

Maximum sinking velocities of suspended particulate matter in a coastal transition zone

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Abstract. Marine coastal ecosystem functioning is crucially linked to the transport and fate of suspended particulate matter (SPM). Transport of SPM is, amongst others, controlled by sinking velocity w_s . Since w_s of cohesive SPM aggregates varies significantly with size and composition of mineral and organic origin, w_s exhibits large spatial variability along gradients of turbulence, SPM concentration (SPMC) and SPM composition. In this study, we retrieved w_s for the German Bight, North Sea, by combining measured vertical turbidity profiles with simulation results for turbulent eddy diffusivity. We analyzed w_s with respect to modeled prevailing dissipation rates ϵ and found that mean w_s were significantly enhanced around $\log_{10}(\epsilon(\text{m}^2 \text{s}^{-3})) \approx -5.5$. This ϵ region is typically found at water depths of approximately 15 m to 20 m along cross-shore transects. Across this zone, SPMC declines drastically towards the offshore and a change in particle composition occurs. This characterizes a transition zone with potentially enhanced vertical fluxes. Our findings contribute to the conceptual understanding of nutrient cycling in the coastal region which is as follows: Previous studies identified an estuarine circulation. Its residual landward-oriented bottom currents are loaded with SPM, particularly within the transition zone. This retains and traps fine sediments and particulate-bound nutrients in coastal waters where organic components of SPM become re-mineralized. Residual surface currents transport dissolved nutrients off shore, where they are again consumed by phytoplankton. Algae excrete extracellular polymeric substances which are known to mediate mineral aggregation and thus sedimentation. This probably takes place particularly in the transition zone and completes the coastal nutrient cycle. The efficiency of the transition zone for retention is thus suggested as an important mechanism that underlies the often observed nutrient gradients towards the coast.

1 Introduction

Biogeochemical cycling and functioning of marine coastal and shelf sea systems crucially relies on particle transport. Vertical fluxes of suspended particulate matter (SPM) are determined by sinking velocity w_s and indirectly affect the horizontal transport. In coastal systems, SPM is composed of living and non-living particulate organic matter (POM) and fine cohesive and non-cohesive resuspended minerals. Fine-grained minerals of sizes typically up to $8 \mu\text{m}$ (Chang et al., 2006) and POM can undergo aggregation and fragmentation processes that change sinking velocity and thus transport properties. As a consequence of flocculation, SPM aggregates ubiquitously possess a broad spectrum of size and composition (Fettweis, 2008). This

heterogeneity between flocs increases the methodological effort to analyze w_s *in situ* (Fettweis, 2008). On larger scales, SPM concentration (SPMC) and composition additionally exhibit strong spatio-temporal variability that are the result of manifold interplaying processes. Tidal and wind-induced currents are the major driver for resuspension and subsequent horizontal transport, while biological processes, such as algae growth and bio-induced sediment stabilization (Black et al., 2002; Stal, 2003), interfere and thus shape the complex distribution of SPM in coastal and estuarine systems. Typically, SPMC and composition show cross-shore gradients (Tian et al., 2009; van der Lee et al., 2009; Li et al., 2010). In shallow waters near the coast, where turbulence and thus resuspension is high, SPMC consists of flocs that are mainly composed of mineral particles with high densities. By contrast, in deeper off-shore regions, SPMC is lower and the flocs are looser and more organic. This general pattern of changing SPMC and composition from coast to open waters is typical for our research area, the German Bight (Eisma and Kalf, 1987), and is observed worldwide in estuaries (e. g. Fugate and Friedrichs, 2003) and across coastal seas (e. g. van der Lee et al., 2009). Since both SPMC and turbulence control w_s of cohesive material (Pejrup and Mikkelsen, 2010), it is likely that these cross-shore SPM gradients induce considerable spatial variability in w_s and thus affect transport and fate of SPM in coastal marine systems. However, to date and to the authors' knowledge, no comprehensive analysis has addressed system-wide cross-shore gradients in sinking velocity, especially in relation to possible drivers.

