe-out zones observed seismically in 1053 the subsurface [Riedel et al., 2002] Riedel et al., 2002], and the 1054 pockmarks at the sediment surface [Hill et al., 2004][Hill et al., 2004]. 1055 On land, the model lacks seasonal melting of surface permafrost (to form 1056 the active layer) and the thaw bulbs underneath lakes and rivers. In the 1057 ocean, the intensity of water column dissolution of rising bubbles depends 1058 on the bubble sizes, which depend on the gas emission rate, ultimately 1059 driven by details of gas transport in the sediment, which are neglected in 1060 the model. These uncertainties all affect the flux of methane to the atmosphere, 1061 1062 which is therefore not well constrained by the model. However, the 1063 model is consistent with observations [Kort et al., 2012] Kort et al., 1064 2012], that the total atmospheric methane flux from the Siberian margin 1065 is a small fraction of the global flux of methane to the atmosphere, and 1066 thus represents only a minor climate forcing. The model would have to 1067 be pushed very hard (as would the measurements) to fundamentally 1068 change this conclusion. 1069 The model bubble flux to the atmosphere in the base case in analog present-day conditions is only 0.02 Tg CH<sub>4</sub> per year, which is an order of 1070 1071 magnitude lower than an estimate of the total methane emission rate 1072 from aircraft [Kort et al., 2012][Kort et al., 2012] of 0.3 Tq  $CH_4$  / yr. 1073 However, tThe model only accounts (crudely) for the bubble flux to the 1074 atmosphere, and does not include gas exchange evasion of methane from 1075 the water column, which could be significant. Concentrations of methane 1076 in the water column of 50 nM are common [Shakhova et al., 1077 2010a] Shakhova et al., 2010a], which, if they were unimpeded by sea 1078 ice, could lead to a flux from the region of 0.4 Tg CH<sub>4</sub> / yr (assuming a 1079 typical gas exchange piston velocity of 3 m/day). Methane fluxes into the water column range up to 0.4 Tg CH<sub>4</sub> / yr during times of relatively 1080 1081 high sea level. Once released to the water column, the fate of a methane 1082 molecule will depend on its lifetime with respect to oxidation, which could 1083 be up to a year in the open water column [Valentine et al.,

1084 2001] Valentine et al., 2001, versus its lifetime with respect to gas exchange, which for ice-unimpeded conditions would be just a few months 1085 1086 for a 50-meter deep water column. Thus the methane in bubbles 1087 dissolving in the water column has some chance of making it to the 1088 atmosphere anyway, depending on stratification in the water column and 1089 the extent of ice, and the gas exchange flux has the potential to be 1090 significant in the regional total flux. Shakhova et al [2010b] proposed that 50 Gton C as methane could erupt 1091 from the Arctic on a time scale of a few years. As has been 1092 1093 acknowledged, the model provides poor constraint on the standing stock 1094 of bubbles or methane hydrate in the sediment column, and neglects 1095 many of the mechanisms that could come into play in transporting methane quickly to the atmosphere, such as faults, channels, and 1096 1097 blowouts of the sediment column. However, one seemingly robust model 1098 result is the thermodynamic exclusion of methane hydrate from the 1099 permafrost zone, by competition for water between ice and hydrate. Thermodynamics does not control everything, especially at low 1100 1101 temperature, but kinetic inhibitions are more often found for nucleation 1102 steps rather than decomposition. To find an accumulation of 1103 "metastable" hydrate would also require some sort of transport 1104 mechanism of hydrate into the region where it is unstable, which does not 1105 exist. There is no reason to imagine that hydrate could form in situ when thermodynamic conditions are wrong for it. A kinetic inhibition of water-1106 1107 ice formation would work, but ice does not tend to super-cool in a dirty, nucleation-site-rich environment like sediments. Therefore it seems as 1108 1109 though methane hydrate should not be expected in sediment depths 1110 shallower than about 300 meters. A warming perturbation at the sea 1111 floor today will not reach this depth for hundreds or thousands of years. 1112 Could an abrupt methane release arise from release of trapped bubbles from melting ice? The model actually does produce a glacial cycle in 1113 bubble inventory, with changes exceeding 50 Gton over a cycle, 1114 1115 apparently driven by methane exclusion from ice formation (Figure 15). 1116 But the model does not deliver an abrupt release in response to 1117 anthropogenic warming for any of its sensitivity studies (Figure 18). 1118 Permafrost melting driven by deglacial sea level rise has already been going on for thousands of years. In this span of time a temperature 1119 1120 anomaly has diffused quite deep into the sediment column. In order for 1121 the abrupt temperature anomaly of global warming to further accelerate

