

Dear Dr Currie,

We thank you very much for your very attentive assistance in improving our manuscript clarity and readability. On behalf of all authors, I am pleased to resubmit our revised version of manuscript **Thermocline mixing and vertical oxygen fluxes in the stratified central North Sea** as an article in *Biogeosciences* within the special issue *Low oxygen environments in marine, fresh and estuarine waters*.

As suggested we have addressed terminology and wording comments throughout the manuscript. We now clearly state that CTD data were Winkler-corrected. We also made sure that figure labels and embedded text are properly referred in the figure caption or in the text. Text on Fig. 2 and Fig.5 were changed accordingly.

MSS profiles, which are the core dataset of this study, were mainly taken at nighttime or dusk/dawn. CTD casts were also collected at daytime, from both our ROV and the ship CTD. From the collected data, we did not observe a clear day/night shift in depth and strength of the oxycline. While we do not claim this was occurring during the short window of our study (we did not specifically look for this), we do propose it as an unconsidered mechanism explaining the overall decreasing O₂ trends in the North Sea. Our proposed migration-driven shift of the O₂ production zone is based on considerations from dedicated phytoplankton studies in the North Sea. This is also the reason we present it only in the discussion section. To clarify it, we now state that we could not observed migration or shifts of the O₂ production in our study.

We are confident the manuscript, after this fifth revision, is now suitable for publication in *Biogeosciences*.

On behalf of all authors

Lorenzo Rovelli

1 Abstract

2 In recent decades, the central North Sea has been experiencing a general trend of
3 decreasing dissolved oxygen (O₂) levels during summer. To understand potential
4 causes driving lower O₂, we investigated a three-day period of summertime
5 turbulence and O₂ dynamics in the thermocline and bottom boundary layer (BBL).
6 The study focuses on coupling biogeochemical with physical transport processes to
7 identify key drivers of the O₂ and organic carbon turnover within the BBL.
8 Combining our flux observations with an analytical process-oriented approach, we
9 resolve drivers that ultimately contribute to determining the BBL O₂ levels. We report
10 substantial turbulent O₂ fluxes from the thermocline into the otherwise isolated
11 bottom water attributed to the presence of a baroclinic near-inertial wave. This
12 contribution to the local bottom water O₂ and carbon budgets has been largely
13 overlooked and is shown to play a role in promoting high carbon turnover in the
14 bottom water while simultaneously maintaining high O₂ concentrations. This process
15 may become suppressed with warming climate and stronger stratification, conditions
16 which could promote migrating algal species that potentially shift the O₂ production
17 zone higher up within the thermocline.

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19 1 Introduction

20 1.1 Hypoxia in shelf seas and coastal regions

21 The distribution of dissolved oxygen (O₂) in marine systems results from a
22 complex interaction between biological processes (photosynthesis and respiration)
23 and physical processes (O₂ flux pathways) occurring within the water column and at
24 the seafloor. O₂ is regarded as an important indicator of ecosystem functioning for
25 aquatic organisms (Best et al., 2007) and for benthic activity (e.g., Glud, 2008).
26 Changes in O₂ distribution, concentrations and supply can therefore have severe
27 impacts on shelf ecosystems. O₂ concentrations below 62.5 μmol L⁻¹, which is
28 generally regarded as the threshold of hypoxia (Vaquer-Sunyer and Duarte, 2008), are
29 shown to significantly stress aquatic communities and increase the mortality among
30 fish communities (Diaz, 2001). These ecological and economic impacts of O₂
31 depletion lead to increasing concern regarding hypoxia occurrence and hypoxic
32 events. As reviewed by Diaz and Rosenberg (2008), hypoxia in coastal environments
33 is spreading and so are the reports of unprecedented occurrence of hypoxia in several

38 shelf seas and coastal regions (Grantham et al., 2004; Chan et al., 2008; Crawford and
39 Pena, 2013).

40

41 **1.2 Hydrodynamics and oxygen depletion in the North Sea**

42 The North Sea is situated on the North–West European continental shelf,
43 between the British Islands and continental Europe (Fig. 1). Its semi-enclosed basin
44 covers an area of 575'300 km², with an average depth of 74 m and a general decrease
45 in depth from North to South (Otto et al., 1990). The central region is characterized
46 by the presence of the Dogger Bank, a shallow sandbank that acts as a hydrological
47 divide. The northern and central North Sea hydrology is mainly dominated by inflow
48 from the North Atlantic Ocean at the northern open boundary, while the southern part
49 relies on inflow from the English Channel (Thomas et al., 2005). Northern and central
50 North Sea areas are characterized by seasonal water column stratification (April to
51 September - October; Meyer et al., 2011). With only weak, wind-driven residual
52 currents (Otto et al., 1990), this stratification leads to isolation of central North Sea
53 bottom water and subsequent O₂ depletion.

54 In the central North Sea, the occurrence of low O₂ levels in bottom waters has
55 been reported (e.g., North Sea Task Force, 1993; Greenwood et al., 2010). Additional
56 monitoring studies in the central North Sea in 2007 and 2008 have shown that O₂
57 concentration in bottom waters at the Oyster Grounds and North Dogger can drop to
58 163 – 169 μmol L⁻¹ (60 – 63% saturation) and ~200 μmol L⁻¹ (71% saturation),
59 respectively (Fig. 1; Greenwood et al., 2010). Comparable field observations were
60 also reported in the summer of 2010 (Queste et al., 2013). The authors also reviewed
61 the available historical O₂ data in the North Sea (1900 – 2010), revealing a clear
62 increase in O₂ depletion after 1990.

63 While the reported O₂ levels were still above the hypoxic threshold, growing
64 concerns of hypoxia developing in the North Sea have highlighted the need for more
65 detailed studies on O₂ dynamics and its driving forces (Kemp et al., 2009). Since
66 | 1984, surface water temperatures in the North Sea have increased by 1 – 2°C, [more](#)
67 than the global mean (OSPAR, 2009, 2010; Meyer et al., 2011). On seasonal time
68 scales, climate projections indicate longer durations of the stratification period and
69 stronger thermocline stability (Lowe et al., 2009; Meire et al., 2013), with some
70 projections suggesting earlier onset of stratification (e.g., Lowe et al., 2009). Due to

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72 the semi-enclosed nature of the North Sea, earlier onset and longer stratification
73 increases the length of time that deep waters are isolated, potentially allowing lower
74 O₂ concentrations to develop (Greenwood et al., 2010).

75

76 **1.3. Physical drivers of oxygen dynamics**

77 The distribution of O₂ and other dissolved constituents within aquatic systems
78 are largely driven by physical transport processes. These include wind driven air –
79 water gas exchange at the sea surface (Wanninkhof, 1992), molecular diffusion at the
80 sediment – water interface (Jørgensen and Revsbech, 1985), horizontal advection
81 (e.g., Radach and Lenhart, 1995) and turbulent transport in the water column, where
82 the latter transport was reported to significantly contribute to constituent balances (see
83 Rippeth, 2005; Fischer et al., 2013; Kreling et al., 2014; Brandt et al., 2015). In shelf
84 seas, the seasonal occurrence of steep thermoclines acts as an important physical
85 barrier separating the surface layer from nutrient-rich deeper waters (Sharples et al.,
86 2001). As measurements of shear and stratification have shown, the central North Sea
87 thermocline is in a state of marginal stability (van Haren et al., 1999). Hence
88 additional sources of shear could trigger shear instability leading to local production
89 of turbulence within the thermocline. This enhanced local turbulence would
90 subsequently enhance the vertical exchange of constituents such as O₂, organic carbon
91 and nutrients. Resolving processes that drive diapycnal (i.e., vertical) fluxes across the
92 thermocline throughout the stratification period is key to understanding the
93 biogeochemical functioning of shelf seas (e.g., Sharples et al., 2001).