Sinking velocity of SPM is determined by floc size and density, both resulting from a complex interplay of processes. Particle size distribution is locally governed by restructuring (Becker et al., 2009), aggregation and processes mainly driven by turbulence-induced shear, $\overline{G} = (\epsilon/\nu)^{1/2}$, (Pejrup and Mikkelsen, 2010). This turbulence is generated by energy dissipation ϵ (Camp and Stein, 1943) with kinematic viscosity ν . Under a given shear regime, aggregation is controlled by floc volume concentration and adhesion properties of the primary particles of mineral and organic origin that form the floc. The higher the volume concentration of flocs, the higher is the encounter rate. Subsequently, adhesion forces of the involved particles eventually determine the probability to stick together. In addition, adhesive forces limit the intrusion of particles during clustering and loosens the floc structure (Meakin and Jullien, 1988). This leads to a floc morphology that shows self-similar fractal scaling (Kranenburg, 1994). As a result, aggregates possess decreased density compared to the comprising primary particles. In addition, adhesion forces between the particles within a floc strengthen the resistance of aggregates against fragmentation (Kranenburg, 1999), while the smallest eddies, with sizes of the Kolmogorov microscale (Kolmogorov, 1941), potentially limit the maximum particle size (Berhane et al., 1997). In sum, w_s is locally governed by turbulence-driven processes whose rates depend on the physico-chemical properties and volume concentration of particles.

To date, there is still a lack of understanding of how environmental conditions and especially biological processes affect physico-chemical properties, and thus w_s of cohesive SPM. Typically, a power law relation between median w_s and SPMC is postulated and found at local measurements (Dyer, 1989). However, these relations vary considerably in their power factor among different systems as summarized by Dyer (1989) for various estuaries. This variation can be attributed to different shear stresses and physico-chemical properties of involved particles. They are particularly subjected to algae and microbial extracellular polymeric substances (EPS) excretions such as transparent exopolymer particles (TEP) that are known to mediate aggregation processes. TEP bridge and glue mineral particles together (Decho, 1990) and potentially increase resistance against particle fragmentation (Fettweis et al., 2014). By these mechanisms, TEP are hypothesized to enable phytoplankton to clear

the water column from suspended sediments (Fettweis et al., 2014). Such clearing ability may contribute to spatial variability of w_s which would affect biogeochemical cycles. So far, knowledge on spatial variations of w_s on a system scale is rare. A study on a transect in San Francisco Bay was carried out (Manning and Schoellhamer, 2013), but generally, there is no systematic understanding about the relevance of sedimentation variability for biogeochemical cycles in coastal ecosystems.

5 Such knowledge, though, is needed to understand and model coastal shelf biogeochemistry and sedimentology. For example, it is still an ongoing scientific discussion which processes and to what extent do they sustain the net-sediment transport towards the Dutch, German and Danish coast into the Wadden Sea (Postma, 1984). We therefore aim at a reconstruction and analyses of w_s of SPM for the German Bight, North Sea. We develop a new approach to retrieve w_s from high resolution turbidity profiles in combination with vertical mixing rates estimated by a hydrodynamical model. Our findings are particularly discussed in

10 light of their relevance for biogeochemical cycling of matter in coastal waters.

2 Methods

In the following section, the study area, sampling and preprocessing of observational data are described. We further explain the procedure for extracting sinking velocities from observations with the aid of hydrodynamical model results.

