

1 **Abstract**

2 In recent decades, the central North Sea has been experiencing a general trend of
3 decreasing dissolved oxygen (O₂) levels during summer. To understand potential
4 causes driving lower O₂, we investigated a three-day period of summertime
5 turbulence and O₂ dynamics in the thermocline and bottom boundary layer (BBL).
6 The study focuses on coupling biogeochemical with physical transport processes to
7 identify key drivers of the O₂ and organic carbon turnover within the BBL.
8 Combining our flux observations with an analytical process-oriented approach, we
9 resolve drivers that ultimately contribute to determining the BBL O₂ levels. We report
10 substantial turbulent O₂ fluxes from the thermocline into the otherwise isolated
11 bottom water attributed to the presence of a baroclinic near-inertial wave. This
12 transient contribution to the local bottom water O₂ and carbon budgets has been
13 largely overlooked and is shown to play a role in promoting high carbon turnover in
14 the bottom water while simultaneously maintaining high O₂ concentrations. This
15 process could become suppressed with warming climate and stronger stratification,
16 conditions which may promote migrating algal species that could potentially shift the
17 O₂ production zone higher up within the thermocline.

18

19 **1 Introduction**

20 **1.1 Hypoxia in shelf seas and coastal regions**

21 The distribution of dissolved oxygen (O₂) in marine systems results from a
22 complex interaction between biological processes (photosynthesis and respiration)
23 and physical processes (O₂ flux pathways) occurring within the water column and at
24 the seafloor. O₂ is regarded as an important indicator of ecosystem functioning for
25 aquatic organisms (Best et al., 2007) and for benthic activity (e.g., Glud, 2008).
26 Changes in O₂ distribution, concentrations and supply can therefore have severe
27 impacts on shelf ecosystems. O₂ concentrations below 62.5 μmol L⁻¹, which is
28 generally regarded as the threshold of hypoxia (Vaquer-Sunyer and Duarte, 2008), are
29 shown to significantly stress aquatic communities and increase the mortality among
30 fish communities (Diaz, 2001). These ecological and economic impacts of O₂
31 depletion lead to increasing concern regarding hypoxia occurrence and hypoxic
32 events. As reviewed by Diaz and Rosenberg (2008), hypoxia in coastal environments
33 is spreading and so are the reports of unprecedented occurrence of hypoxia in several

34 shelf seas and coastal regions (Grantham et al., 2004; Chan et al., 2008; Crawford and
35 Pena, 2013).

36

37 **1.2 Hydrodynamics and oxygen depletion in the North Sea**

38 The North Sea is situated on the North–West European continental shelf,
39 between the British Islands and continental Europe (Fig. 1). Its semi-enclosed basin
40 covers an area of 575'300 km², with an average depth of 74 m and a general decrease
41 in depth from North to South (Otto et al., 1990). The central region is characterized
42 by the presence of the Dogger Bank, a shallow sandbank that acts as a hydrological
43 divide. The northern and central North Sea hydrology is mainly dominated by inflow
44 from the North Atlantic Ocean at the northern open boundary, while the southern part
45 relies on inflow from the English Channel (Thomas et al., 2005). Northern and central
46 North Sea areas are characterized by seasonal water column stratification (April to
47 September - October; Meyer et al., 2011). With only weak, wind-driven residual
48 currents (Otto et al., 1990), this stratification leads to isolation of central North Sea
49 bottom water and subsequent O₂ depletion.

50 In the central North Sea, the occurrence of low O₂ levels in bottom waters has
51 been reported (e.g., North Sea Task Force, 1993; Greenwood et al., 2010). Additional
52 monitoring studies in the central North Sea in 2007 and 2008 have shown that O₂
53 concentration in bottom waters at the Oyster Grounds and North Dogger can drop to
54 163 – 169 μmol L⁻¹ (60 – 63% saturation) and ~200 μmol L⁻¹ (71% saturation),
55 respectively (Fig. 1; Greenwood et al., 2010). Comparable field observations were
56 also reported in the summer of 2010 (Queste et al., 2013). The authors also reviewed
57 the available historical O₂ data in the North Sea (1900 – 2010), revealing a clear
58 increase in O₂ depletion after 1990.

59 While the reported O₂ levels were still above the hypoxic threshold, growing
60 concerns of hypoxia developing in the North Sea have highlighted the need for more
61 detailed studies on O₂ dynamics and its driving forces (Kemp et al., 2009). Since
62 1984, surface water temperatures in the North Sea have increased by 1 – 2°C, greater
63 than the global mean (OSPAR, 2009, 2010; Meyer et al., 2011). On seasonal time
64 scales, climate projections indicate longer durations of the stratification period and
65 stronger thermocline stability (Lowe et al., 2009; Meire et al., 2013), with some
66 projections suggesting earlier onset of stratification (e.g., Lowe et al., 2009). Due to

67 the semi-enclosed nature of the North Sea, earlier onset and longer stratification
68 increases the length of time that deep waters are isolated, potentially allowing lower
69 O₂ concentrations to develop (Greenwood et al., 2010).

70

71 **1.3. Physical drivers of oxygen dynamics**

72 The distribution of O₂ and other dissolved constituents within aquatic systems
73 are largely driven by physical transport processes. These include wind driven air –
74 water gas exchange at the sea surface (Wanninkhof, 1992), molecular diffusion at the
75 sediment – water interface (Jørgensen and Revsbech, 1985), horizontal advection
76 (e.g., Radach and Lenhart, 1995) and turbulent transport in the water column, where
77 the latter transport was reported to significantly contribute to constituent balances (see
78 Rippeth, 2005; Fischer et al., 2013; Kreling et al., 2014; Brandt et al., 2015). In shelf
79 seas, the seasonal occurrence of steep thermoclines acts as an important physical
80 barrier separating the surface layer from nutrient-rich deeper waters (Sharples et al.,
81 2001). As measurements of shear and stratification have shown, the central North Sea
82 thermocline is in a state of marginal stability (van Haren et al., 1999). Hence
83 additional sources of shear could trigger shear instability leading to local production
84 of turbulence within the thermocline. This enhanced local turbulence would
85 subsequently enhance the vertical exchange of constituents such as O₂, organic carbon
86 and nutrients. Resolving processes that drive diapycnal (i.e., vertical) fluxes across the
87 thermocline throughout the stratification period is key to understanding the
88 biogeochemical functioning of shelf seas (e.g., Sharples et al., 2001).

89

90 **1.4 Present study**

91 The goal of this study is to provide understanding of key turbulent processes
92 driving O₂ fluxes across the thermocline during the summertime stratification period
93 in the central North Sea using data from a 3-day process study. We investigate and
94 describe O₂ dynamics and fluxes to the bottom waters and discuss their potentially
95 influence on the seasonal O₂ balance. Using the resolved O₂ fluxes, we perform a
96 simple 1-D mass balance model to quantify O₂ sources and sinks. Finally, processes
97 that could further promote hypoxia in the central North Sea in a warming climate are
98 discussed.

99

100 **2 Methods**

101 **2.1 Study site**

102 We performed O₂ and turbulence measurements in the Norwegian sector of
103 the central North Sea, N. 1/9, at the Tommeliten site (56°29'30" N, 2°59'00" E; Fig. 1)
104 from 8 – 11 August 2009 aboard the R/V *Celtic Explorer* (cruise CE0913). The site,
105 located ~100 km northeast from the northern Dogger Bank, and its surroundings are
106 characterized by shallow waters (~70 m) relatively far from coastal areas (on average
107 ~300 km). The site is known for the presence of buried salt diapirs, methane (CH₄)
108 seeps and bacterial mats (Hovland and Judd, 1988). Bathymetric surveys from
109 Schneider von Deimling et al. (2010) revealed a rather flat sandy seabed with almost
110 no features, with the exception of cm-sized ripples (McGinnis et al., 2014).