94

95 **1.4 Present study**

96 The goal of this study is to provide understanding of key turbulent processes
97 driving O₂ fluxes across the thermocline during the summertime stratification period
98 in the central North Sea using data from a 3-day process study. We investigate and
99 describe O₂ dynamics and fluxes to the bottom waters and discuss their potentially
100 influence on the seasonal O₂ balance. Using the resolved O₂ fluxes, we perform a
101 simple 1-D mass balance model to quantify O₂ sources and sinks. Finally, processes
102 that could further promote hypoxia in the central North Sea in a warming climate are
103 discussed.

104

105 **2 Methods**

106 **2.1 Study site**

107 We performed O₂ and turbulence measurements in the Norwegian sector of
108 the central North Sea, N. 1/9, at the Tommeliten site (56°29'30" N, 2°59'00" E; Fig. 1)
109 from 8 – 11 August 2009 aboard the R/V *Celtic Explorer* (cruise CE0913). The site,
110 located ~100 km northeast from the northern Dogger Bank, and its surroundings are
111 characterized by shallow waters (~70 m) relatively far from coastal areas (on average
112 ~300 km). The site is known for the presence of buried salt diapirs, methane (CH₄)
113 seeps and bacterial mats (Hovland and Judd, 1988). Bathymetric surveys from
114 Schneider von Deimling et al. (2010) revealed a rather flat sandy seabed with almost
115 no features, with the exception of cm-sized ripples (McGinnis et al., 2014).

116 The currents of the central North Sea are predominantly driven by the semi-
117 diurnal lunar tide (M₂) (Otto et al., 1990). Seasonal stratification begins in April
118 around Julian day 100 and lasts until the end of September or early October, Julian
119 days 270- 290 (e.g. Meyer et al., 2011). The thermocline has been identified as an
120 important zone for the establishment of primary production and the O₂ maximum
121 layer (see Pingree et al., 1978). In fact, the North Sea deep chlorophyll maximum
122 (DCM) is estimated to account for 58% of the water column primary production and
123 37% of the annual new production for the summer stratified North Sea (Weston et al.,
124 2005). The development of the associated O₂ maximum due to this production is thus
125 important and so far not considered in the overall O₂ balance of the central North Sea.

126

127 **2.2 Instrumental setup**

128 High resolution (mm scale) turbulent shear and temperature profiles were
129 obtained with a MSS90-L microstructure turbulence profiler (Sea and Sun
130 Technology, Trappenkamp, Germany). The MSS90-L is a free-falling, loosely-
131 tethered profiler which samples at 1024 Hz with 16 channels and is designed for an
132 optimal sink rate of 0.5 – 0.6 m s⁻¹. The probe was equipped with two air-foil shear
133 probes, an accelerometer (to correct for probe pitch, roll, and vibration), a fast
134 temperature sensor (FP07, 7–12 ms response time), standard CTD sensors
135 (temperature, pressure, conductivity), and a fast (0.2 s response time) galvanic O₂
136 sensor (AMT, Analysenmesstechnik GmbH, Rostock, Germany). Absolute O₂

137 concentrations were calibrated against shipboard CTD O₂ profiles and Winkler
138 titrations using discrete water samples (see below).

139 Water column hydrodynamics were characterized with the compact benthic
140 Paleoceanography (POZ) lander, which was deployed using a video guided launcher
141 (Pfannkuche and Linke, 2003). The POZ lander recorded 3-dimensional current
142 velocity profiles and acoustic backscatter information throughout the water column
143 using a 300 kHz acoustic Doppler current profiler (ADCP; Workhorse Sentinel,
144 Teledyne RD Instruments, Poway, United States), which sampled every 15 s with a
145 bin size of 0.5 m starting from 2.75 m from the bottom. A conductivity-temperature-
146 depth (CTD) logger (XR-420 CT logger, RBR, Kanata, Canada) recorded
147 temperature, conductivity and pressure (Digiquartz, Paroscientific, Redmond, United
148 States) every 2 s near the seafloor (~0.3 m distance). The POZ lander was also
149 equipped with a Winkler-calibrated O₂ optode sensor (Aanderaa Data Instruments
150 AS, Bergen, Norway), which recorded BBL O₂ concentration at 1 min intervals.

151 Water column profiles were obtained using a SBE9plus CTD-rosette system
152 (Seabird, Washington, United States). The CTD sampled at 24 Hz and was equipped
153 with standard temperature, conductivity, pressure, O₂ and light transmission sensors.
154 The rosette system mounted 12 Niskin bottles (10 L each) for discrete water
155 sampling. Each water sample was subsampled with three Winkler bottles of known
156 volume (~62 mL on average) upon recovery, and the samples were immediately fixed
157 on deck and titrated manually within 24 h after the sampling (see Winkler 1888).
158 CTD O₂ readings were calibrated against O₂ concentrations from Winkler samples;
159 values deviated by <5%.

160

161 2.3 Hydrodynamic data evaluation

162 The main tidal directions (major and minor axis of the tidal ellipsoid) were
163 determined by performing a variance analysis on the ADCP velocity time series. The
164 u and v velocities were rotated over a stepwise increasing rotation angle (r) as
165 $u_{rot} = u \cdot \cos(-r) - v \cdot \sin(-r)$ and $v_{rot} = u \cdot \sin(-r) - v \cdot \cos(-r)$, and the
166 variance computed at each step. The angle with the largest variance is the main tidal
167 direction. Barotropic and baroclinic flow contributions of tides were separated by
168 least-square fitting the detrended velocity time series to harmonics $u = A \cdot \cos(\omega \cdot$
169 $t + \varphi)$ with A , ω , φ being the amplitude, frequency, and the phase lag, respectively.

171 In the analysis below, the barotropic semi-diurnal principle lunar tide (M_2) and
 172 diurnal declination tide (K_1) contributions had frequencies of 1.93227 cycles per day
 173 (cpd) and 1.00274 cpd, respectively, and were subtracted from the time series to
 174 analyze residual flow. For barotropic contributions, the fit was applied to the depth
 175 average of the time series, while baroclinic contributions were obtained by fitting the
 176 harmonics to the velocity time series from each 0.5 m ADCP bin. The occurrence of
 177 enhanced shear in the stratified water column was investigated by calculating the
 178 vertical shear of horizontal velocity, S , from the vertical gradients between adjacent
 179 bins of east and north velocity (0.5 m resolution) as $S = \sqrt{(du/dz)^2 + (dv/dz)^2}$.
 180 Frequency spectra of the time series of horizontal velocity and vertical shear of
 181 horizontal velocity were used to identify the tidal and non-tidal flow components. The
 182 spectra were calculated using fast-Fourier transforms combined with a $1/2$ -cosine
 183 taper (Hanning window), which was applied to the first and last 10% of the time
 184 series data.

185 Turbulent kinetic energy dissipation rate (ε) was quantified from airfoil shear
 186 readings by integrating shear wavenumber spectra assuming isotropic turbulence
 187 (Batchelor, 1953):

$$\varepsilon = 7.5\mu \int_{k_{min}}^{k_{max}} E_{du'/dz}(k) dk \quad (1)$$

188 where μ is the dynamic viscosity of seawater. Shear spectra $E_{du'/dz}(k)$ were
 189 calculated from one-second ensembles (1024 values) and integrated between a lower
 190 $k_{min} = 3$ cycles per minute (cpm) and an upper wavenumber k_{max} that varied between
 191 14 cpm and 30 cpm depending on the Kolmogorov wavenumber. Here, a Bartlett
 192 window was applied to the whole ensemble prior to spectral decomposition. Loss of
 193 variance due to limited wavenumber band was taken into account by fitting the
 194 observed shear spectra to the universal Nasmyth spectrum. Similarly, corrections for
 195 the loss of variance due to finite sensor tip of the airfoil probes were applied (see
 196 Schafstall et al., 2010). The detection limit, or noise level, of the used profiler for ε
 197 was inferred to be 1×10^{-9} W kg⁻¹ (Schafstall et al., 2010); the upper detection limit is
 198 a function of the shear sensor geometry (up to 10^{-4} W kg⁻¹; Prandke and Stips, 1998).