2.1 Study area

15 The surveyed German Bight (Fig. 1) is located in the south-east of the North Sea and features a typical depth of about 30 m to a maximum of 50 m. The North Sea is a shallow shelf sea with an average depth of 80 m (Sündermann and Pohlmann, 2011) and is connected to the North Atlantic via the English Channel in the south-west and opens towards the north, spanning the European continental shelf. The North Sea is exposed to tides whose tidal range is between approximately 1.8 m and 3.4 m in the southeastern German Bight. The main North Sea tidal wave turns anti-clockwise and drives the large scale current system

20 (Sündermann and Pohlmann, 2011). SPM originating from the British coast is transported eastwards by the East Anglian Plume and occasionally reaches the surveyed area (Fettweis et al., 2012; Pietrzak et al., 2011). In the southern North Sea, SPM and dissolved nutrients that enters though the river Rhine are transported eastward along the Dutch and German East-Frisian shore by this large-scale current system. In addition, the German rivers Ems, Weser and Elbe discharge into the German Bight. As a result of the riverine nutrient input, the German Bight features a generally high primary productivity (Joint and Pomroy,

25 1993) while it possesses pronounced heterogeneity with strong cross-shore gradients in i) nutrients (Brockmann et al., 1990; Ebenhöf et al., 2004), and ii) SPMC and composition (Eisma and Kalf, 1987). The near-coastal waters of the German Bight are characterized by relatively high SPMC of few to several hundred gram per cubic meter in near-bottom regions. SPMC possesses a pronounced seasonality showing higher values in winter than in summer. This seasonal pattern is reflected in the shallow parts of the tidal back barriers of the Wadden Sea by the bottom fraction of fine, cohesive sediments accumulating

30 during spring and summer due to aggregation and subsequent increased sedimentation (Chang et al., 2006). The sediment catchment area for the Dutch and East-Frisian Wadden Sea system, an area sheltered from the North Sea by a chain of islands, is hypothesized to be defined by a *line-of-no-return* located alongshore off the coast (Postma, 1984). The *line-of-no-return*

is conceptually described as an imaginary line beyond which particles escape coastal trapping mechanisms such as density gradient-driven undercurrents (Postma, 1984).

2.2 Sampling and processing of observational data

Measurements were carried out as part of the 'Coastal Observing System for Northern and Arctic Seas' project (COSYNA: www.cosyna.de; see also Baschek et al., in review, 2016). Sensors were mounted on board a towed vehicle (ScanFish, Mark II, EIVA a/s, Denmark). The ScanFish was remotely operated behind a vessel sailing with a speed of 6 to 8 knots. Cruises were carried out during May, July/August, September (2009), March, May, July, September (2010) and April, June and September (2011) during moderate weather conditions. No winter measurements were carried out. The transect grid covered the German Exclusive Economic Zone (Fig. 1). The ScanFish was forced to operate in nearly V-curved undulating path mode between approximately 3m below the water surface and sea bottom, with a vertical speed of 0.4 m s^{-1} and a sampling rate of 11 Hz, yielding a vertical data spacing of about 0.04m. Our analysis considered measurements of specific conductance (Conductivity Sensor, ADM Elektronik, Germany), water temperature T (PT100, ADM Elektronik, Germany), pressure p (PA-7, Keller AG, Switzerland), optical turbidity (Seapoint Turbidity Meter 880nm, Seapoint Sensors inc., USA), and chlorophyll a fluorescence F (TriOS MicroFlu-chl, TRIOS Inc., Germany). Specific conductivity and T were calibrated against reference standards directly before and after the ship surveys. Potential water density ρ_θ , expressed as $\sigma_T = \rho_\theta - 1000 \text{ kg m}^{-3}$, was calculated according to the EOS-80 equations of state (Fofonoff and Millard, 1983). Thermal lag effects were visible in σ_T near the thermoclines, but this effect was neglected since the selection criteria for profiles diminished their relevance for this analysis, as described below. During post-processing, data generally underwent tests for any stuck values and spikes, and were finally visually inspected to remove remaining faults. Turbidity, measured in formazine turbidity unit (FTU), was converted to SPMC using a factor of $1.08 \text{ g (dry-weight) m}^{-3} \text{ FTU}^{-1}$ that was determined by linear regression ($r^2=0.967$; Röttgers et al., 2011). Laboratory investigations revealed a sensitivity of the fluorescence signal to turbidity. Therefore, a correction factor was applied to measured fluorescence. A factor of $(0.32 + (1 - 0.32) \exp(-0.025 \text{ turbidity}))^{-1}$ was experimentally determined with commercially available chlorophyll a dissolved in varying formazine concentrations. As a final step, data sets were split into up- and down casts between consecutive vertical extrema of the undulating flight path. Casts with deficient pressure data were discarded.