1122 the ongoing ice or hydrate melting, it will have to diffuse down in the 1123 sediment column to where the ice still is. We would get a faster initial 1124 response to global warming if the transition from glacial to global warming 1125 sediment surface temperatures hadn't mostly happened thousands of 1126 vears ago. 1127 In the real world, geological features such as faults and permeable lavers 1128 dominate the methane cycle in the sediments. A continuum model such 1129 as this one predicts a smooth methane release response to a warming, growing in on some e-folding time-scale. A world dominated by features 1130 1131 that each represent a small fraction of the total methane reservoir will 1132 release methane more episodically, but the statistical distribution of the 1133 response in time should still show the e-folding time scale of the 1134 underlying driving mechanism, the diffusion of heat into the sediment column. The way to deliver 50 Gton of methane to the atmosphere is for 1135 1136 it all to be released from a single geologic feature pent up by ice. But 50 Gton of C represents a large fraction of all the traditional natural gas 1137 deposits on Earth (about 100 Gton C). The place to look for such a large 1138 unstable gas reservoir is in the field, not in this model, but until such a 1139 1140 thing is found it remains conjecture. 1141 One Another probably robust feature of the model is the dominant impact 1142 of sea level inundation of the sediment column on the atmospheric 1143 methane flux. The methane flux is highest during cold times, because sea level is low, rather than providing a positive climate feedback of releasing 1144 methane during warm (high sea level) intervals. There is a warming 1145 1146 positive feedback in the simulated future from climate warming, but it is 1147 much smaller than the impact of sea level changes in the past. The potential for future sea level change is much higher for the deep future, 1148 thousands of years from now, than the forecast for the year 2100, 1149 1150 because it takes longer than a century for ice sheets to respond to 1151 changes in climate. The model finds that for the future, if sea level 1152 changes by tens of meters, as guided by paleoclimate reconstructions [Archer and Brovkin, 2008][Archer and Brovkin, 2008], the impact of sea 1153 1154 level rise could overwhelm the impact of warming. The dominance of sea 1155 level over temperature in the model of this area is due to dissolution of 1156 methane in the water column, rather than a pressure effect on hydrate stability, which is generally a weaker driver than ocean temperature in 1157 1158 deeper-water settings [Mienert et al., 2005] [Mienert et al., 2005].