199 Estimates of turbulent eddy diffusivities of mass (K_ρ) were obtained from
 200 measurements of ε as

$$K_\rho = \gamma \varepsilon / N^2 \quad (2)$$

201 where γ is the mixing efficiency and N^2 the water column stability. This method,
 202 proposed by Osborn (1980), approximates K_ρ under the assumption of a local
 203 equilibrium of production and dissipation of turbulent kinetic energy. Values for N^2
 204 were calculated from temperature, salinity and pressure data using the adiabatic
 205 method (Fofonoff, 1985) as $N^2 = -g(\rho^{-1}\partial\rho/\partial z - g/c^2)$, where ρ , g , and c are the
 206 density, the earth's gravitational acceleration and speed of sound. Mixing efficiency
 207 values in stratified waters range from 0.1 to 0.2 (Ivey and Imberger, 1991) and
 208 decreases in weakly stratified waters such as within the BBL (Lorke et al., 2008). To
 209 account for this decrease, we used the γ and K_ρ parameterization of Shih et al. (2005).
 210 Based on the turbulence activity parameter $\varepsilon/\nu N^2$, with the kinematic viscosity, ν ,
 211 the authors found that in energetic regimes, i.e., $\varepsilon/\nu N^2 > 100$, eddy diffusivities are
 212 better estimated as $K_\rho = 2\nu(\varepsilon/\nu N^2)^{1/2}$. As horizontal density gradients at the study
 213 site were deemed to be small compared to vertical gradients (see Discussion), we
 214 equated diapycnal eddy diffusivities with vertical diffusivities (i.e., $K_\rho = K_z$).

215 To obtain representative mean turbulent eddy diffusivities, the data were
 216 evaluated in ensembles of three to four consecutive profiles and averaged in depth and
 217 time to reduce uncertainties due to the patchiness of turbulence, temporal fluctuation
 218 of N^2 , and temporal γ variations (see Smyth et al., 2001). As proposed by Ferrari and
 219 Polzin (2005), the level of uncertainty of the averaged K_z can be quantified as:

$$\Delta K_z = K_z \left[\left(\frac{\Delta\gamma}{\gamma} \right)^2 + \left(\frac{\Delta\varepsilon}{\varepsilon} \right)^2 + \left(\frac{\Delta N^2}{N^2} \right)^2 \right]^{1/2} \quad (3)$$

220 with Δ being the absolute uncertainty of the various average terms. Here, the
 221 uncertainties are evaluated in the region of strong vertical O_2 gradients and in 2 m
 222 depth bins. The absolute uncertainty for $\Delta\gamma$ was assumed to be 0.04 (see St. Laurent
 223 and Schmitt, 1999). The absolute uncertainty on N^2 (ΔN^2) was determined by the
 224 standard error over the 2 m average, computed as the standard deviation divided by
 225 the square root of the number of estimates. Finally, the statistical uncertainty of ε for
 226 each bin was calculated using a bootstrap method (10^4 resamples) (Efron, 1979).

227 The vertical O_2 fluxes F_θ were then obtained from K_z and the O_2 concentration
 228 gradients $\partial[O_2]/\partial z$ as

$$F_{\theta} = K_z \frac{\partial [O_2]}{\partial z} \quad (4)$$

229 Accordingly, the uncertainty of averaged turbulent O₂ fluxes were given by:

$$\Delta F_{\theta} = F_{\theta} \left[\left(\frac{\Delta K_z}{K_z} \right)^2 + \left(\frac{\Delta \partial_z [O_2]}{\partial_z [O_2]} \right)^2 \right]^{1/2} \quad (5)$$

230 where $\Delta \partial_z [O_2]$ denotes the standard error of mean vertical gradients of O₂
 231 concentrations. It should be noted that the analysis did not include biases or
 232 uncertainties due to measurement errors.

233

234 **3 Results**

235 During the three-day observational period (8 – 11 August 2009), we collected
 236 39 high-resolution MSS profiles in consecutive sets of three to five profiles at 5 – 10
 237 min intervals. Most of the profiles were in the evening (profiles 1 – 8, 26 – 28, 36 –
 238 39) or at night (9 – 15, 29 – 35) with the remaining profiles acquired in the morning
 239 (6 to 9 AM). One shipboard CTD profile was performed prior to the actual MSS
 240 profiles to provide hydrographic information, the water turbidity and O₂
 241 concentrations, and discrete water samples for subsequent onboard Winkler titrations.
 242 Hydroacoustic water column current measurements were carried out continuously
 243 throughout the observational period. The following results are structured to first
 244 present a characterization of the site's physical settings and turbulence drivers,
 245 followed by the O₂ fluxes and O₂ BBL budget.

246

247 **3.1 Water column structure**

248 The ~70 m deep water column was characterized by a stable, well-defined
 249 four-layer temperature structure (Fig. 2a). The well-mixed surface boundary layer
 250 (SBL) and bottom boundary layer (BBL), 15 m and 30 m thick, respectively, were
 251 separated by a weakly-stratified transition layer (15 – 25 m depth) and a strongly
 252 stratified interior layer (25 – 40 m depth). The stratified interior layer was
 253 characterized by two very steep thermoclines situated in the upper (27 – 30 m depth)
 254 and lower (36 – 39 m depth) region of the layer, with vertical temperature gradients of
 255 up to 4°C m⁻¹. The average salinity was 35.08 with little variation throughout the
 256 water column (35.04 – 35.1). The light transmission profile from the ship CTD ranged
 257 from 89% to 96% (Fig. 2b). The most turbid layer (89%) was observed at the lower

258 boundary of the interior layer (at 40 m depth) suggesting the presence of the deep
259 chlorophyll maximum, phytoplankton, zooplankton and suspended particles.

260 The O₂ profiles were generally characterized by near saturation in the SBL
261 and transition layers (238 – 243 μmol kg⁻¹) and undersaturated (~80%) in the BBL
262 (~243 μmol kg⁻¹) (Fig. 2c,d). The stratified interior layer was oversaturated by up to
263 115%, with a well-established O₂ maximum at ~39 m depth with concentrations up to
264 ~315 μmol kg⁻¹. Below that maximum, at the thermocline-BBL interface, we
265 observed a 2 – 3 m thick steep oxycline, with an O₂ gradient of 34 μmol kg⁻¹ m⁻¹. We
266 did not detect significant variations in the O₂ values, or gradient over our limited
267 measurement range (Fig. 2c,d). We resolve the O₂ flux into the BBL associated with
268 this oxycline.