2.3 Data processing and sinking velocity extraction

Sinking velocities were obtained by fitting an analytical solution of the vertical distribution of SPMC to observations. Assuming steady state and neglecting horizontal advection, the SPM transport equation in time t for concentration C reduces to

$$\frac{\partial}{\partial t} C = \frac{\partial}{\partial z} k_v \frac{\partial}{\partial z} C - \frac{\partial}{\partial z} (w_s C) = 0 \quad (1)$$

This describes the balance between fluxes in the positively downward pointing vertical direction z due to sinking and turbulent eddy diffusivity k_v . If we assume that the sinking time scale is larger than the tidal period, we can simplify Eq. (1) by using the vertically averaged k_v of a profile. This assumes a rather homogeneous turbulence intensity which we account for in the

data processing described below. If we further allow fluxes across profiles borders, which cancel out under steady state, we can derive an analytical model ($C_m(z)$) for depth-dependent SPMC

$$C_m(z) = \frac{\lambda \exp(\lambda z^*/H_p)}{\exp(\lambda) - 1} \langle C \rangle \quad \text{where } \lambda = \frac{w_s H_p}{\langle k_v \rangle} . \quad (2)$$

Here, H_p represents the profile height and $z^* = z - z_1$, where z_1 is the depth at which the profile starts. Application of the analytical solution to observed SPMC profiles requires information on corresponding k_v values. These and energy dissipation rate ϵ and water density σ_T were obtained from hydrodynamical simulations of Gräwe et al. (2015), based on a 1 nautical mile numerical model of the North Sea and the Baltic Sea. In the vertical, 42 terrain-following levels were used. Vertical mixing was parametrized by means of a two-equation $k - \epsilon$ turbulence model coupled to an algebraic second-moment closure (Canuto et al., 2001). The implementation of the turbulence module was done via the General Ocean Turbulence Model (GOTM; Umlauf and Burchard, 2005). The hydrodynamic parameters were stored as 2-hourly snapshots.

Hydrodynamic model results for k_v , ϵ , and σ_T were linearly interpolated in space and time to extract a corresponding value for every measured data point. Even though state-of-the-art hydrodynamical models generally perform very well, they may locally exhibit discrepancies with observations. To discriminate the lack of congruency, both observed and modeled σ_T were interpolated on a common vertical grid of $\Delta z = 0.05$ m and filtered by applying a least-square straight line fit to the data and a 'natural' cubic spline interpolant. The latter had a weight of 80 % to produce a smooth interpolation curve $\sigma'_T(z')$. Subtracting the respective vertical mean resulted in profiles $\widetilde{\sigma}_T$ for both observed and modeled data. We further constrained our analysis to casts that satisfy two criteria. First, after subtraction of modeled $\widetilde{\sigma}_T$ from observed data, the standard deviation should not exceed $\text{std}(\Delta \widetilde{\sigma}_T) \leq 0.015 \text{ kg m}^{-3}$. Second, the density gradients defined as $\delta_z \sigma_T = (\langle \sigma'_T(\text{last meter}) \rangle - \langle \sigma'_T(\text{first meter}) \rangle) / H$, where $\langle x \rangle$ represents the mean of a variable x and specifically in this case, the vertical mean, should not exceed a difference of $\Delta \delta_z \sigma_T \leq 0.015 \text{ kg m}^{-3} \text{ m}^{-1}$ between observed and modeled data. Both limits for $\text{std}(\Delta \widetilde{\sigma}_T)$ and $\Delta \delta_z \sigma_T$ were applied to select similar vertically structured observed and modeled density profiles. The chosen values, however, were somewhat arbitrary and therefore considered in a Monte-Carlo-type sensitivity simulation (see below).