111 The currents of the central North Sea are predominantly driven by the semi-
112 diurnal lunar tide (M₂) (Otto et al., 1990). Seasonal stratification begins in April
113 around Julian day 100 and lasts until the end of September or early October, Julian
114 days 270- 290 (e.g. Meyer et al., 2011). The thermocline has been identified as an
115 important zone for the establishment of primary production and the O₂ maximum
116 layer (see Pingree et al., 1978). In fact, the North Sea deep chlorophyll maximum
117 (DCM) is estimated to account for 58% of the water column primary production and
118 37% of the annual new production for the summer stratified North Sea (Weston et al.,
119 2005). The development of the associated O₂ maximum due to this production is thus
120 important and so far not considered in the overall O₂ balance of the central North Sea.

121

122 **2.2 Instrumental setup**

123 High resolution (mm scale) turbulent shear and temperature profiles were
124 obtained with a MSS90-L microstructure turbulence profiler (Sea and Sun
125 Technology, Trappenkamp, Germany). The MSS90-L is a free-falling, loosely-
126 tethered profiler which samples at 1024 Hz with 16 channels and is designed for an
127 optimal sink rate of 0.5 – 0.6 m s⁻¹. The probe was equipped with two air-foil shear
128 probes, an accelerometer (to correct for probe pitch, roll, and vibration), a fast
129 temperature sensor (FP07, 7–12 ms response time), standard CTD sensors
130 (temperature, pressure, conductivity), and a fast (0.2 s response time) galvanic O₂
131 sensor (AMT, Analysenmesstechnik GmbH, Rostock, Germany). Absolute O₂

132 concentrations were calibrated against shipboard CTD O₂ profiles and Winkler
133 titrations using discrete water samples (see below).

134 Water column hydrodynamics were characterized with the compact benthic
135 Paleooceanography (POZ) lander, which was deployed using a video guided launcher
136 (Pfannkuche and Linke, 2003). The POZ lander recorded 3-dimensional current
137 velocity profiles and acoustic backscatter information throughout the water column
138 using a 300 kHz acoustic Doppler current profiler (ADCP; Workhorse Sentinel,
139 Teledyne RD Instruments, Poway, United States), which sampled every 15 s with a
140 bin size of 0.5 m starting from 2.75 m from the bottom. A conductivity-temperature-
141 depth (CTD) logger (XR-420 CT logger, RBR, Kanata, Canada) recorded
142 temperature, conductivity and pressure (Digiquartz, Paroscientific, Redmond, United
143 States) every 2 s near the seafloor (~0.3 m distance). The POZ lander was also
144 equipped with a Winkler-calibrated O₂ optode sensor (Aanderaa Data Instruments
145 AS, Bergen, Norway), which recorded BBL O₂ concentration at 1 min intervals.

146 Water column profiles were obtained using a SBE9plus CTD-rosette system
147 (Seabird, Washington, United States). The CTD sampled at 24 Hz and was equipped
148 with standard temperature, conductivity, pressure, O₂ and light transmission sensors.
149 The rosette system mounted 12 Niskin bottles (10 L each) for discrete water
150 sampling. Each water sample was subsampled with three Winkler bottles of known
151 volume (~62 mL on average) upon recovery, and the samples were immediately fixed
152 on deck and titrated manually within 24 h after the sampling (see Winkler 1888).
153 CTD O₂ concentrations deviated from Winkler values by <5%.

154

155 **2.3 Hydrodynamic data evaluation**

156 The main tidal directions (major and minor axis of the tidal ellipsoid) were
157 determined by performing a variance analysis on the ADCP velocity time series. The
158 u and v velocities were rotated over a stepwise increasing rotation angle (r) as
159 $u_{rot} = u \cdot \cos(-r) - v \cdot \sin(-r)$ and $v_{rot} = u \cdot \sin(-r) - v \cdot \cos(-r)$, and the
160 variance computed at each step. The angle with the largest variance is the main tidal
161 direction. Barotropic and baroclinic flow contributions of tides were separated by
162 least-square fitting the detrended velocity time series to harmonics $u = A \cdot \cos(\omega \cdot$
163 $t + \varphi)$ with A , ω , φ being the amplitude, frequency, and the phase lag, respectively.
164 In the analysis below, the barotropic semi-diurnal principle lunar tide (M₂) and

165 diurnal declination tide (K_1) contributions had frequencies of 1.93227 cycles per day
 166 (cpd) and 1.00274 cpd, respectively, and were subtracted from the time series to
 167 analyze residual flow. For barotropic contributions, the fit was applied to the depth
 168 average of the time series, while baroclinic contributions were obtained by fitting the
 169 harmonics to the velocity time series from each 0.5 m ADCP bin. The occurrence of
 170 enhanced shear in the stratified water column was investigated by calculating the
 171 vertical shear of horizontal velocity, S , from the vertical gradients between adjacent
 172 bins of east and north velocity (0.5 m resolution) as $S = \sqrt{(du/dz)^2 + (dv/dz)^2}$.
 173 Frequency spectra of the time series of horizontal velocity and vertical shear of
 174 horizontal velocity were used to identify the tidal and non-tidal flow components. The
 175 spectra were calculated using fast-Fourier transforms combined with a $\frac{1}{2}$ -cosine
 176 taper (Hanning window), which was applied to the first and last 10% of the time
 177 series data.

178 Turbulent kinetic energy dissipation rate (ε) was quantified from airfoil shear
 179 readings by integrating shear wavenumber spectra assuming isotropic turbulence
 180 (Batchelor, 1953):

$$\varepsilon = 7.5\mu \int_{k_{min}}^{k_{max}} E_{du'/dz}(k) dk \quad (1)$$

181 where μ is the dynamic viscosity of seawater. Shear spectra $E_{du'/dz}(k)$ were
 182 calculated from one-second ensembles (1024 values) and integrated between a lower
 183 $k_{min} = 3$ cycles per minute (cpm) and an upper wavenumber k_{min} that varied between
 184 14 cpm and 30 cpm depending on the Kolmogorov wavenumber. Here, a Bartlett
 185 window was applied to the whole ensemble prior to spectral decomposition. Loss of
 186 variance due to limited wavenumber band was taken into account by fitting the
 187 observed shear spectra to the universal Nasmyth spectrum. Similarly, corrections for
 188 the loss of variance due to finite sensor tip of the airfoil probes were applied (see
 189 Schafstall et al., 2010). The detection limit, or noise level, of the used profiler for ε
 190 was inferred to be 1×10^{-9} W kg⁻¹ (Schafstall et al., 2010); the upper detection limit is
 191 a function of the shear sensor geometry (up to 10^{-4} W kg⁻¹; Prandke and Stips, 1998).

192 Estimates of turbulent eddy diffusivities of mass (K_ρ) were obtained from
 193 measurements of ε as

$$K_\rho = \gamma\varepsilon/N^2 \quad (2)$$

194 where γ is the mixing efficiency and N^2 the water column stability. This method,
 195 proposed by Osborn (1980), approximates K_ρ under the assumption of a local
 196 equilibrium of production and dissipation of turbulent kinetic energy. Values for N^2
 197 were calculated from temperature, salinity and pressure data using the adiabatic
 198 method (Fofonoff, 1985) as $N^2 = -g(\rho^{-1}\partial\rho/\partial z - g/c^2)$, where ρ , g , and c are the
 199 density, the earth's gravitational acceleration and speed of sound. Mixing efficiency
 200 values in stratified waters range from 0.1 to 0.2 (Ivey and Imberger, 1991) and
 201 decreases in weakly stratified waters such as within the BBL (Lorke et al., 2008). To
 202 account for this decrease, we used the γ and K_ρ parameterization of Shih et al. (2005).
 203 Based on the turbulence activity parameter $\varepsilon/\nu N^2$, with the kinematic viscosity, ν ,
 204 the authors found that in energetic regimes, i.e., $\varepsilon/\nu N^2 > 100$, eddy diffusivities are
 205 better estimated as $K_\rho = 2\nu(\varepsilon/\nu N^2)^{1/2}$. As horizontal density gradients at the study
 206 site were deemed to be small compared to vertical gradients (see Discussion), we
 207 equated diapycnal eddy diffusivities with vertical diffusivities (i.e., $K_\rho = K_z$).