269

270 3.2 Hydrodynamics

271 The hydrostatic pressure dataset (POZ lander) revealed that the tidal water
272 level ranged from 0.6 to 0.9 m (Fig. 3a). Variance analysis on the ADCP velocity data
273 identified the major and minor axis of the tidal ellipsoid components to occur at 45°
274 and 135° from true north, respectively. Along these axes, the current amplitudes were
275 0.21 m s⁻¹ and 0.04 m s⁻¹, indicating a narrow tidal current ellipsoid, as reported by
276 Otto et al. (1990). The site was characterized by a negative tide polarity (anti-
277 cyclonic) for the semi-diurnal tides. A dominance of the barotropic M₂ current
278 amplitude at all depths was clearly visible in the velocity time series (Fig. 3b, c) and
279 the harmonic analyses. East (zonal) and north (meridional) barotropic M₂-current
280 amplitudes were 0.12 m s⁻¹ and 0.17 m s⁻¹, respectively, while K₁-current amplitudes
281 were only 0.005 m s⁻¹ and 0.03 m s⁻¹.

282 Although the limited length of the ADCP velocity time series did not allow for
283 full separation of the M₂ and *f* frequencies, the spectral density functions indicated
284 maximum energy at frequencies of about the semi-diurnal tide. This maximum varied
285 little with depth, indicating barotropic M₂ motions. Superimposed on those barotropic
286 currents, we observed the presence of baroclinic velocity contributions (Fig. 3b, c).
287 Additionally, near-inertial motions were also detected.

288 The occurrence of near-inertial motions was most pronounced in the
289 thermocline (32 – 39 m; Fig. 3e). Lower, but still elevated, energy densities at the

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292 near-inertial band were found in the SBL and BBL. Moreover, the near-inertial
293 currents exhibited a distinct 180° phase shift between the SBL and the thermocline as
294 well as between the thermocline and the BBL, suggesting a second vertical mode
295 nature of these fluctuations. Average amplitudes of the near-inertial fluctuations in the
296 thermocline obtained from least-square fitting were 0.11 m s⁻¹. In the BBL and SBL,
297 average amplitudes were reduced to 0.06 m s⁻¹ and 0.04 m s⁻¹, respectively,
298 suggesting that f oscillations might account for enhanced shear in the thermocline.

299 | Enhanced vertical shear of horizontal velocity was found at the interior layer –
300 | transition layer and at the interior layer – BBL interfacial regions (Fig. 3d). As
301 | indicated by the spectral density function of the shear time series from the interior
302 | layer interfacial layers (SI Fig. 1), the shear exhibited near-inertial frequencies
303 | (1.6722 cpd), and resulted from the baroclinic near-inertial wave. The high vertical
304 | resolution (0.5 m) of our velocity data allowed the resolution of the interfacial shear
305 | layers, which were typically 2 to 3 m thick with elevated values of up to 0.05 s⁻¹.
306 | Comparisons with CTD data showed that they are collocated with the two enhanced
307 | temperature gradients layers in the thermocline (27 – 30 m and 36 – 39 m depth; Fig.
308 | 2a).

309 | The dissipation rates (ϵ) of turbulent kinetic energy (TKE) determined from
310 | microstructure shear probes were particularly low in the center of the stratified
311 | interior layer ($2 - 5 \times 10^{-9}$ W kg⁻¹) but still above the MSS detection limit. However, ϵ
312 | increased to 5×10^{-9} W kg⁻¹ and 2×10^{-8} W kg⁻¹ at the upper and lower interior layer
313 | limits, respectively (Fig. 4a). These coincided with the depth range of the interfacial
314 | shear layers (Fig. 3d) at the strong temperature gradients (Fig. 2a) and resulting water
315 | column stability maxima ($\sim 1 \times 10^{-3}$ s⁻²).

316 | Bin-averaged values of K_z varied by a factor of 5, ranging from 6×10^{-7} m² s⁻¹
317 | in the central interior layer to 3×10^{-6} m² s⁻¹ in the lower region of the transition layer
318 | (Fig. 4b). In the upper interface (thermocline – transition layer), where ϵ was elevated
319 | with respect to the central interior layer but reduced compared to the lower interfacial
320 | layer, stronger stratification (i.e., larger N^2 values up to 10^{-3} s⁻²) reduced the eddy
321 | diffusivities. At the interior layer – BBL interface, higher K_z values ($\sim 2 \times 10^{-5}$ m² s⁻¹)
322 | resulted from increased turbulence and weaker stratification. This enhanced turbulent
323 | transport was located where the vertical O₂ gradient was the strongest (Fig. 2d).

324

326 3.3 Oxygen fluxes and budget

327 With the fast responding AMT galvanic O₂ sensor and rapid sampling rate, we
328 were able to resolve the O₂ gradient with high precision. Figure 4c shows the 2 m bin
329 average O₂ fluxes for the interior [layer](#) together with the averages from each
330 ensemble. Small O₂ fluxes (~1 mmol m⁻² d⁻¹) were estimated for the center and upper
331 region of the interior [layer](#); suggesting that relatively little O₂ is transported upward
332 from the O₂ maximum to the upper interior [layer](#). In contrast, a substantial O₂ flux
333 ranging from 9 – 134 mmol m⁻² d⁻¹ (average of 54 mmol m⁻² d⁻¹) was identified from
334 the lower thermocline towards the BBL. The confidence interval associated with the
335 uncertainties of the O₂ flux estimates was 18 – 74 mmol m⁻² d⁻¹. Although the O₂
336 fluxes to the BBL water from the thermocline were variable in magnitude (Fig. 4c)
337 and the measurements limited to the observational period (Fig. 3), their magnitude
338 nevertheless suggests an important, yet overlooked, O₂ pathway.

339 We performed a simple 1-D BBL mass balance to investigate the relevance to
340 the local O₂ balance during our observational period. Here, we defined the apparent
341 (measured) O₂ loss rate in the BBL $\partial[O_2]/\partial t$ as the consequence of O₂ replenishment
342 from F_θ and the O₂ utilization via sediment O₂ uptake rate (SUR) and water column
343 organic matter respiration (R) expressed as

$$\frac{\partial[O_2]V}{\partial t A} = |F_\theta| - |SUR| - |R| \quad \{mmol m^{-2} d^{-1}\} \quad (6)$$

344 The mass balance was constrained to the (assumed) well-mixed 35 m deep BBL
345 section of area, $A = 1 m^2$ with a volume $V = 35 m^3$. We further assumed negligible
346 horizontal O₂ gradients (as observed from the CTD casts), and thus a net zero
347 horizontal O₂ advective transport.

348 The average SUR for the same time period and location, obtained from parallel
349 eddy correlation measurements, was ~-10 mmol m⁻² d⁻¹ (McGinnis et al., 2014). The
350 SUR was consistent with the average SUR at Oyster Grounds reported by Neubacher
351 et al. (2011), -9.8 mmol m⁻² d⁻¹, and with modeled $SURs$ at the same site (average -8.6
352 mmol m⁻² d⁻¹; Meire et al., 2013). The apparent BBL O₂ loss of -0.42 μmol kg⁻¹ d⁻¹
353 was determined from the POZ lander O₂ optode time series (Fig. 5a) over 52 hours,
354 ($R^2=0.60$). Though limited to our short observational period, the vertically integrated
355 apparent BBL O₂ loss was about -15 mmol m⁻² d⁻¹ and thus within 2% of the nearby
356 North Dogger average presented by Greenwood et al. (2010). Based on Eq. (6) and

357 using the observed BBL O₂ loss rate, F_{θ} and SUR , the water column respiration, R
358 was calculated to be $\sim 60 \text{ mmol m}^{-2} \text{ d}^{-1}$. This implies that without the O₂
359 replenishment, the apparent BBL O₂ loss would be $\sim 2 \text{ } \mu\text{mol kg}^{-1} \text{ d}^{-1}$ and thus four
360 times higher than observed. Our results indicated that the total respiration in the
361 bottom water was therefore $\sim 70 \text{ mmol m}^{-2} \text{ d}^{-1}$ ($SUR + R$), with about 14% of the
362 organic carbon mineralization occurring in the sediment and 86% in the BBL.