ScanFish cruises covered well-mixed coastal but partly also stratified waters in the inner German Bight. A consistent approach to retrieve w_s was required to meet both conditions. When needed, we chose to split casts into subprofiles and to fit the single analytical solution, Eq. (2), to (sub-) profiles. If observed $\delta_z \sigma_T < 5 \cdot 10^{-4} \text{ kg m}^{-3} \text{ m}^{-1}$, the whole profile was used for fitting, otherwise the cast was split. Since the critical value was set by visual inspection, sensitivity of this critical value was also considered in the Monte-Carlo-type simulation described below. Strong gradients in σ_T are typically considered to indicate dampening of turbulent mixing. The squared buoyancy frequency (N^2)

$$N^2 = \frac{g}{\rho} \frac{\partial \rho}{\partial z} , \quad (3)$$

where g denotes the gravitational acceleration constant, is related to the vertical eddy diffusivity k_v by (Osborn and Cox, 1972)

$$k_v = c \frac{\epsilon}{N^2} \quad (4)$$

with the current standard value for $c = 0.2$ (Lindborg and Brethouwer, 2008). A strong vertical sea water density gradient thus reflects low turbulent diffusion, which would imply to split profiles in stratified waters at the maximum density gradient(s). However, as recently discussed by Franks (2014), mixed layer depths defined by density gradients are a rather inadequate proxy for turbulence intensity. Since particle properties such as size and density are a function of shear rate and SPM components, it is reasonable to expect them to differ with vertical turbulence intensity, thus leading to different sinking velocities. For example, Leipe et al. (2000) reported vertically varying mean aggregate sizes for the Baltic Sea. While we generally expect vertical gradients in SPMC, strong vertical gradients in SPMC potentially reflect weak mixing and are thus possible indication for changes in turbulence intensities. The co-occurrence of strong gradients in σ_T and SPMC can thus signal changes in turbulence intensity and particle property. Splitting at these points allowed us to apply the analytical SPMC model (Eq. (2)), where we assumed a rather homogeneous turbulent diffusion by using the vertically averaged k_v . Thus, end points of subprofiles were set, where regions in the profile with

$$\left| \frac{1}{C'} \frac{\partial C'}{\partial z'} \cdot \frac{1}{\sigma'_T} \frac{\partial \sigma'_T}{\partial z} \right| > \left\langle \left| \frac{1}{C'} \frac{\partial C'}{\partial z'} \right| \right\rangle \left\langle \left| \frac{1}{\sigma'_T} \frac{\partial \sigma'_T}{\partial z} \right| \right\rangle \quad (5)$$

occur and reach their maximum. Here, $\langle x \rangle$ denotes the vertical average and C' is the SPMC smoothing spline-interpolated analogously to σ_T . Start points of the subprofiles were either the first data point from surface or where regions defined by Eq. (5) end with increasing depth. Only subprofiles longer than 4 m were considered for further analysis.

The analytical model, Eq. (2), was fitted to the original observation of each (sub-) profile with N_p data points, and the cost function

$$\text{err} = \frac{1}{N_p} \sum_{i=1}^{N_c} \frac{(C(z_i) - C_m(z_i))^2}{0.5(\delta_R C^2 + \delta_P C^2)} \quad (6)$$