- Another seemingly robust model result is the exclusion of methane 1159 hydrate from the permafrost zone, by competition for water between ice 1160 and hydrate. This behavior implies that any hydrate on the continental 1161 shelf must be at least 300 meters below the sea floor in the sediment 1162 1163 column (as opposed to as shallow as 100 meters when the water activity is held constant [Romanovskii et al., 2005]). The insulating layer of 1164 sediment should act to slow the time scale for the response to ocean 1165 warming, to thousands of years [Archer, 2007]. Shakhova et al [2010b] 1166 1167 proposed that 50 Gton C as methane could erupt from the Arctic on a time scale of a few years. As has been acknowledged, the model 1168 provides poor constraint on the standing stock of bubbles or methane 1169 hydrate in the sediment column, and neglects many of the mechanisms 1170 that could come into play in transporting methane quickly to the 1171 atmosphere, such faults, channels, and blowouts of the sediment column. 1172 However, it is clear that the time scale for ocean warming to perturb 1173 1174 methane hydrates that are at least 300 meters below the sea floor will be much slower than a few years, leading to the expectation, consistent with 1175 1176 the global warming simulations of the model, that the methane cycle on the shelf should take thousands of years to respond to a climate change. 1177 The magnitude, and the time scale, of the model response ensures that 1178 nothing the model can do would generate such a large abrupt methane 1179 1180 eruption. An abrupt release would therefore require a large, contiguous gas pocket suddenly released by melting permafrost, like Arctic lakes that 1181 1182 drain away overnight. But 50 Gton of C represents a large fraction of all the traditional natural gas deposits on Earth (about 100 Gton C). The 1183 place to look for such a large unstable gas reservoir is in the field, not in 1184 this model, but until such a thing is found it remains conjecture. 1185
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1515	
1516	6. Figure Captions
1517 1518 1519 1520 1521 1522	Figure 1. Thermodynamics of hydrate and ice. Top) Colors are salinities, which range from fresh if there is no solid phase, to saltier as the freezing point depression of the solid phase follows the in situ temperature. Contours indicate the extent of thermal disequilbrium, $\Delta T_{eq} = T - T_{eq}$ . a) For the system of ice and fluid. b) Considering hydrate and fluid phases, excluding ice formation and assuming equilibrium with methane gas. c)

- 1523 Combined ice + hydrate + fluid system, where the salinity is controlled by the most stable solid phase. Solid contours are ΔT<sub>eq hydrate</sub>, dashed ΔT<sub>eq ice</sub>. 1524 1525 d and e) Colors are  $\Delta T_{eq}$ , where 0 (purple) indicates stability, and contours are the excess salinity relative to a solid phase, e.g. S<sub>max</sub> - S<sub>eq. hydrate</sub> in (d). 1526 1527 for hydrate, and e) ice. f) Phase diagram for the ice + hydrate + brine 1528 system. Hydrate is excluded from the ice phase space by the high salinity 1529 of the brine. Ice is ideally also excluded from part of the hydrate stability 1530 zone by a similar mechanism, but this would only happen in nature under conditions of unlimited methane availability. Thus it is easier to envision 1531 1532 coexistence of hydrate and ice within the hydrate stability zone, under 1533 conditions of limited methane availability, than it is to imagine hydrate in 1534 the permafrost zone, where ice has no impediment for formation. 1535 Figure 21. Domain of the model as applied to the Laptev Sea continental 1536 shelf and slope. This is the result of 62 million years of sediment 1537 accumulation on the crust, isostatic subsidence, pore fluid flow, and 1538 thermal diffusion, used as the initial condition for glacial / interglacial 1539 cycle and climate change simulations. Color indicates temperature. a) 1540 Full view. Black line shows the bottom of the crust, which grades 1541 smoothly from continental on the left into ocean crust through most of 1542 the domain on the right. b) Zoom in to see increased model resolution in 1543 the upper kilometer of the sediment column. 1544 Figure 32. Particulate Organic Carbon (POC) concentration. Highest 1545 values are found in the sediment depocenter just off the continental shelf 1546 break. 1547 Figure 43. Pore water salinity a) The fully marine case, in which the
- 1548 sediment column has always been submerged underneath a time-invariant
- sea level. b) Result of sediment column freshening by hydrological
- 1550 groundwater flow, driven by the pressure head resulting from a water
- 1551 table higher than sea level. A movie of the transition from marine to
- 1552 <u>freshened (the origin of b)</u>a to b can be seen at
- $1553 \qquad \text{http://geosci.uchicago.edu/} \sim \text{archer/spongebob\_arctic/} \\ \frac{\text{sal.65e6.nc.drop1}}{\text{sal.65e6.nc.drop1}}$
- 1554 **20**fig4.movie.gif
- 1555 Figure <u>5</u>4. Initial distribution of dissolved methane. a) Concentration in
- 1556 moles/m<sup>3</sup>. b-d)  $\Omega = CH_4 / CH_{4(sat)}$  deviation from equilibrium, b) of the
- 1557 Marine (salty) initial condition; c) of the pre-freshened initial condition
- 1558 (note depletion in near-surface near-shore sediments in the upper left);