208 To obtain representative mean turbulent eddy diffusivities, the data were
 209 evaluated in ensembles of three to four consecutive profiles and averaged in depth and
 210 time to reduce uncertainties due to the patchiness of turbulence, temporal fluctuation
 211 of N^2 , and temporal γ variations (see Smyth et al., 2001). As proposed by Ferrari and
 212 Polzin (2005), the level of uncertainty of the averaged K_z can be quantified as:

$$\Delta K_z = K_z \left[\left(\frac{\Delta\gamma}{\gamma} \right)^2 + \left(\frac{\Delta\varepsilon}{\varepsilon} \right)^2 + \left(\frac{\Delta N^2}{N^2} \right)^2 \right]^{1/2} \quad (3)$$

213 with Δ being the absolute uncertainty of the various average terms. Here, the
 214 uncertainties are evaluated in the region of strong vertical O_2 gradients and in 2 m
 215 depth bins. The absolute uncertainty for $\Delta\gamma$ was assumed to be 0.04 (see St. Laurent
 216 and Schmitt, 1999). The absolute uncertainty on N^2 (ΔN^2) was determined by the
 217 standard error over the 2 m average, computed as the standard deviation divided by
 218 the square root of the number of estimates. Finally, the statistical uncertainty of ε for
 219 each bin was calculated using a bootstrap method (10^4 resamples) (Efron, 1979).

220 The vertical O_2 fluxes F_θ were then obtained from K_z and the O_2 concentration
 221 gradients $\partial[O_2]/\partial z$ as

$$F_\theta = K_z \frac{\partial[O_2]}{\partial z} \quad (4)$$

222 Accordingly, the uncertainty of averaged turbulent O_2 fluxes were given by:

$$\Delta F_{\theta} = F_{\theta} \left[\left(\frac{\Delta K_z}{K_z} \right)^2 + \left(\frac{\Delta \partial_z [O_2]}{\partial_z [O_2]} \right)^2 \right]^{1/2} \quad (5)$$

223 where $\Delta \partial_z [O_2]$ denotes the standard error of mean vertical gradients of O_2
 224 concentrations. It should be noted that the analysis did not include biases or
 225 uncertainties due to measurement errors.

226

227 **3 Results**

228 During the three-day observational period (8 – 11 August 2009), we collected
 229 39 high-resolution MSS profiles in consecutive sets of three to five profiles at 5 – 10
 230 min intervals. Most of the profiles were in the evening (profiles 1 – 8, 26 – 28, 36 –
 231 39) or at night (9 – 15, 29 – 35) with the remaining profiles acquired in the morning
 232 (6 to 9 AM). One shipboard CTD profile was performed prior to the actual MSS
 233 profiles to provide hydrographic information, the water turbidity and O_2
 234 concentrations, and discrete water samples for subsequent onboard Winkler titrations.
 235 Hydroacoustic water column current measurements were carried out continuously
 236 throughout the observational period. The following results are structured to first
 237 present a characterization of the site’s physical settings and turbulence drivers,
 238 followed by the O_2 fluxes and O_2 BBL budget.

239

240 **3.1 Water column structure**

241 The ~70 m deep water column was characterized by a stable, well-defined
 242 four-layer temperature structure (Fig. 2a). The well-mixed surface boundary layer
 243 (SBL) and bottom boundary layer (BBL), 15 m and 30 m thick, respectively, were
 244 separated by a weakly-stratified transition layer (15 – 25 m depth) and a strongly
 245 stratified interior layer (25 – 40 m depth). The stratified interior layer was
 246 characterized by two very steep thermoclines situated in the upper (27 – 30 m depth)
 247 and lower (36 – 39 m depth) region of the layer, with vertical temperature gradients of
 248 up to 4°C m^{-1} . The average salinity was 35.08 with little variation throughout the
 249 water column (35.04 – 35.1). The light transmission profile from the ship CTD ranged
 250 from 89% to 96% (Fig. 2b). The most turbid layer (89%) was observed at the lower
 251 boundary of the interior layer (at 40 m depth) suggesting the presence of the deep
 252 chlorophyll maximum, phytoplankton, zooplankton and suspended particles.

253 The O₂ profiles were generally characterized by near saturation in the SBL
254 and transition layers (238 – 243 μmol kg⁻¹) and undersaturated (~80%) in the BBL
255 (~243 μmol kg⁻¹) (Fig. 2c,d). The stratified interior was oversaturated by up to 115%,
256 with a well-established O₂ maximum at ~39 m depth with concentrations up to ~315
257 μmol kg⁻¹. Below that maximum, at the thermocline-BBL interface, we observed a 2 –
258 3 m thick steep oxycline, with an O₂ gradient of 34 μmol kg⁻¹ m⁻¹ and exhibiting very
259 limited day/night, depth and thickness variation. We resolve the O₂ flux into the BBL
260 associated with this oxycline.

261

262 **3.2 Hydrodynamics**

263 The hydrostatic pressure dataset (POZ lander) revealed that the tidal water
264 level ranged from 0.6 to 0.9 m (Fig. 3a). Variance analysis on the ADCP velocity data
265 identified the major and minor axis of the tidal ellipsoid components to occur at 45°
266 and 135° from true north, respectively. Along these axes, the current amplitudes were
267 0.21 m s⁻¹ and 0.04 m s⁻¹, indicating a narrow tidal current ellipsoid, as reported by
268 Otto et al. (1990). The site was characterized by a negative tide polarity (anti-
269 cyclonic) for the semi-diurnal tides. A dominance of the barotropic M₂ current
270 amplitude at all depths was clearly visible in the velocity time series (Fig. 3b, c) and
271 the harmonic analyses. East (zonal) and north (meridional) barotropic M₂-current
272 amplitudes were 0.12 m s⁻¹ and 0.17 m s⁻¹, respectively, while K₁-current amplitudes
273 were only 0.005 m s⁻¹ and 0.03 m s⁻¹.

274 Although the limited length of the ADCP velocity time series did not allow for
275 full separation of the M₂ and *f* frequencies, the spectral density functions indicated
276 maximum energy at frequencies of about the semi-diurnal tide. This maximum varied
277 little with depth, indicating barotropic M₂ motions. Superimposed on those barotropic
278 currents, we observed the presence of baroclinic velocity contributions (Fig. 3b, c).
279 Additionally, near-inertial motions were also detected.

280 The occurrence of near-inertial motions was most pronounced in the
281 thermocline (32 – 39 m; Fig. 3e). Lower, but still elevated, energy densities at the
282 near-inertial band were found in the SBL and BBL. Moreover, the near-inertial
283 currents exhibited a distinct 180° phase shift between the SBL and the thermocline as
284 well as between the thermocline and the BBL, suggesting a second vertical mode
285 nature of these fluctuations. Average amplitudes of the near-inertial fluctuations in the

286 thermocline obtained from least-square fitting were 0.11 m s^{-1} . In the BBL and SBL,
287 average amplitudes were reduced to 0.06 m s^{-1} and 0.04 m s^{-1} , respectively,
288 suggesting that f oscillations might account for enhanced shear in the thermocline.