was calculated accordingly. $\delta_P C^2$ and $\delta_R C^2$ represent the variance in concentration of a profile and in a region, respectively. The latter was defined as the variance for measurements around the profile within ± 5 min. to account for higher variability in coastal than in open water regions found in a pre-analysis of the data. Only profiles below a cost function value of 0.05, chosen by visual inspection, were considered for the analysis. The splitting and subsequent fitting of the analytical model is visualized in Fig. 2. The variables SPMC and fluorescence-to-SPMC ratio (F/SPMC) were vertically averaged for the respective (sub-) profiles. Afterward, variables were binned with respect to modeled ϵ and eventually averaged bin-wise, resulting in respective bin-wise ensemble means $\langle \text{SPMC} \rangle$, $\langle \text{F/SPMC} \rangle$, $\langle w_s \rangle$, and $\langle \epsilon \rangle$. To test for significant changes of binned $\langle w_s \rangle$ with $\langle \epsilon \rangle$, we applied a Mann-Whitney-U test with a significance level of $p < 0.05$ for each binned $\langle w_s \rangle$ against each other. In summary, approximately 67 % of initially measured ≈ 68000 profiles passed the congruency check with modeled ones. After application of the cost function threshold, this resulted in about 12260 w_s -values.

All applied threshold values were carefully selected by visual inspection during each filtering step. To quantify their influence on the result, we performed a Monte-Carlo-type simulation with a variation by $\pm 50\%$ for the following parameters: $\text{std}(\Delta \widetilde{\sigma_T})$, $\Delta \delta_z \sigma_T$, $\delta_z \sigma_T$, and the cost function error. The analysis was repeated for variations of all parameters against each other. For each parameter variation, binned mean values for $\langle w_s \rangle$ versus ϵ were calculated, eventually averaged and their standard deviation σ quantified (see Fig. 4 C).

2.4 Conceptual cross-coastal sinking velocity model

We introduce and apply a conceptual model for the cross-shore variation of w_s as a tool to interpret our results. According to Stokes (1851), w_s of a particle of diameter D can be described by

$$w_s = \frac{1}{18\mu} (\rho_f - \rho) g D^2, \quad (7)$$

- 5 where μ is the dynamic viscosity, ρ is the water density, and g the gravitational acceleration constant. The floc density ρ_f is strongly related to the particles' composition and structure. The latter can be described as fractal dimension d_f according to Kranenburg (1994), who derived the excess floc density $\Delta\rho_f = \rho_f - \rho$ for a particle composed of primary particles of diameter D_p and density ρ_p :

$$\Delta\rho_f = (\rho_p - \rho) \left(\frac{D_p}{D} \right)^{3-d_f}. \quad (8)$$

- 10 Following the approach of Kranenburg (1994) and assuming equally sized primary particles of diameter D_p , Maggi (2009) derived the excess floc density for an aggregate composed of mineral and organic particles with density ρ_s and ρ_o , respectively,

$$\Delta\rho_f = (\omega\Delta\rho_s + (1-\omega)\Delta\rho_o) \cdot \left(\frac{D_p}{D} \right)^{3-d_f}, \quad (9)$$

For w_s of an aggregate, it follows

$$15 \quad w_s = \frac{1}{18\mu} (\omega\Delta\rho_s + (1-\omega)\Delta\rho_o) g D_p^{3-d_f} D^{d_f-1}, \quad (10)$$

where $\omega = n_s/(n_s + n_o)$ and n_s and n_o are the number of mineral and organic primary particles, respectively.