- d) including permeable channels every five grid points, plus pre-
- 1560 freshening.
- 1561 Figure 65. Freshening the sediment column by hydrological groundwater
- 1562 flushing. Color indicates salinity. Solid black line represents sea level in
- 1563 the ocean (white space), and the equilibrium fresh-salty boundary given a
- 1564 snapshot of the pressure head (the Ghyben-Herzberg relation). Left side:
- results of dropping sea level 30 meters and holding it there. A freshwater
- 1566 lens forms and strives to reach Ghyben Herzberg equilibrium as the
- sediment column subsides, where atmospheric exposure decreases its
- 1568 buoyancy and stops sediment accumulation. After the sediment column
- subsides beneath the still-lowered sea level, the fresh water lens remains
- 1570 for millions of years. A movie can be seen at
- 1571 <a href="http://geosci.uchicago.edu/~archer/spongebob\_arctic/sal.zoom.65e6.nc">http://geosci.uchicago.edu/~archer/spongebob\_arctic/sal.zoom.65e6.nc</a>
- 1572 <u>.drop030/fig6a.movie.qif</u> . Right side: Result of dropping sea level 120
- meters and holding it there forever. Movie at
- 1574 <a href="http://geosci.uchicago.edu/~archer/spongebob\_arctic/fig6sal.zoom.65e">http://geosci.uchicago.edu/~archer/spongebob\_arctic/fig6sal.zoom.65e</a>
- 1575 <u>6.nc.drop12b0.movie.gif</u>
- 1576 Figure 76. Time scale of depleting the salinity of the continental shelf
- 1577 sediment column after an instantaneous sea level drop of 30 meters. The
- 1578 effect of lateral canyons is to provide a pathway for saline fluid to be
- 1579 replaced by fresh groundwater in sediments above sea level. If the lateral
- 1580 canyon spacing is 10 km, they can have a significant impact on the time
- 1581 constant for ground water flushing. A more conservative 100-km canyon
- 1582 is adopted for the rest of the simulations.
- 1583 Figure <u>87</u>. Dissolved methane impact by hydrological freshening of the
- 1584 sediment column as described in Figure 5.  $\Omega = CH_4 / CH_{4(sat)}$ . Movies can
- 1585 be seen at
- 1586 <a href="http://geosci.uchicago.edu/~archer/spongebob\_arctic/fig8ad\_ch4.65e6.">http://geosci.uchicago.edu/~archer/spongebob\_arctic/fig8ad\_ch4.65e6.</a>
- 1587 <u>nc.drop030.movie.gif</u> and
- 1588 <a href="http://geosci.uchicago.edu/~archer/spongebob\_arctic/fig8b/d\_ch4.65e">http://geosci.uchicago.edu/~archer/spongebob\_arctic/fig8b/d\_ch4.65e</a>
- 1589 <u>6.nc.drop120.movie.gif</u>
- 1590 Figure <u>98</u>. Time-dependent forcing for the glacial / interglacial
- 1591 simulations and the global warming scenarios. a) Sea level is imposed as
- 1592 a sawtooth 100-kyr cycle, with interglacial intervals shaded. The GW+S
- simulation tracks potential changes in sea level on long time scales due to
- 1594 fossil fuel CO<sub>2</sub> release, following a covariation from the geologic past of