289 Enhanced vertical shear of horizontal velocity was found at the interior –
290 transition layer and at the interior – BBL interfacial regions (Fig. 3d). As indicated by
291 the spectral density function of the shear time series from the interior interfacial layers
292 (SI Fig. 1), the shear exhibited near-inertial frequencies (1.6722 cpd), and resulted
293 from the baroclinic near-inertial wave. The high vertical resolution (0.5 m) of our
294 velocity data allowed the resolution of the interfacial shear layers, which were
295 typically 2 to 3 m thick with elevated values of up to 0.05 s^{-1} . Comparisons with CTD
296 data showed that they are collocated with the two enhanced temperature gradients
297 layers in the thermocline ($27 - 30 \text{ m}$ and $36 - 39 \text{ m}$ depth; Fig. 2a).

298 The dissipation rates (ε) of turbulent kinetic energy (TKE) determined from
299 microstructure shear probes were particularly low in the center of the stratified
300 interior ($2 - 5 \times 10^{-9} \text{ W kg}^{-1}$) but still above the MSS detection limit. However, ε
301 increased to $5 \times 10^{-9} \text{ W kg}^{-1}$ and $2 \times 10^{-8} \text{ W kg}^{-1}$ at the upper and lower interior layer
302 limits, respectively (Fig. 4a). These coincided with the depth range of the interfacial
303 shear layers (Fig. 3d) at the strong temperature gradients (Fig. 2a) and resulting water
304 column stability maxima ($\sim 1 \times 10^{-3} \text{ s}^{-2}$).

305 Bin-averaged values of K_z varied by a factor of 5, ranging from $6 \times 10^{-7} \text{ m}^2 \text{ s}^{-1}$
306 in the central interior to $3 \times 10^{-6} \text{ m}^2 \text{ s}^{-1}$ in the lower region of the transition layer (Fig.
307 4b). In the upper interface (thermocline – transition layer), where ε was elevated with
308 respect to the central interior but reduced compared to the lower interfacial layer,
309 stronger stratification (i.e., larger N^2 values up to 10^{-3} s^{-2}) reduced the eddy
310 diffusivities. At the interior-BBL, higher K_z values ($\sim 2 \times 10^{-5} \text{ m}^2 \text{ s}^{-1}$) resulted from
311 increased turbulence and weaker stratification. This enhanced turbulent transport was
312 located where the vertical O_2 gradient was the strongest (Fig. 2d).

313

314 **3.3 Oxygen fluxes and budget**

315 With the fast responding AMT galvanic O_2 sensor and rapid sampling rate, we
316 were able to resolve the O_2 gradient with high precision. Figure 4c shows the 2 m bin
317 average O_2 fluxes for the interior together with the averages from each ensemble.

318 Small O₂ fluxes (~1 mmol m⁻² d⁻¹) were estimated for the center and upper region of
 319 the interior; suggesting that relatively little O₂ is transported upward from the O₂
 320 maximum to the upper interior. In contrast, a substantial O₂ flux ranging from 9 – 134
 321 mmol m⁻² d⁻¹ (average of 54 mmol m⁻² d⁻¹) was identified from the lower thermocline
 322 towards the BBL. The confidence interval associated with the uncertainties of the O₂
 323 flux estimates was 18 – 74 mmol m⁻² d⁻¹. Although the O₂ fluxes to the BBL water
 324 from the thermocline were variable in magnitude (Fig. 4c) and the measurements
 325 limited to the observational period (Fig. 3), their magnitude nevertheless suggests an
 326 important, yet overlooked, O₂ pathway.

327 We performed a simple 1-D BBL mass balance to investigate the relevance to
 328 the local O₂ balance during our observational period. Here, we defined the apparent
 329 (measured) O₂ loss rate in the BBL $\partial[O_2]/\partial t$ as the consequence of O₂ replenishment
 330 from F_θ and the O₂ utilization via sediment O₂ uptake rate (SUR) and water column
 331 organic matter respiration (R) expressed as

$$\frac{\partial[O_2]V}{\partial t A} = |F_\theta| - |SUR| - |R| \quad \{mmol\ m^{-2}\ d^{-1}\} \quad (6)$$

332 The mass balance was constrained to the (assumed) well-mixed 35 m deep BBL
 333 section of area, $A = 1\ m^2$ with a volume $V = 35\ m^3$. We further assumed negligible
 334 horizontal O₂ gradients (as observed from the CTD casts), and thus a net zero
 335 horizontal O₂ advective transport.

336 The average SUR for the same time period and location, obtained from parallel
 337 eddy correlation measurements, was ~-10 mmol m⁻² d⁻¹ (McGinnis et al., 2014). The
 338 SUR was consistent with the average SUR at Oyster Grounds reported by Neubacher
 339 et al. (2011), -9.8 mmol m⁻² d⁻¹, and with modeled SUR s at the same site (average -8.6
 340 mmol m⁻² d⁻¹; Meire et al., 2013). The apparent BBL O₂ loss of -0.42 μmol kg⁻¹ d⁻¹
 341 was determined from the POZ lander O₂ optode time series (Fig. 5a) over 52 hours,
 342 ($R^2=0.60$). Though limited to our short observational period, the vertically integrated
 343 apparent BBL O₂ loss was about -15 mmol m⁻² d⁻¹ and thus within 2% of the nearby
 344 North Dogger average presented by Greenwood et al. (2010). Based on Eq. (6) and
 345 using the observed BBL O₂ loss rate, F_θ and SUR , the water column respiration, R
 346 was calculated to be ~-60 mmol m⁻² d⁻¹. This implies that without the O₂
 347 replenishment, the apparent BBL O₂ loss would be ~-2 μmol kg⁻¹ d⁻¹ and thus four
 348 times higher than observed. Our results indicated that the total respiration in the

349 bottom water was therefore $\sim 70 \text{ mmol m}^{-2} \text{ d}^{-1}$ ($SUR + R$), with about 14% of the
350 organic carbon mineralization occurring in the sediment and 86% in the BBL.

351

352 **4 Discussion**

353 During our three-day observational period, we found that the baroclinic near-
354 inertial wave in the interior was the main contributor to the detected enhanced shear
355 (Fig. 3d) and the observed elevated vertical O_2 flux to the BBL (Fig. 6). As near-
356 inertial waves decay after a few weeks, it should be noted that we observed a rather
357 special situation, and that vertical O_2 fluxes will not likely be as highly elevated
358 during periods when near-inertial waves are not present.

359 Within this context, we: 1) discuss the turbulent mechanisms leading to these
360 thermocline O_2 fluxes and mechanisms promoting the formation of the O_2 maximum
361 zone in terms of primary productivity; 2) discuss the implication for the local O_2 BBL
362 dynamics and carbon budget; 3) elaborate on factors that can ultimately influence O_2
363 depletion in the North Sea and other seasonally stratified shelf seas.

364

365 **4.1 Thermocline mixing**

366 The expansive North Sea thermocline ($1 - 5 \times 10^5 \text{ km}^2$; Meyer et al., 2011)
367 has been regarded as being in a state of marginal stability, where additional sources of
368 shear could lead to increased thermocline mixing (e.g., van Haren et al., 1999).
369 Itsweire et al. (1989) showed that layers of strong shear are likely to be found where
370 strong stratification occurs. Generally, in the absence of varying topography and
371 sloping boundaries, the major sources of shear in the thermocline are considered to be
372 internal tides and near-inertial oscillations (see Rippeth, 2005). Sharples et al. (2007)
373 demonstrated that internal tidally-driven thermocline mixing enhanced diapycnal
374 nutrient fluxes, the overall productivity in the thermocline, and the associated carbon
375 export to the BBL.

376 The occurrence of near-inertial oscillations in shelf seas during the stratified
377 season has been reported in several studies from the North Sea (van Haren et al.,
378 1999; Knight et al., 2002) and in other shelf seas (e.g., Rippeth et al., 2002;
379 McKinnon and Gregg, 2005). During the presence of baroclinic inertial waves in the
380 water column, periods of enhanced shear have been observed in the western Irish Sea

381 (Rippeth et al., 2009), the Celtic Sea (Palmer et al., 2008) and the northern North Sea
382 (Burchard and Rippeth, 2009). These take the form of shear spikes, which occur
383 approximately every inertial period and in bursts lasting several days.