363

364 **4 Discussion**

365 During our three-day observational period, we found that the baroclinic near-
366 inertial wave in the interior [layer](#) was the main contributor to the detected enhanced
367 shear (Fig. 3d) and the observed elevated vertical O₂ flux to the BBL (Fig. 6). As
368 near-inertial waves decay after a few weeks, it should be noted that we observed a
369 rather special situation, and that vertical O₂ fluxes will not likely be as highly elevated
370 during periods when near-inertial waves are not present.

371 Within this context, we: 1) discuss the turbulent mechanisms leading to these
372 thermocline O₂ fluxes and mechanisms promoting the formation of the O₂ maximum
373 zone in terms of primary productivity; 2) discuss the implication for the local O₂ BBL
374 dynamics and carbon budget; 3) elaborate on factors that can ultimately influence O₂
375 depletion in the North Sea and other seasonally stratified shelf seas.

376

377 **4.1 Thermocline mixing**

378 The expansive North Sea thermocline ($1 - 5 \times 10^5 \text{ km}^2$; Meyer et al., 2011)
379 has been regarded as being in a state of marginal stability, where additional sources of
380 shear could lead to increased thermocline mixing (e.g., van Haren et al., 1999).
381 Itsweire et al. (1989) showed that layers of strong shear are likely to be found where
382 strong stratification occurs. Generally, in the absence of varying topography and
383 sloping boundaries, the major sources of shear in the thermocline are considered to be
384 internal tides and near-inertial oscillations (see Rippeth, 2005). Sharples et al. (2007)
385 demonstrated that internal tidally-driven thermocline mixing enhanced diapycnal
386 nutrient fluxes, the overall productivity in the thermocline, and the associated carbon
387 export to the BBL.

388 The occurrence of near-inertial oscillations in shelf seas during the stratified
389 season has been reported in several studies from the North Sea (van Haren et al.,
390 1999; Knight et al., 2002) and in other shelf seas (e.g., Rippeth et al., 2002;
391 McKinnon and Gregg, 2005). During the presence of baroclinic inertial waves in the
392 water column, periods of enhanced shear have been observed in the western Irish Sea
393 (Rippeth et al., 2009), the Celtic Sea (Palmer et al., 2008) and the northern North Sea
394 (Burchard and Rippeth, 2009). These take the form of shear spikes, which occur
395 approximately every inertial period and in bursts lasting several days.

396 While we mainly attributed the observed enhanced turbulent mixing to the
397 occurrence of a near-inertial wave, the site's physical setting has further implications
398 for mixing processes in the thermocline. In the northern hemisphere, sites with anti-
399 cyclonic tides, such as Tommeliten, are often characterized by an increased vertical
400 extension of the BBL, and higher BBL dissipation rates than comparable cyclonic
401 sites (see Simpson and Tinker, 2009). As a result of this enhanced BBL thickness, we
402 observed sporadically elevated thermocline turbulence resulting from tidal-driven
403 bottom turbulence propagating vertically towards the thermocline (Fig. 5b). A study
404 by Burchard and Rippeth (2009) also reported that short lived thermocline shear
405 spikes can arise due to the alignment of the surface wind stress, bulk shear, and bed
406 stress vectors in the presence of baroclinic near-inertial motions and barotropic tidal
407 currents. These mechanisms are stronger with anti-cyclonic tides. Although all the
408 features required for shear spike generation were present during the observational
409 period, the two-layer mechanism described by these authors would require a more
410 complex water column structure to be applicable to the Tommeliten site.

411 The site's water column structure clearly showed the occurrence of a 10 m
412 thick transition layer (Fig. 2a). This layer represents the region of the water column
413 | where mixing turns from elevated in the SBL to strongly reduced in the interior layer
414 | (Ferrari and Boccaletti, 2004). The transition layer is therefore an obligate pathway
415 | for solute and heat exchange between SBL and the interior layer (Ferrari and
416 | Boccaletti, 2004; Rhein et al., 2010) and has been reported to be a region of enhanced
417 shear and near-inertial wave activity (Dohan and Davis, 2011). Although the
418 presented data did not allow quantification of the O₂ exchange across the transition
419 layer, such contribution might be considerable and thus highly relevant for the cycling

420 of O₂ and CO₂ in the upper water column, which in turn could have direct biological
421 implications.

422

423 **4.2 BBL O₂ dynamics**

424 Ultimately, observed O₂ depletion in the BBL of the central North Sea
425 depends on the supply of organic matter, the rate of carbon mineralization, and the
426 flux of O₂ to the bottom water either from horizontal advection or turbulent vertical
427 transport. Our study investigated the significance of turbulent vertical O₂ fluxes to the
428 BBL, which has been previously overlooked in shelf sea carbon balances. Studies
429 focusing on O₂ replenishment in the BBL through the thermocline are limited to
430 freshwater systems (e.g. Bouffard et al., 2013; Kreling et al., 2014). In a large
431 stratified water body such as Lake Erie, O₂ transport from the thermocline to the
432 hypolimnion was found to be substantial, with a magnitude comparable to ~18% of
433 the hypolimnetic O₂ utilization rate over the whole stratification period (Bouffard et
434 al., 2013).

435 Horizontal O₂ gradients and associated horizontal advective O₂ fluxes were
436 not quantified in this study. Our data suggest, however, that such fluxes would not
437 significantly contribute to the O₂ balance at the Tommeliten site. BBL O₂
438 concentration time series (Fig. 5a) did not show any variability at the tidal and or
439 inertial frequencies, implying that horizontal O₂ gradients were small. Additionally,
440 mean currents in the BBL were small (~2 cm s⁻¹) compared to the tidal amplitudes.
441 This, in conjunction with weak horizontal O₂ gradients, suggests that horizontal
442 advective O₂ fluxes during our observational period are negligible compared to the
443 turbulent O₂ flux from the thermocline.

444 Based on the above, we can argue that the O₂ dynamics during the stratified
445 period are more complicated than previously regarded. To maintain an excess of O₂ in
446 the thermocline, primary producers require adequate nutrient entrainment from the
447 bottom water to fuel potential new production. The resulting increase in (new)
448 productivity and subsequent export to the bottom water could boost the carbon
449 turnover estimates substantially. Using a 1:1 O₂ utilization – carbon re-mineralization
450 (see Canfield, 1993), Greenwood et al. (2010) inferred the average BBL carbon re-
451 mineralization rate at the nearby North Dogger to be 15 mmol m⁻² d⁻¹, or 180 mg C m⁻²
452 d⁻¹. Similar results for a typical NW European shelf sea were obtained via modeling

453 by Sharples (2008), who reported rates ranging from ~ 35 to ~ 200 $\text{mg m}^{-2} \text{d}^{-1}$ for neap
454 and spring tide, respectively. Their study, however, did not include the daily tidal
455 variation, and thus rates could be much higher on shorter timescales.

456 With the absence of targeted long-term studies focusing on O_2 and carbon
457 dynamics in the thermocline and BBL, we can only speculate on the long-term fate of
458 the BBL O_2 and its replenishment from the thermocline by vertical O_2 fluxes (F_θ). It
459 seems possible that the overall net BBL water column O_2 respiration, R , is higher than
460 previously thought, suggesting a much higher carbon turnover than inferred from the
461 apparent O_2 loss rate. Based on Eq. (6), the BBL carbon re-mineralization (and export
462 to the BBL) would be on the order of nearly 850 $\text{mg C m}^{-2} \text{d}^{-1}$, a factor of nearly 5
463 higher than reported by Greenwood et al. (2010). However, the same turbulent
464 transport that supports the O_2 export from the DCM to the BBL also supports BBL
465 nutrient import to the DCM (Fig. 6). The higher import of nutrients to the DCM likely
466 promotes additional primary production and a subsequent increase in organic matter
467 (OM) export to the BBL. In such a scenario, the O_2 flux to the BBL presented in this
468 “snapshot” study will be associated with additional OM to the BBL, and therefore
469 lead to a temporary increased re-mineralization that offsets the increased F_θ . While
470 the overall effect is an increase in carbon turnover, this process would not result in
471 any observable change in the decreasing O_2 trend (apparent O_2 loss rate).