1595 15 meters / °C. The GW and Control simulations hold sea level at interglacial levels. b) Ocean temperature forcings. 1596 1597 Figure 109. Colors indicate salinity in the unfrozen pore fluid of the 1598 sediment column. Thin solid black contours show the frozen fraction of 1599 the pore space. Heavy black stippled contour shows the stability 1600 boundary of methane hydrate as a function of temperature, pressure, and 1601 unfrozen pore fluid salinity. Left side: previously pre-freshened initial 1602 condition. Right side: Pure marine initial condition. c-d) Lowered sea level 1603 (from 70 kyr in Figure 8) but warm air temperatures prevent permafrost 1604 formation. e-f) Glacial conditions of lowered sea level (70 kyr) and 1605 atmospheric temperature of -17 °C driving permafrost formation. The 1606 pre-freshened and the marine initial conditions differ in the frozen fraction 1607 of sediment, but the salinity of the unfrozen fluid, a correlate of the 1608 activity of water, depends only the temperature. g-h) Rising sea level (at 1609 90 kyr in Figure 8) into an interglacial interval. A-Movies of the glacial 1610 cycles (GL) movie of with the pre-freshened simulation initial condition can 1611 be seen at 1612 http://geosci.uchicago.edu/~archer/spongebob\_arctic/fig10asal.zoom.6 1613 5e6.nc.ld2.gl.pfb.movie.gif, and the marine simulation initial condition at 1614 http://geosci.uchicago.edu/~archer/spongebob\_arctic/fig10bsal.zoom.6 5e6.nd.nc.ld2.al.pfb.movie.qif. 1615 Figure 110. Pore fluid pressure forcing and flow through the glacial 1616 cycles. Left) Colors indicate P<sub>excess</sub> + P<sub>head</sub>, solid contours are ice fraction, 1617 dashed contours are P<sub>head</sub>. Right) Colors indicate P<sub>excess</sub> + P<sub>head</sub>, note 1618 different color scale from Left. Initial refers to the prefreshened initial 1619 condition. "Low Sea Level" refers to simulation SL. "Glacial" and 1620 1621 "Interglacial" refer to simulation GL. Dashed contours indicate ice 1622 fraction, vectors fluid velocity. Movies of the prefreshened initial 1623 condition and glacial cycles (GL) can be seen at 1624 http://geosci.uchicago.edu/~archer/spongebob\_arctic/press\_uw.65e6.n 1625 c.ld2.gl.pf\_eg.gw.comp.movie.gif and 1626 http://geosci.uchicago.edu/~archer/spongebob\_arctic/pressure\_flow.65 1627 e6.nc.ld2.gl.pf\_eq.gw.comp.movie.gif. 1628 Figure 12. Sensitivities of the hydrate stability zone. Impact of the 1629 competition between ice and hydrate phases (a-d), and the geothermal 1630 temperature gradient (e-f). When ice is included as a potential solid 1631 phase, the pore waters are salty in the permafrost zone (a), restricting hydrate stability to at least 300 meters below sea level thoughout the 1632