384 While we mainly attributed the observed enhanced turbulent mixing to the
385 occurrence of a near-inertial wave, the site's physical setting has further implications
386 for mixing processes in the thermocline. In the northern hemisphere, sites with anti-
387 cyclonic tides, such as Tommeliten, are often characterized by an increased vertical
388 extension of the BBL, and higher BBL dissipation rates than comparable cyclonic
389 sites (see Simpson and Tinker, 2009). As a result of this enhanced BBL thickness, we
390 observed sporadically elevated thermocline turbulence resulting from tidal-driven
391 bottom turbulence propagating vertically towards the thermocline (Fig. 5b). A study
392 by Burchard and Rippeth (2009) also reported that short lived thermocline shear
393 spikes can arise due to the alignment of the surface wind stress, bulk shear, and bed
394 stress vectors in the presence of baroclinic near-inertial motions and barotropic tidal
395 currents. These mechanisms are stronger with anti-cyclonic tides. Although all the
396 features required for shear spike generation were present during the observational
397 period, the two-layer mechanism described by these authors would require a more
398 complex water column structure to be applicable to the Tommeliten site.

399 The site's water column structure clearly showed the occurrence of a 10 m
400 thick transition layer (Fig. 2a). This layer represents the region of the water column
401 where mixing turns from elevated in the SBL to strongly reduced in the interior
402 (Ferrari and Boccaletti, 2004). The transition layer is therefore an obligate pathway
403 for solute and heat exchange between SBL and the interior (Ferrari and Boccaletti,
404 2004; Rhein et al., 2010) and has been reported to be a region of enhanced shear and
405 near-inertial wave activity (Dohan and Davis, 2011). Although the presented data did
406 not allow quantification of the O₂ exchange across the transition layer, such
407 contribution might be considerable and thus highly relevant for the cycling of O₂ and
408 CO₂ in the upper water column, which in turn could have direct biological
409 implications.

410

411 **4.2 BBL O₂ dynamics**

412 Ultimately, observed O₂ depletion in the BBL of the central North Sea
413 depends on the supply of organic matter, the rate of carbon mineralization, and the

414 flux of O₂ to the bottom water either from horizontal advection or turbulent vertical
415 transport. Our study investigated the significance of turbulent vertical O₂ fluxes to the
416 BBL, which has been previously overlooked in shelf sea carbon balances. Studies
417 focusing on O₂ replenishment in the BBL through the thermocline are limited to
418 freshwater systems (e.g. Bouffard et al., 2013; Kreling et al., 2014). In a large
419 stratified water body such as Lake Erie, O₂ transport from the thermocline to the
420 hypolimnion was found to be substantial, with a magnitude comparable to ~18% of
421 the hypolimnetic O₂ utilization rate over the whole stratification period (Bouffard et
422 al., 2013).

423 Horizontal O₂ gradients and associated horizontal advective O₂ fluxes were
424 not quantified in this study. Our data suggest, however, that such fluxes would not
425 significantly contribute to the O₂ balance at the Tommeliten site. BBL O₂
426 concentration time series (Fig. 5a) did not show any variability at the tidal and or
427 inertial frequencies, implying that horizontal O₂ gradients were small. Additionally,
428 mean currents in the BBL were small (~2 cm s⁻¹) compared to the tidal amplitudes.
429 This, in conjunction with weak horizontal O₂ gradients, suggests that horizontal
430 advective O₂ fluxes during our observational period are negligible compared to the
431 turbulent O₂ flux from the thermocline.

432 Based on the above, we can argue that the O₂ dynamics during the stratified
433 period are more complicated than previously regarded. To maintain an excess of O₂ in
434 the thermocline, primary producers require adequate nutrient entrainment from the
435 bottom water to fuel potential new production. The resulting increase in (new)
436 productivity and subsequent export to the bottom water could boost the carbon
437 turnover estimates substantially. Using a 1:1 O₂ utilization – carbon re-mineralization
438 (see Canfield, 1993), Greenwood et al. (2010) inferred the average BBL carbon re-
439 mineralization rate at the nearby North Dogger to be 15 mmol m⁻² d⁻¹, or 180 mg C m⁻²
440 d⁻¹. Similar results for a typical NW European shelf sea were obtained via modeling
441 by Sharples (2008), who reported rates ranging from ~35 to ~200 mg m⁻² d⁻¹ for neap
442 and spring tide, respectively. Their study, however, did not include the daily tidal
443 variation, and thus rates could be much higher on shorter timescales.

444 With the absence of targeted long-term studies focusing on O₂ and carbon
445 dynamics in the thermocline and BBL, we can only speculate on the long-term fate of
446 the BBL O₂ and its replenishment from the thermocline by vertical O₂ fluxes (F_{θ}). It

447 seems possible that the overall net BBL water column O₂ respiration, R , is higher than
448 previously thought, suggesting a much higher carbon turnover than inferred from the
449 apparent O₂ loss rate. Based on Eq. (6), the BBL carbon re-mineralization (and export
450 to the BBL) would be on the order of nearly 850 mg C m⁻² d⁻¹, nearly a factor of 5
451 higher than reported by Greenwood et al. (2010). However, the same turbulent
452 transport that supports the O₂ export from the DCM to the BBL also supports BBL
453 nutrient import to the DCM (Fig. 6). The higher import of nutrients to the DCM likely
454 promotes additional primary production and a subsequent increase in organic matter
455 (OM) export to the BBL. In such a scenario, the transient O₂ flux to the BBL
456 presented in this study will be associated with additional OM to the BBL, and
457 therefore lead to a temporary increased re-mineralization that offsets the increased F_{θ} .
458 While the overall effect is an increase in carbon turnover, this process would not
459 result in any observable change in the decreasing O₂ trend (apparent O₂ loss rate).

460

461 **4.3 Causes and controls on BBL O₂ depletion**

462 According to Boers (2005), for BBL O₂ to decrease throughout the stratified
463 season, there must be suitable physical conditions, biomass production, nutrient input
464 and continued benthic O₂ uptake. SUR , and thus the sediment nutrient release and
465 organic carbon mineralization have been shown to be strongly tidal-driven (McGinnis
466 et al., 2014). Therefore, we briefly discuss the potential tidal impact driving the
467 overall carbon cycling and suggest factors that may promote the development of
468 lower BBL O₂ concentrations during the stratification period.

469 Tidal forcing on diapycnal constituent fluxes and primary production have
470 been explored by e.g., Sharples et al. (2007, 2008). The authors showed that spring-
471 neap tide drives nutrient fluxes between the BBL and the DCM at the thermocline,
472 and the carbon export. Based on our velocity measurements and estimated O₂ fluxes,
473 we can expect similar patterns corresponding to semidiurnal tidal fluctuations. Blauw
474 et al. (2012) investigated fluctuating phytoplankton concentrations in relation to tidal
475 drivers and found that in the southern North Sea, chlorophyll fluctuations correlated
476 with the typical tidal current speed periods, the semidiurnal tidal cycle, in addition to
477 the day-night and spring-neap periods. During most of the year, chlorophyll and
478 suspended particulate matter fluctuated in phase with tidal current speed and indicated
479 alternating periods of sinking and vertical mixing of algae and suspended matter with

480 tidal cycles. Thus, these results suggest that in addition to the spring-neap tidal cycles,
481 we can expect a semidiurnal tidal-driven export of carbon and O₂ from the DCM to
482 the BBL, and entrainment of nutrients that strongly vary based on a timescale related
483 to the semi-diurnal tidal cycle.