472

473 4.3 Causes and controls on BBL O_2 depletion

474 According to Boers (2005), for BBL O_2 to decrease throughout the stratified
475 season, there must be suitable physical conditions, biomass production, nutrient input
476 and continued benthic O_2 uptake. SUR , and thus the sediment nutrient release and
477 organic carbon mineralization have been shown to be strongly tidal-driven (McGinnis
478 et al., 2014). Therefore, we briefly discuss the potential tidal impact driving the
479 overall carbon cycling and suggest factors that may promote the development of
480 lower BBL O_2 concentrations during the stratification period.

481 Tidal forcing on diapycnal constituent fluxes and primary production have
482 been explored by e.g., Sharples et al. (2007, 2008). The authors showed that spring-
483 neap tide drives nutrient fluxes between the BBL and the DCM at the thermocline,
484 and the carbon export. Based on our velocity measurements and estimated O_2 fluxes,
485 we can expect similar patterns corresponding to semidiurnal tidal fluctuations. Blauw

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488 et al. (2012) investigated fluctuating phytoplankton concentrations in relation to tidal
489 drivers and found that in the southern North Sea, chlorophyll fluctuations correlated
490 with the typical tidal current speed periods, the semidiurnal tidal cycle, in addition to
491 the day-night and spring-neap periods. During most of the year, chlorophyll and
492 suspended particulate matter fluctuated in phase with tidal current speed and indicated
493 alternating periods of sinking and vertical mixing of algae and suspended matter with
494 tidal cycles. Thus, these results suggest that in addition to the spring-neap tidal cycles,
495 we can expect a semidiurnal tidal-driven export of carbon and O₂ from the DCM to
496 the BBL, and entrainment of nutrients that strongly vary based on a timescale related
497 to the semi-diurnal tidal cycle.

498 The flux of O₂ from the DCM production zone downward to the BBL could
499 set the lower limit of the BBL O₂ concentration, and thus the O₂ depletion level,
500 during the stratification period. If there is little isolation between the zone of
501 production and the zone of mineralization, then the net O₂ production and O₂
502 utilization would nearly balance. In such case, the apparent O₂ loss in the BBL would
503 either be negligible or very small, depending whether the *SUR*, which is largely
504 particulate organic matter driven, will be balanced by the ventilation from the
505 thermocline. However, historically decreasing BBL O₂ concentrations within the
506 North Sea (Queste et al., 2013) point to an increasing disconnect between the primary
507 O₂ production zone and the mineralization zone. Greenwood et al. (2010) state that
508 stratification is an important factor which determines susceptibility to O₂ depletion,
509 especially in their nearby study site Oyster Grounds.

510 Surveys on the North Sea have shown that the regions with the lowest BBL O₂
511 concentrations are generally characterized by the strongest stratification (see Queste
512 et al., 2013), with the lowest values (~100 μmol kg⁻¹) reported to occur during
513 particularly calm and warm weather (see Boers, 2005; Weston et al., 2008). Strong
514 gradients in the thermocline associated with warmer temperature are suggested to
515 limit the O₂ flux to the BBL (Weston et al., 2008). This points to potential future O₂
516 depletion resulting from increasing temperatures leading to both stronger stratification
517 and a longer stratification season (Lowe et al., 2009). However, it could be argued
518 that if O₂ fluxes between the DCM and BBL were suppressed, then the upward
519 nutrient fluxes would be similarly suppressed, thus inhibiting primary production and
520 reducing the potential for O₂ deficits.

521

522 4.4 Biological perspective

523 The occurrence of stronger stratification and subsequently reduced turbulent
524 mixing could alter algal populations (Hickman et al., 2009), potentially favoring
525 migrating/swimming phytoplankton. An example of these migrating phytoplankton
526 species, armored dinoflagellates, are extensively found in the DCM of the central and
527 northern North Sea during the summer months; their abundance was found to be
528 largely determined by the local hydrodynamic conditions (Reid et al., 1990). In calm
529 conditions, which are typically associated with stronger stratification, there are often
530 blooms of migrating dinoflagellates which have access to the large nutrient pool in the
531 deeper water and can out-compete non-migrating species for both light and nutrients.
532 Stronger turbulent mixing, in contrast, has been suggested to interfere with their
533 swimming abilities (see Jephson et al., 2012 and references therein).

534 Due to the limited observational period of this study and our measurement
535 approach, we were not able to observe algal migration. However, based on the above
536 literature observations of algae in the central North Sea, we propose that algal
537 migration could be an over-looked mechanism that contributes to diminishing bottom
538 water O₂ levels. By promoting an upward shift of the DCM and moving the associated
539 O₂ production higher in the thermocline where turbulence levels are reduced, while
540 still maintaining comparable production rates, even a few meters, upward shift would
541 substantially reduce turbulent O₂ fluxes to the BBL, while maintaining similar rates of
542 carbon export (settling armored dinoflagellates). Studies on climate change impacts
543 on the North Sea have suggested that O₂ loss in the bottom waters would mainly
544 result from a strengthening of the stratification and O₂ solubility reduction with
545 increasingly warmer waters (e.g., Meire et al., 2013). In those scenarios, the intricate
546 interplay between local tidally-driven processes, water column structure,
547 biogeochemical cycling and active phytoplankton migration have not been considered
548 nor quantified. The proposed mechanism could contribute to the observed decreasing
549 O₂ levels in the North Sea water column, however, further detailed studies are
550 obviously necessary to test and quantify this effect, and the results described in this
551 study, at the seasonal level.