1633 simulation (c). When ice is forbidden to form, hydrate can be stable 1634 nearly to the sediment surface during the height of the glaciation (b and 1635 d). The base of the stability zone is sensitive to the geothermal 1636 temperature gradient, while the shallowest reach of the stability zone 1637 does not respond to changing heat fluxes, because the temperatures are 1638 "anchored" at the ocean value at the top of the sediment column. 1639 Figure 134. Dissolved methane concentration relative to equilibrium ( $\Omega =$ CH<sub>4</sub> / CH<sub>4(sat)</sub>). Solid contours indicate ice fraction, dashed contours show 1640 the methane hydrate stability boundary. Movies for the left, center, and 1641 1642 right columns, respectively can be seen at 1643 http://geosci.uchicago.edu/~archer/spongebob\_arctic/fig13ad\_ch4.65e 1644 6.nc.al.pfb.movie.aif. http://geosci.uchicago.edu/~archer/spongebob\_arctic/fig13bd\_ch4.65e 1645 6.nc.ld2.gl.pfb.movie.gif, and 1646 http://geosci.uchicago.edu/~archer/spongebob\_arctic/fig13cd\_ch4.65e 1647 6.nd.nc.ld2.gl.pfb.movie.gif for the left, center, and right columns, 1648 1649 respectively. Figure 142. Carbon cycle through glacial cycles from a prefreshened 1650 1651 initial condition. Solid contours: Ice Fraction. Dashed contours: Methane hydrate stability zone. Left) Particulate organic carbon (POC) 1652 1653 concentration. Movie at 1654 http://geosci.uchicago.edu/~archer/spongebob\_arctic/fig14apoc.65e6.n c.ld2.gl.pfb.movie.gif. Center) Biological methane production rate. Movie 1655 1656 at 1657 http://geosci.uchicago.edu/~archer/spongebob\_arctic/fig14bch4\_src.65 1658 e6.nc.ld2.gl.pfb.movie.gif Right) Methane hydrate concentration. Movie 1659 1660 http://geosci.uchicago.edu/~archer/spongebob\_arctic/fig14chydrate\_bu bbles\_zoom.65e6.nc.ld2.gl.pfb.movie.gif-. Movies of methane hydrate 1661 1662 stability and concentration are given for the sensitivity studies, in the 1663 supplemental material and at 1664 http://geosci.uchicago.edu/~archer/spongebob/. Figure 153. Glacial cycle of methane hydrate inventory on the 1665 continental shelf for the simulation including permafrost formation. a) 1666 1667 Effects of salt and ice. b) Sensitivity to methaneogenesis rates. c) Sensitivity to the column temperature gradient, d) Glacial cycles of shelf 1668

bubble inventories, effects of salt and ice.

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Figure 164. Spatial distribution and sea level impact of methane fluxes to 1670 1671 the atmosphere. a-d) Solid line shows the elevation of the sediment 1672 surface relative to the sea level at the time. Grey lines (scale to right) show the efficiency of bubble transport through the water column, 1673 1674 assuming a flux attenuation length scale of 30 meters. e-k) Dashed line: 1675 Methane bubble flux across the sediment surface. Solid line: Methane 1676 bubble flux to the atmosphere (dashed line multiplied by transport 1677 efficiency). Most of the methane flux in the model occurs near the shelf 1678 break, and submergence in the ocean has a strong impact on the flux to the atmosphere. A related movie can be seen at 1679 1680 http://geosci.uchicago.edu/~archer/spongebob\_arctic/fig16bubb\_atm.6 1681 5e6.nc.ld2.ql.pfb.qw.comp.movie.qif. 1682 Figure 175. Glacial / interglacial cycle of methane fluxes on the continental margin of the model. a) Sea level driver for reference at top, 1683 1684 grey regions indicate interglacial intervals, pink the Anthropocene. b) Biological methane production on the shelf. Green is case SL (sea level 1685 changes only), black is SL+LD (adding land deposition of carbon-rich 1686 soils), red is GL (adding permafrost formation). ac-fe) Cumulative 1687 1688 methane fluxes. Red lines show production rate. Brown regions show 1689 lateral transport of dissolved methane. Grey shows oxidation by SO<sub>4</sub>2- in 1690 the sediment column. Blue shows bubble flux to the water column. During 1691 interglacial times (e.g. far left) there is a small onshore transport of 1692 methane, which is represented by a negative starting point for the 1693 oxidation (grey) region. In equilibrium, the colored areas should fill in the region under the red curve. 1694 1695 Figure 186. Methane fluxes to the atmosphere. Sea level at the top. 1696 interglacial intervals in vertical grey bars, the Anthropocene in pink, a) 1697 From a pre-freshened initial condition, with and without permafrost formation. b) From a pure marine initial condition. c and d) Sensitivity to 1698 terrestrial organic carbon deposition during low sea-level stands, and to 1699 thermogenic methane flux. e) Sensitivity to the impact of ice fraction on 1700 bubble mobility. 1701 Impact of sediment column pre-freshening and sea level changes on 1702 model methane fluxes to the atmosphere. Black lines show results 1703 without permafrost formation, red lines are with permafrost, a) From a 1704 pre-freshened initial condition, lowering the mean salinity of the sediment 1705

column. b) From a pure marine initial condition.