484 The flux of O₂ from the DCM production zone downward to the BBL could
485 set the lower limit of the BBL O₂ concentration, and thus the O₂ depletion level,
486 during the stratification period. If there is little isolation between the zone of
487 production and the zone of mineralization, then the net O₂ production and O₂
488 utilization would nearly balance. In such case, the apparent O₂ loss in the BBL would
489 either be negligible or very small, depending whether the *SUR*, which is largely
490 particulate organic matter driven, will be balanced by the ventilation from the
491 thermocline. However, historically decreasing BBL O₂ concentrations within the
492 North Sea (Queste et al., 2013) point to an increasing disconnect between the primary
493 O₂ production zone and the mineralization zone. Greenwood et al. (2010) state that
494 stratification is an important factor which determines susceptibility to O₂ depletion,
495 especially in their nearby study site Oyster Grounds.

496 Surveys on the North Sea have shown that the regions with the lowest BBL O₂
497 concentrations are generally characterized by the strongest stratification (see Queste
498 et al., 2013), with the lowest values (~100 μmol kg⁻¹) reported to occur during
499 particularly calm and warm weather (see Boers, 2005; Weston et al., 2008). Strong
500 gradients in the thermocline associated with warmer temperature are suggested to
501 limit the O₂ flux to the BBL (Weston et al., 2008). This points to potential future O₂
502 depletion resulting from increasing temperatures leading to both stronger stratification
503 and a longer stratification season (Lowe et al., 2009). However, it could be argued
504 that if O₂ fluxes between the DCM and BBL were suppressed, then the upward
505 nutrient fluxes would be similarly suppressed, thus inhibiting primary production and
506 reducing the potential for O₂ deficits.

507

508 **4.4 Biological perspective**

509 The occurrence of stronger stratification and subsequently reduced turbulent
510 mixing could alter algal populations (Hickman et al., 2009), potentially favoring
511 migrating/swimming phytoplankton. An example of these migrating phytoplankton
512 species, armored dinoflagellates, are extensively found in the DCM of the central and

513 northern North Sea during the summer months; their abundance was found to be
514 largely determined by the local hydrodynamic conditions (Reid et al., 1990). In calm
515 conditions, which are typically associated with stronger stratification, there are often
516 blooms of migrating dinoflagellates which have access to the large nutrient pool in the
517 deeper water and can out-compete non-migrating species for both light and nutrients.
518 Stronger turbulent mixing, in contrast, has been suggested to interfere with their
519 swimming abilities (see Jephson et al., 2012 and references therein).

520 Algal migration could promote an upward shift of the DCM and move the
521 associated O₂ production higher in the thermocline where turbulence levels are
522 reduced, while still maintaining comparable production rates. Even by a few meters,
523 such an upward shift would substantially reduce turbulent O₂ fluxes to the BBL and
524 likely further isolate the BBL from the potential O₂ supply in the thermocline,
525 although maintaining similar rates of carbon export (settling armored dinoflagellates).
526 Studies on climate change impacts on the North Sea have suggested that O₂ loss in the
527 bottom waters would mainly result from a strengthening of the stratification and O₂
528 solubility reduction with increasingly warmer waters (e.g., Meire et al., 2013). In
529 those scenarios, the intricate interplay between local tidally-driven processes, water
530 column structure, biogeochemical cycling and active phytoplankton migration have
531 not been considered nor quantified. The proposed mechanism could contribute to the
532 observed decreasing O₂ levels in the North Sea water column, however, further
533 detailed studies are obviously necessary to validate and fully quantify this effect, and
534 the results described in this study, at the seasonal level.

535

536 **Acknowledgements**

537 We are thankful to the captain and crewmembers of the R/V *Celtic Explorer*
538 for their outstanding collaboration and support during the survey, Uwe Koy and
539 Rudolf Link for their logistic support, and Jens Schafstall, Tim Fischer and Markus
540 Faulhaber for their help in data collection and processing. We are grateful for the
541 technical development and support in deployment of the benthic chamber by Ralf
542 Schwarz, Sergiy Cherednichenko and the ROV Kiel 6000 team. Financial support was
543 provided by the Sonderforschungsbereich (SFB) 754 “Climate – Biogeochemistry in
544 the tropical Ocean”, SFB 574 “Volatiles and Fluids in Subduction Zones”, and by the
545 Excellence Cluster “Future Ocean” (project 2009/1 CP 0915, LR), supported by the

546 Deutsche Forschungsgemeinschaft (DFG). Additional founding was provided by the
547 National Environmental Research Council (NERC, project NE/J011681/1). The cruise
548 was financed by Wintershall within the Fluid and Gas Seepage in the Southern
549 German North Sea (SDNS) project.