552

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576

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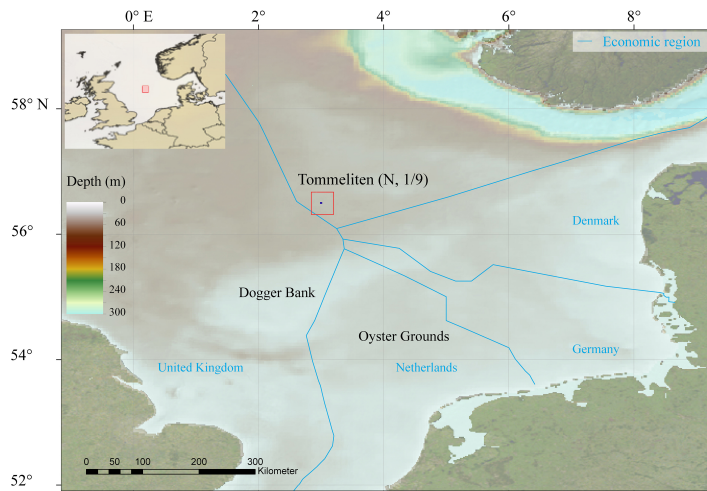
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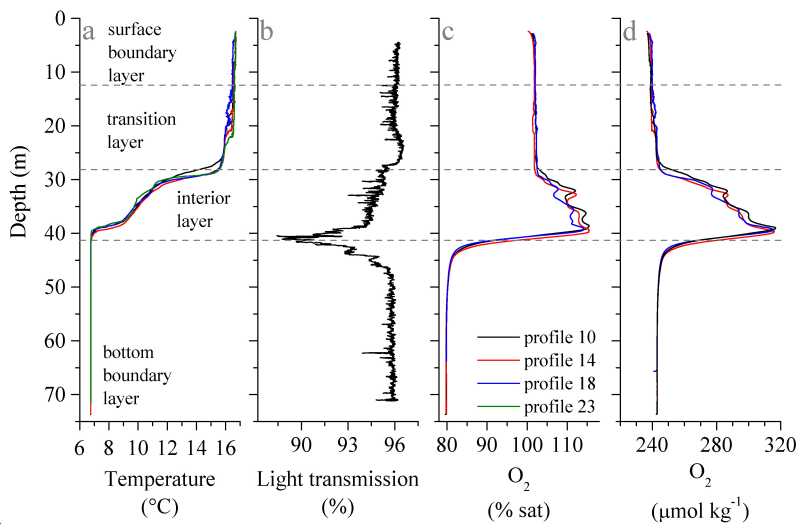
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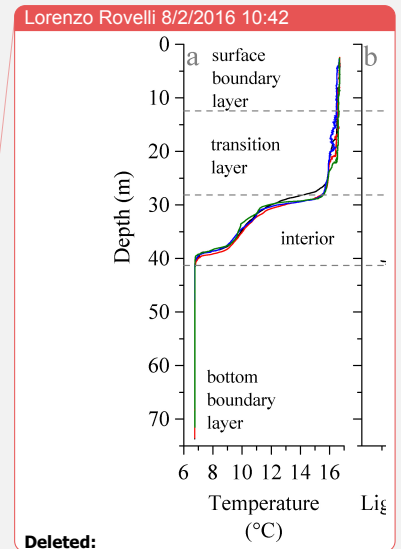
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786 Figure 1. Map of the North Sea indicating the water depths and location of the
787 Tommeliten site and the borders of the economic regions of the surrounding European
788 countries.

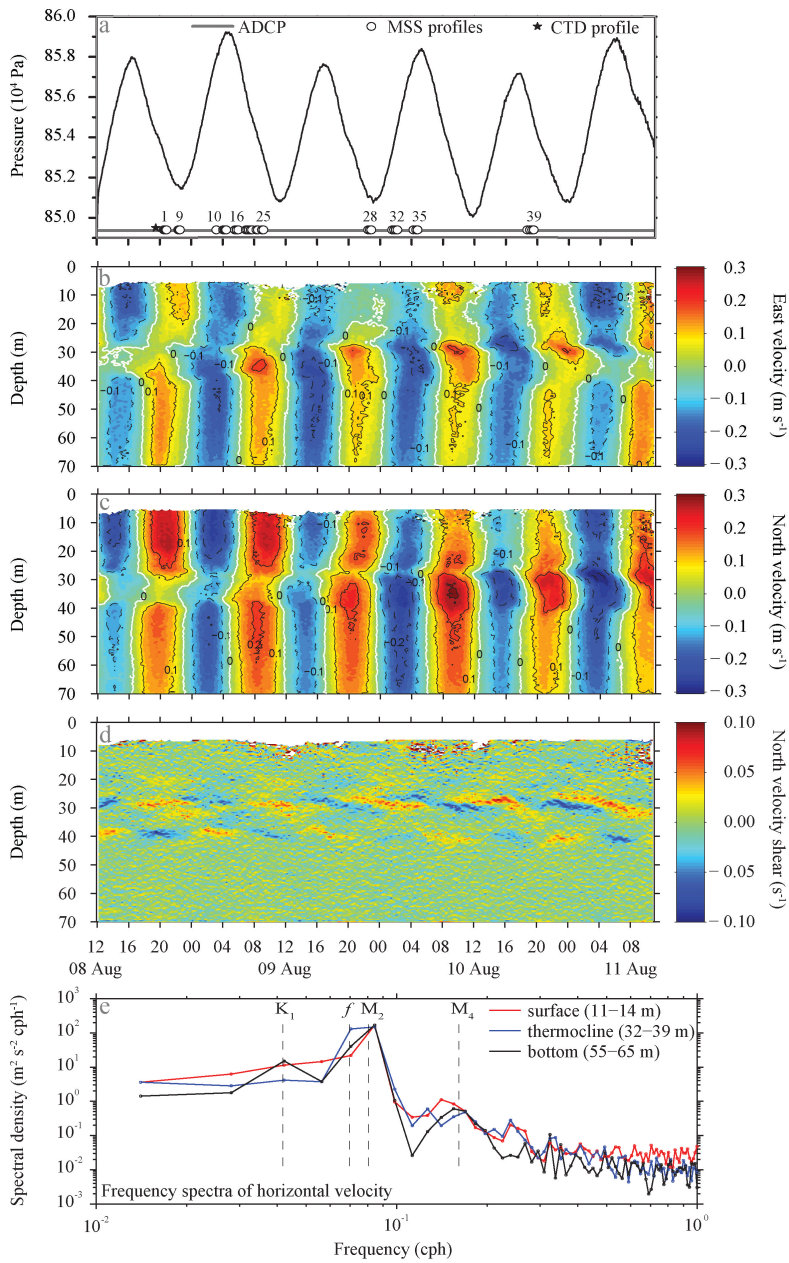


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791 Figure 2. Water column profiles, based on high-resolution MSS profiles (a, c, d) and
 792 ship CTD profile (b). (a) Potential temperature profiles. Water column layers were
 793 identified based on the temperature profiles. A 0.2°C and 1.5°C decrease from the
 794 surface boundary layer average temperature (3–6 m depth) was used to determine the
 795 depth of the surface boundary layer – transition layer interface and the transition layer
 796 – interior layer interface, respectively. Correspondingly, a 0.2°C from a 50–60 m
 797 depth average temperature was used to locate the interior layer – bottom boundary
 798 layer interface. (b) Light transmission profile. (c, d) O₂ saturation profiles and
 799 associated absolute concentrations.



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805 Figure 3. Current regime at the Tommeliten site from ADCP measurements (a - d)
 806 and spectral analysis (e). (a) Sea surface elevation relative to average level during the
 807 observational period (elevation = 0 m) and schedule of different instrument

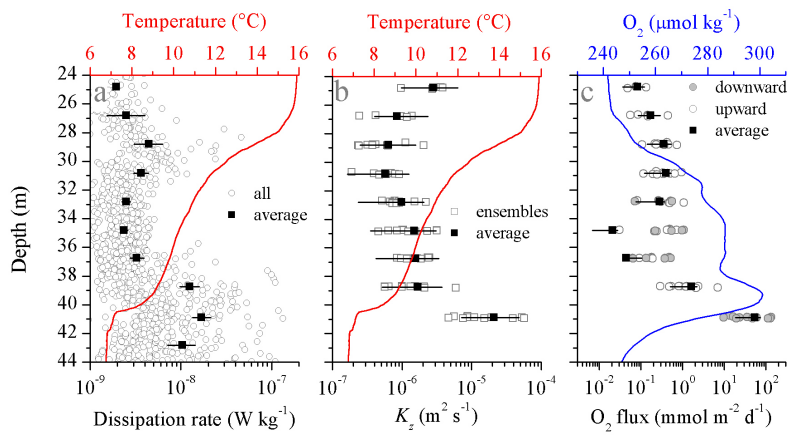
808 deployments. Numbers on the MSS markers indicate the profile number. (b, c)
809 Horizontal velocities, showing 20 min averaged east (b) and north (c) velocities. (d)
810 Vertical shear of north velocity, dv/dz , calculated from the ADCP velocity data (see
811 panels b, c). Note that panels a - d have the same time axis. (e) Frequency spectra of
812 horizontal velocity calculated from the ADCP data for selected depth ranges for the
813 SBL (surface; red line), thermocline (blue line), and BBL (bottom; black line). The
814 inertial frequency f , and the frequency for the diurnal declination tide, K_1 , semi-
815 diurnal lunar tide, M_2 and shallow water tidal constituent M_4 are marked.