- This simple w_s model, Eq. (10), allows for calculation of trends in w_s on a cross-shore transect. Based on observations, we presume the general following parameter changes from coastal to open waters: *i*) average primary particle size and aggregate size increase, *ii*) particles become more organic, and *iii*) more porous, thus d_f decreases. Based on a first order approach,
- 20 we assume a linear decrease or increase of parameters with distance from coast. For boundary conditions in coastal waters, we therefore assume an average $\langle D_p \rangle = 4\mu\text{m}$ due to the dominance of cohesive sediment (Winterwerp, 1998), $\langle D \rangle = 100\mu\text{m}$ (Fettweis et al., 2006), $\omega = 0.9$ to account for a loss on ignition (LoI) of ≈ 0.04 (Fettweis, 2008), and $d_f = 2.0$ as an average coastal fractal dimension (Winterwerp, 1998). By contrast, for open waters we presume $\langle D_p \rangle = 10\mu\text{m}$ for algae-dominated particles, $\langle D \rangle = 500\mu\text{m}$ since we assume similar trends as in van der Lee et al. (2009) for the Irish Sea, $\omega = 0.05$, to parallel
- 25 a LoI ≈ 0.89 (Eisma and Kalf, 1987), and $d_f = 1.6$ as a typical value for fluffy aggregates (Logan and Alldredge, 1989; Li and Logan, 1995). The densities of primary particles are set to $\rho_p = 2650\text{kg m}^{-3}$ (Maggi, 2009) and $\rho_o = 1100\text{kg m}^{-3}$, which is computed as an average of the range given in Fettweis (2008). The water density and the dynamic viscosity are set to $\rho = 1000\text{kg m}^{-3}$ and $\mu = 0.001\text{kg m}^{-1}\text{s}^{-1}$, respectively. The sensitivity of w_s to changes in each parameter is assessed by varying each parameter separately while keeping the other parameter parameters at their typical values for coastal waters, e. g.
- 30 for an organic-rich aggregate with an assumed diameter of $D = 100\mu\text{m}$, and $\omega = 0.5$, $D_p = 4\mu\text{m}$, and $d_f = 2$.

3 Results

3.1 Cross-coastal gradients

Energy dissipation rates ϵ generally possess spatio-temporal variability. Most relevant for this study, vertically and temporally averaged model-derived ϵ exhibit a cross-shore gradient from the inner German Bight towards the coast. The temporal averaging is carried out for the time of the cruises while accounting for the length of the tidal cycle. Values range between $\epsilon \approx 10^{-9} \text{ m}^2 \text{ s}^{-3}$ and $\epsilon \approx 10^{-4} \text{ m}^2 \text{ s}^{-3}$, respectively (Fig. 3).

Along this range, mean SPMC first increases from values below 1 g m^{-3} to approximately 10 g m^{-3} . Above $\log_{10}(\epsilon) = -5.5$, mean SPMC moderately decreases to approximately 6 g m^{-3} and increases again to 10 g m^{-3} with further increase of ϵ (Fig. 4 A) towards the coast.

Under the assumption that fluorescence can be applied as a proxy for POM, the potential significance of algae and their products for particle composition is depicted by the mean ratio between measured fluorescence and SPMC. High F/SPMC ratios indicate rather organic-rich particles, while low F/SPMC ratios indicate rather mineral particle-dominated flocs. The F/SPMC ratio shows rather high variability for $\log_{10}(\epsilon) < -6$ (Fig. 4 B). With a further increase in ϵ , the ratio drops and levels at approximately a fourth of the open German Bight value for $\log_{10}(\epsilon) > -5$.

Average sinking velocities $\langle w_s \rangle$ show very low values of order 10^{-6} m s^{-1} to 10^{-5} m s^{-1} in regions of $\log_{10}(\epsilon) < -7.5$ (Fig. 4 C). At higher ϵ , $\langle w_s \rangle$ increases significantly reaching a maximum of about $7 \cdot 10^{-4} \text{ m s}^{-1}$ around $\log_{10}(\epsilon) \approx -5.5$. This maximum $\langle w_s \rangle$ around $\log_{10}(\epsilon) \approx -5.5$ coincides with an increase of SPMC, and a drop in fluorescence to SPMC ratio. By further increasing ϵ , $\langle w_s \rangle$ decreases again to values of $\langle w_s \rangle \approx 3 \cdot 10^{-4} \text{ m s}^{-1}$. The Monte-Carlo-type simulation to estimate parameter uncertainties exhibits the same pattern, expressed as σ and 2σ confidence levels, 68 % and 95 % respectively (Fig. 4C), around their ensemble mean, and thus underpins the results of $\langle w_s \rangle$ along ϵ .