550

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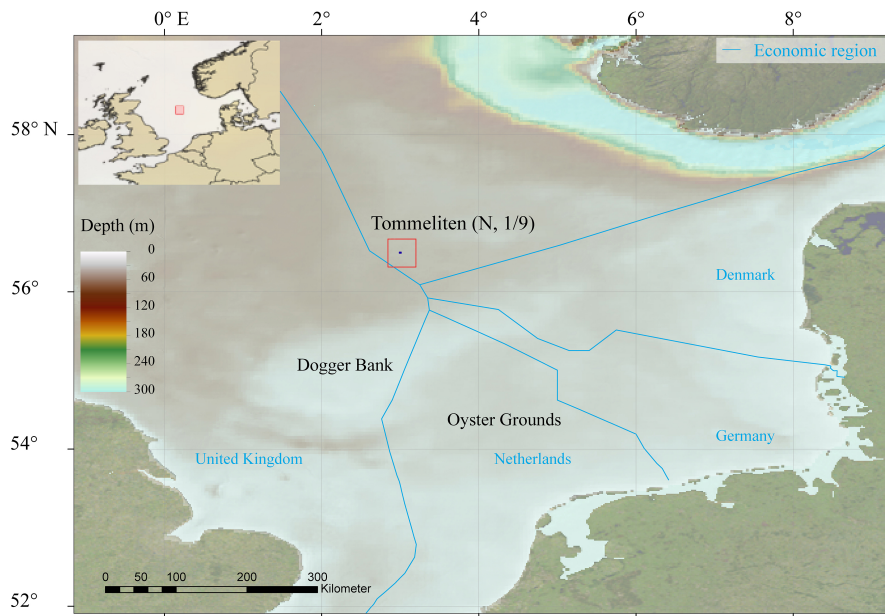
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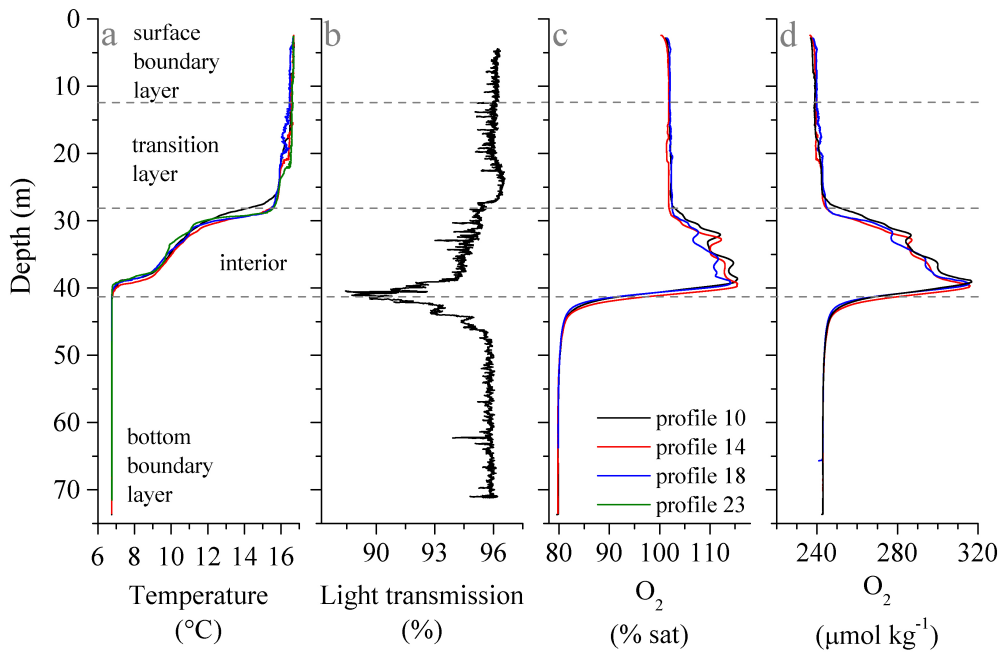
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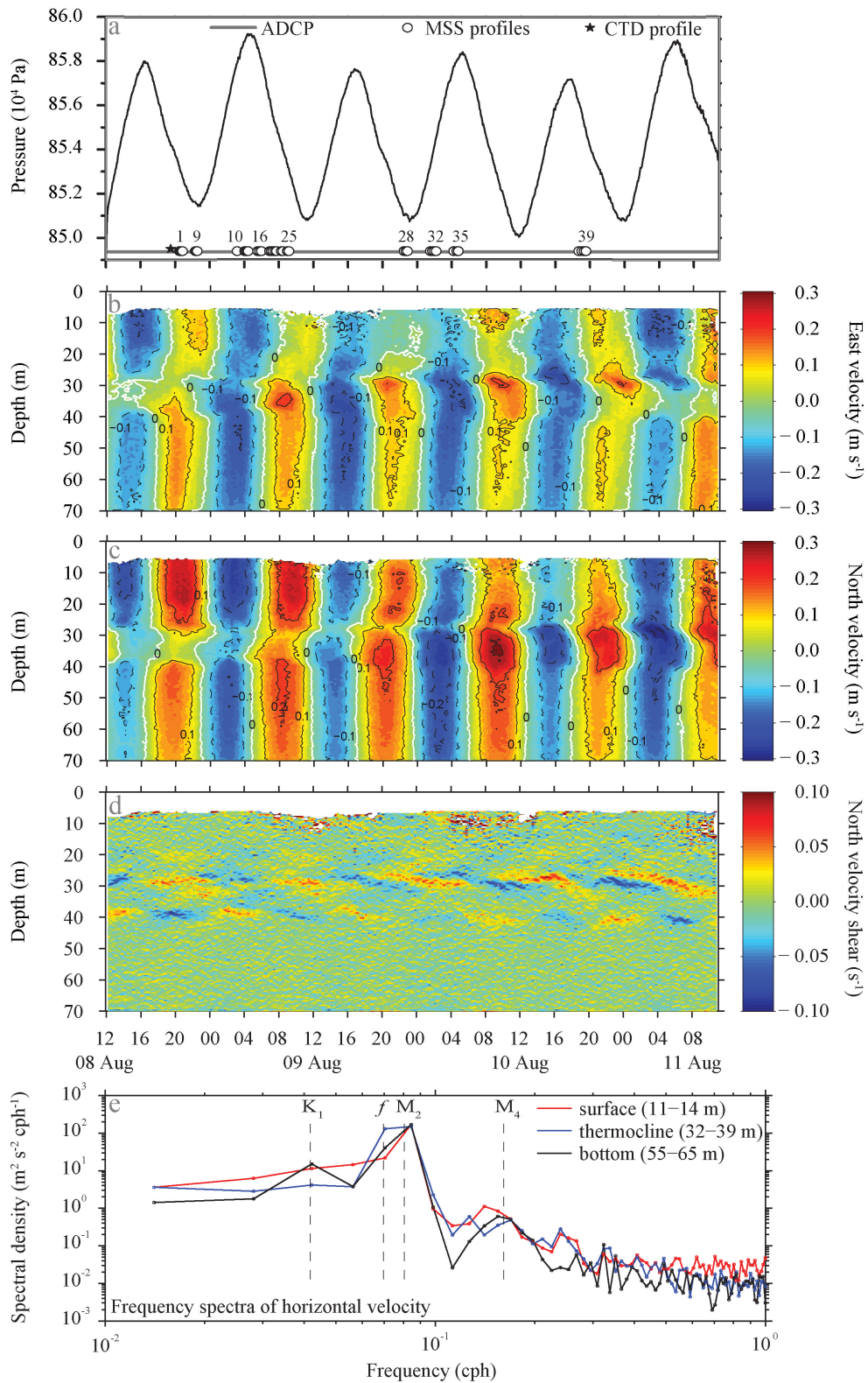
760 Figure 1. Map of the North Sea indicating the water depths and location of the
761 Tommeliten site and the borders of the economic regions of the surrounding European
762 countries.



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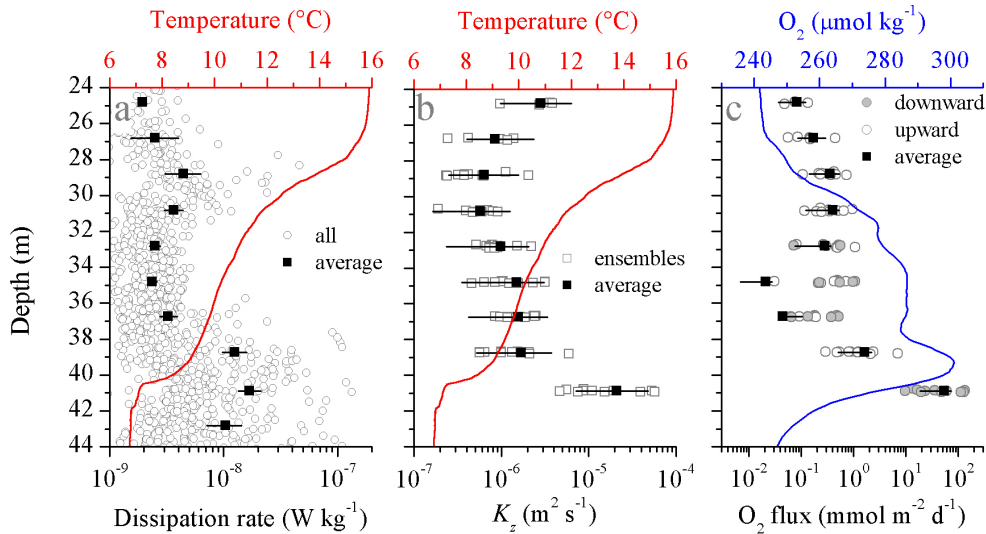
765 Figure 2. Selected water column profiles based on based on high-resolution MSS
 766 profiles (a, c, d) and ship CTD profile (b). (a) Potential temperature profiles. Water
 767 column layers were identified based on the temperature profiles. A 0.2°C and 1.5°C
 768 decrease from the surface boundary layer average temperature (3–6 m depth) was
 769 used determine the depth of the surface boundary layer – transition layer interface and
 770 the transition layer – interior interface, respectively. Correspondingly, a 0.2°C from a
 771 50-60 m depth average temperature was used to locate the interior – bottom boundary
 772 layer interface. (b) Light transmission profile. (c, d) O₂ saturation profiles and
 773 associated absolute concentrations.



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776 Figure 3. Current regime at the Tommeliten site from ADCP measurements (a - d)
777 and spectral analysis (e). (a) Sea surface elevation relative to average level during the
778 observational period (elevation = 0 m) and schedule of different instrument

779 deployments. Numbers on the MSS markers indicate the profile number. (b, c)
780 Horizontal velocities, showing 20 min averaged east (b) and north (c) velocities. (d)
781 Vertical shear of North velocity, dv/dz , calculated from the ADCP velocity data (see
782 panels b, c). Note that panels a - d have the same time axis. (e) Frequency spectra of
783 horizontal velocity calculated from the ADCP data for selected depth ranges for the
784 SBL (surface; red line), thermocline (blue line), and BBL (bottom; black line). The
785 inertial f , K_1 , M_2 and M_4 frequencies are marked.