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1707 1708 1709	Figure 1 <u>9</u> 7. Impact of anthropogenic warming on the methane cycle in the model. <u>a) Base cases, a warming scenario (GW), without and with a geological time-scale sea level rise scenario (+SLR), and extended integral side and the little relationship.</u>
1710	interglacial control (Ctl). Warming plus increasing sea level decreases the
1711	methane flux overall, due to bubble dissolution in a deeper water column.
1712	b) Altered model physics impacts. c and d) Altered methanogenesis
1713	rates. e) Sensitivity to the ice fraction at which bubble mobility is
1714	assumed stopped.
1715 1716 1717	The effect of the warming by itself is a slight increase in the methane flux to the atmosphere. Warming + increasing sea level attenuates the methane flux, due to bubble dissolution in a deeper water column.
1718	Tables

1719 <u>Table 1. Summary of model runs.</u>

SL	Sea level changes with constant air and water temperatures
GL	SL + glacial cycles in air and water temperature
GW	A long-term global warming scenario, a peak and long tail temperature perturbation consistent with CO <sub>2</sub> release and cessation of the glacial sawtooth forcing.
+SLR	Adds geologic-timescale sea level rise due to anthropogenic climate change, based on correlation between temperature and sea level in the geologic past (10 meters / °C).
<u>Ctl</u>	An extended interglacial with no CO <sub>2</sub> release forcing.
+ LD	Land deposition of carbon-rich Yedoma. Base case is 10 m / 100 kyr, with sensitivity runs using 30 and 100 m / 100 kyr accumulation of 30% POC material. Movies in the supplemental material are identified by the tags Land30 and Land100.

+ TG	Thermogenic methane production rate sensitivity runs, scaling the rate from the spinup result by factors of 10 and 100.  Movies in the supplemental material are identified by the tags TGenX10 and TGenX100.
+ Geotherm	Sensitivity of ice and hydrate cycles on the geothermal temperature gradient.  Temperatures from the Base simulation were adjusted when calculating the stability of ice and hydrate, to simulate the impact of geothermal heat fluxes on hydrate stability.  Note that other aspects of the sediment column, including the solubility of methane, retained the original temperatures. Heat fluxes simulated include 25 mW/m2, 37.5, 50 (Base), 62.5, and 75. Movies of the non-base runs are identified by tags HF050, HF075, HF125, and HF150.
Ice and Bubble Transport	When the ice fraction exceeds a threshold value methane gas flow is disabled. Base case is 50%, variants 10%, 30%, 70%, and 90%, identified with tags Ice10, Ice30, Ice70, and Ice90.
No Ice	The ice phase is disallowed in the thermodynamic calculation. Movies in the supplemental material include salinity. The files are tagged as Nolce
No Salt from Ice	Ice is allowed to form, but it does not affect the salinity as it determines methane hydrate stability. Movie files are tagged as NoSalFromIce.
Permeable Channels	Increasing vertical permeability by a factor of 10 every 5 <sup>th</sup> grid cell, to generate

	heterogeneity in the flow. Tagged as PermChan
No Horizontal Flow	Horizontal flow is disabled. Tagged as NoHFlow.

1720	Movies comparing altered scenario runs with the Base scenario are given
1721	in the supplemental material, and at
1722	http://geosci.uchicago.edu/~archer/spongebob/. Movies named
1723	hydrate* and bubbles* show methane hydrate and bubble inventories and
1724	stability zone changes. Files entitled salinity* show salinities, and
1725	bubb_atm* show bubble fluxes through and out of the sediment column,
1726	into the ocean, and into the atmosphere, through time.