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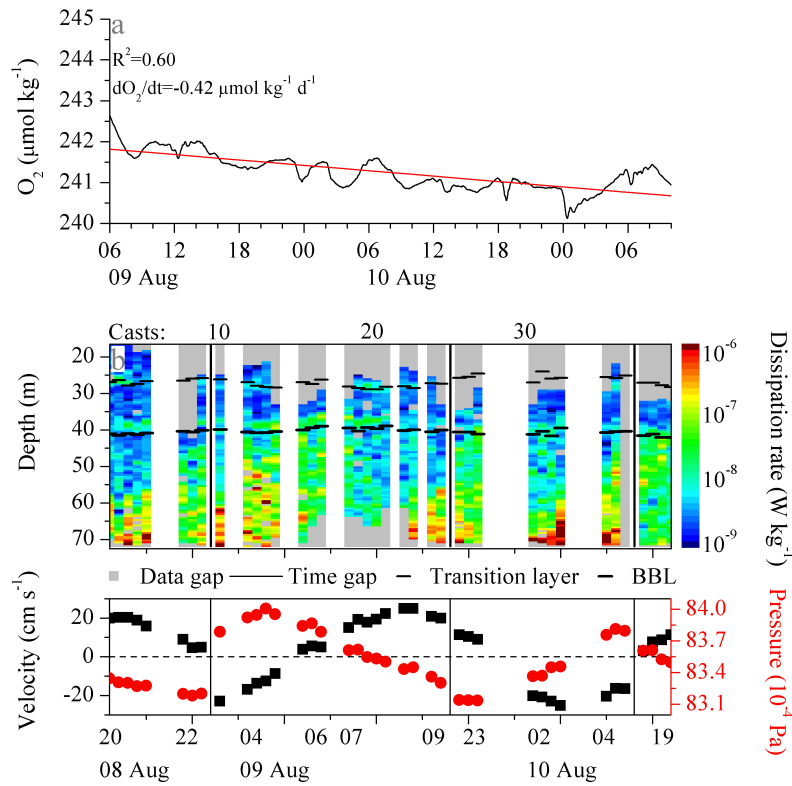


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820 | Figure 4. Overview of turbulent transport and O₂ fluxes within the interior layer
 821 | (defined in Fig. 2). Each panel is overlaid with temperature (a, b) and O₂
 822 | concentration (c) profiles. (a) Dissipation from all profiles (open dots) together with
 823 | the arithmetic mean (solid squares). (b) Average vertical eddy diffusion coefficient K_z
 824 | with uncertainties bars and the K_z values for every ensemble (open squares), which
 825 | represent the average over 3 to 4 consecutive profiles. (c) Calculated average O₂ flux
 826 | over 2 m bins with the respective uncertainties intervals (solid square and black line).
 827 | The values for each profile cluster are shown both downward and upward fluxes (grey
 828 | solid and open dots, respectively). Note that in the center interior layer (33 – 37 m)
 829 | the average reflects the combination of the variability of the observed upward and
 830 | downwards fluxes.

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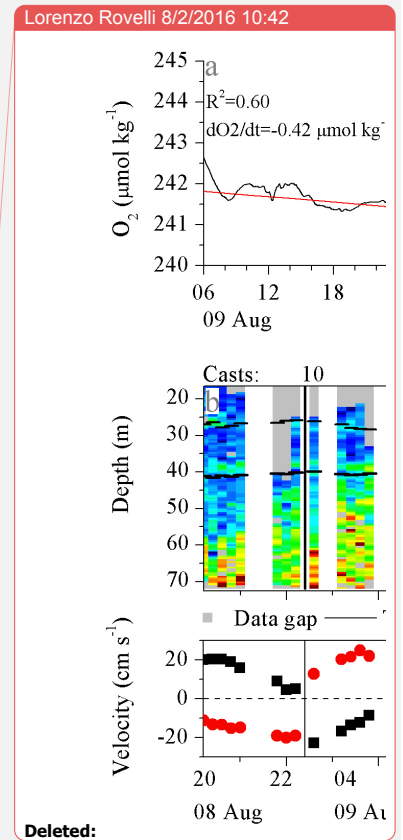


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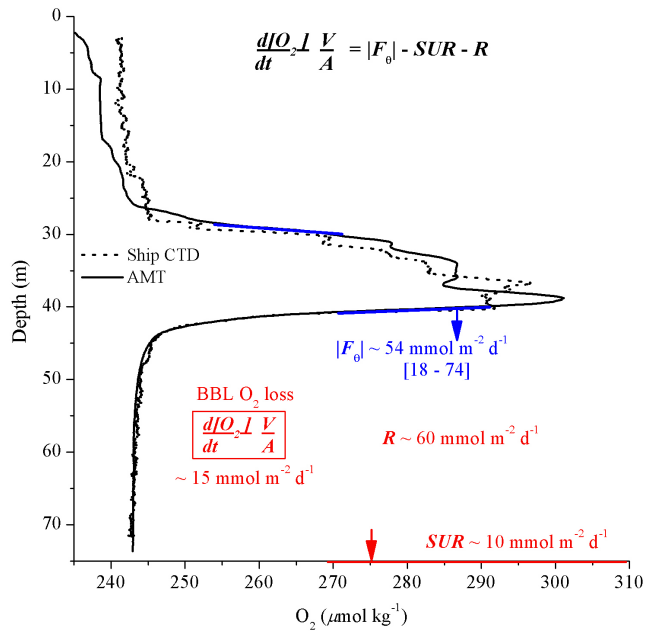
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834 Figure 5. BBL dissolved oxygen time series and turbulence contour. (a) Near-seafloor
 835 BBL O_2 concentration changes over the observational period from the POZ-Lander.
 836 Red line indicates the estimated apparent linear O_2 loss. (b, top) Turbulence contour
 837 plot of all MSS90 casts together with the temperature layers. Thin and thick dashed
 838 lines represent the transition layer – interior layer interface and the interior layer –
 839 BBL interface, respectively. Gray spots indicate data missing due to uncompleted
 840 profiles (casts 16-23), unsuccessful profiles (cast 36), or flagged as bad based on
 841 spikes, collisions and suspected contamination due to ship activity. The vertical black
 842 lines indicate the transition (time gaps) between consecutive profile ensembles. (b,
 843 bottom) Background information on bottom current, and hydrostatic pressure during
 844 the casts. Both velocity and pressure data were collected by the deployed POZ lander.
 845 Note that as a result of the time gaps between the consecutive MSS90 casts (see Fig.
 846 3a) the time scale is not linear.

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852 | Figure 6. Main O₂ fluxes in this study. The ranges shown for the interior layer O₂
 853 | fluxes refer to the associated uncertainty and intermittency levels. The sediment O₂
 854 | uptake rates (*SUR*) are based on eddy correlation (EC) measurements (McGinnis et
 855 | al., 2014), while central North Sea apparent BBL O₂ loss is based on Greenwood et al.
 856 | (2010) and this study. Representative O₂ profiles are based on the AMT sensor on the
 857 | MSS profiler (solid line) and ship CTD (dotted line). Note that while the O₂ profiles
 858 | showed differences in absolute concentration within the thermocline, the actual O₂
 859 | gradients within the thermocline-BBL oxycline are comparable.

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