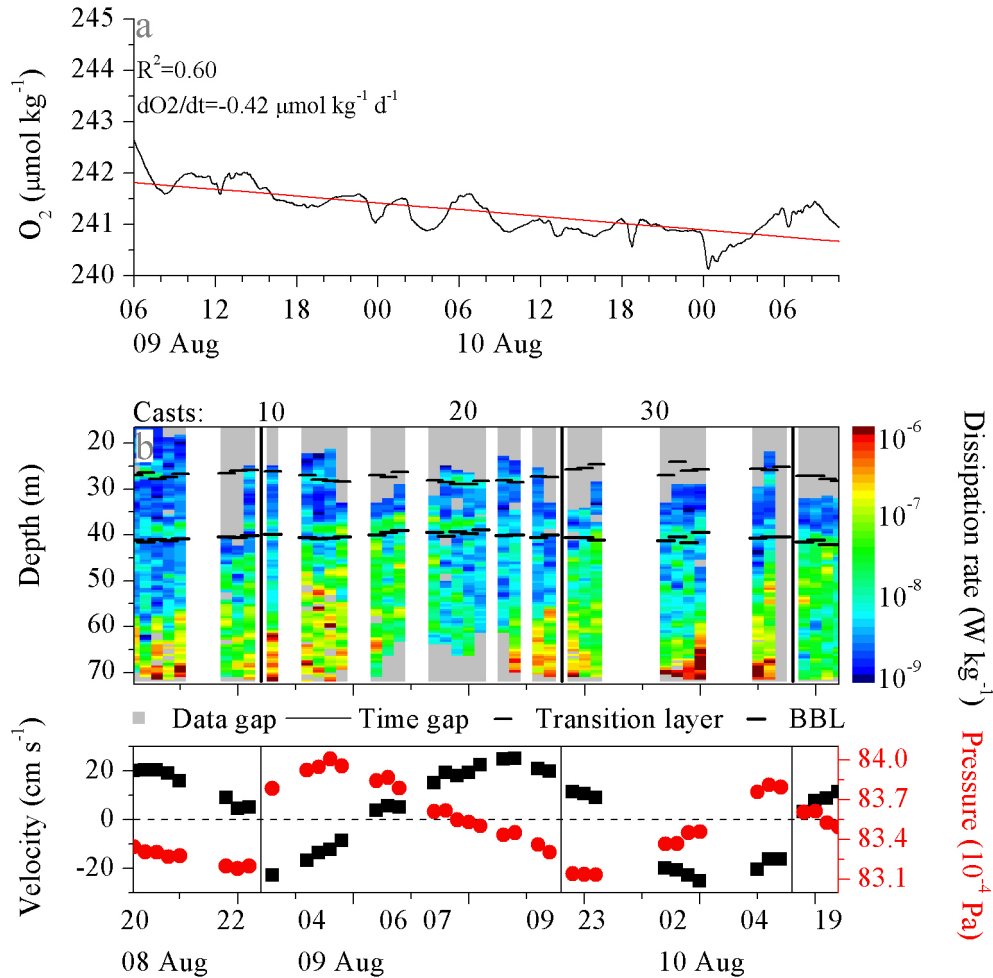


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788 Figure 4. Overview of turbulent transport and O₂ fluxes within the interior (defined in
 789 Fig. 2). Each panel is overlaid with temperature (a, b) and O₂ concentration (c)
 790 profiles. (a) Dissipation from all profiles (open dots) together with the arithmetic
 791 mean (solid squares). (b) Average vertical eddy diffusion coefficient K_z with
 792 uncertainties bars and the K_z values for every ensemble (open squares), which
 793 represent the average over 3 to 4 consecutive profiles. (c) Calculated average O₂ flux
 794 over 2 m bins with the respective uncertainties intervals (solid square and black line).
 795 The values for each profile cluster are shown both downward and upward fluxes (grey
 796 solid and open dots, respectively). Note that in the center interior (33 – 37 m) the
 797 average reflects the combination of the variability of the observed upward and
 798 downwards fluxes.

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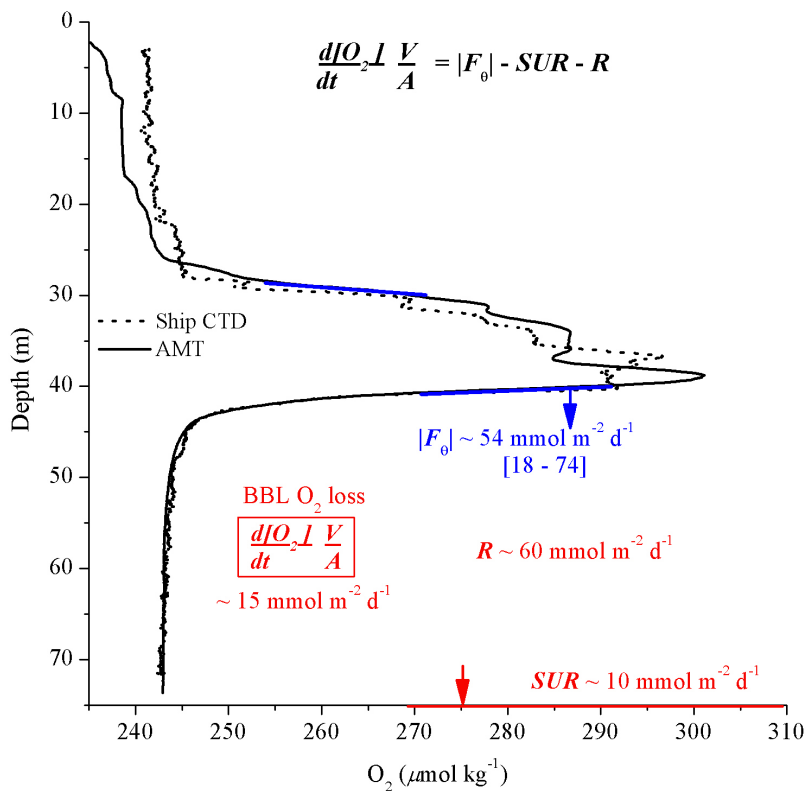


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802 Figure 5. BBL dissolved oxygen time series and turbulence contour. (a) Near-seafloor
 803 BBL O₂ concentration changes over the observational period from the POZ-Lander.
 804 Red line indicates the estimated apparent linear O₂ loss. (b, top) Turbulence contour
 805 plot of all MSS90 casts together with the temperature layers. Thin and thick dashed
 806 lines represent the transition layer – interior interface and the interior – BBL interface,
 807 respectively. Gray spots indicate data missing due to uncompleted profiles (casts 16-
 808 23), unsuccessful profiles (cast 36), or flagged as bad based on spikes, collisions and
 809 suspected contamination due to ship activity. The vertical black lines indicate the
 810 transition (time gaps) between consecutive profile ensembles. (b, bottom) Background
 811 information on bottom current, and hydrostatic pressure during the casts. Both
 812 velocity and pressure data were collected by the deployed POZ lander. Note that as a
 813 result of the time gaps between the consecutive MSS90 casts (see Fig. 3a) the time
 814 scale is not linear.

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819 Figure 6. Main O₂ fluxes in this study. The ranges shown for the interior O₂ fluxes
 820 refer to the associated uncertainty and intermittency levels. The sediment O₂ uptake
 821 rates (SUR)

822 are based on eddy correlation (EC) measurements (McGinnis et al., 2014), while
 823 central North Sea apparent BBL O₂ loss is based on Greenwood et al. (2010) and this
 824 study. Representative O₂ profiles are based on the AMT sensor on the MSS profiler
 825 (solid line) and ship CTD (dotted line). Note that while the O₂ profiles showed
 826 differences in absolute concentration within the thermocline, the actual O₂ gradients
 827 within the thermocline-BBL oxycline are comparable.