3.2 Conceptual model

Sinking velocity of an average particle (as parametrized in Sec. 2.4) changes almost linearly with variations of single parameters from open to coastal waters (Fig. 5, upper four panels). Flocs would sink less rapid because of the decreasing diameter of the aggregate or primary particles in the coastal region. By contrast, increasing amounts of mineral particles and fractal dimension would both lead to increased w_s . As an overall effect, when considering all parameters, w_s reaches a maximum in between the near-shore and open waters (lowest panel of Fig. 5).

4 Discussion

4.1 Sinking velocity on a gradient of prevailing energy dissipation rate

Sinking velocities were determined along a cross-shore transect in the German Bight defined by the prevailing ϵ . For the reconstruction of $\langle w_s \rangle$, we had to assume congruency between *in situ* measurements and hydrodynamic model results within

a range defined by the applied filters. The estimated $\langle w_s \rangle$ of order 10^{-6} m s^{-1} to $5 \cdot 10^{-4} \text{ m s}^{-1}$ in low and high energy dissipation regions, respectively (Fig. 4), compare very well with former *in situ* studies in the German Bight and adjacent areas. Puls et al. (1995) found $w_s \approx 10^{-6} \text{ m s}^{-1}$ to $2 \cdot 10^{-5} \text{ m s}^{-1}$ for the open German Bight, a low ϵ region, and Pejrup and Mikkelsen (2010) reported median w_s of $1 \cdot 10^{-4} \text{ m s}^{-1}$ to $11 \cdot 10^{-4} \text{ m s}^{-1}$ for the Rømø and Høyer Dyb, Danish Wadden Sea where ϵ is comparably high. Our estimated $\langle w_s \rangle$ are also well in the range found in other regions, e. g. in Chesapeake Bay (Gibbs, 1985) or at the Belgian Coast (Fettweis and Baeye, 2015). In agreement with former studies (e. g. compiled in Dyer, 1989), $\langle w_s \rangle$ increases with increasing $\langle \text{SPMC} \rangle$ (Fig. 6), which gives further confidence in the methodological approach. However, the correlation is rather poor compared to former local studies and can be explained by the heterogeneity of the German Bight system, seasonal effects and the intrinsically high variability of sinking velocities (Fettweis, 2008) and of turbulence (Pejrup and Mikkelsen, 2010). The significant maximum of $\langle w_s \rangle$ at $\langle \epsilon \rangle \approx 10^{-5.5} \text{ m}^2 \text{ s}^{-3}$ can be explained by the balance between aggregation and fragmentation. Both processes are controlled by turbulent shear that is generated by energy dissipation, SPM volume concentration, and adhesion properties of particles involved. Previous flocculation modeling studies point to a formation of w_s maximum along ϵ (Winterwerp, 1998; Pejrup and Mikkelsen, 2010). Our findings of a maximum $\langle w_s \rangle$ between $\epsilon \approx 10^{-6} \text{ m}^2 \text{ s}^{-3}$ and $10^{-5} \text{ m}^2 \text{ s}^{-3}$, which translates to shear rates of about $1 \text{ s}^{-1} - 3.2 \text{ s}^{-1}$ for $\nu = 10^{-6} \text{ m}^2 \text{ s}^{-1}$, should thus be comparable to former theoretical studies. For example, in a field work-based modeling approach Pejrup and Mikkelsen (2010) found the maximum of sinking velocity to occur under higher shear rates of about 8.5 s^{-1} . By contrast, Maerz et al. (2011) calculated highest mean sinking velocities of about $5 \cdot 10^{-4} \text{ m s}^{-1}$, similar to our findings, for a jar test device-based laboratory experiment with natural SPM during shear rates of about 1 s^{-1} . Given the underlying uncertainties in calculating the shear rate in both cases, our study agrees well with former theoretical studies. Hence, as suggested by Dyer (1989) and Pejrup and Mikkelsen (2010), in addition to SPMC, turbulent shear can be regarded as the major determinant for w_s .

4.2 Sinking velocity as a result of SPMC, composition and a turbulence gradient

Our study additionally suggests a change in particle composition with ϵ as proxied by the fluorescence-to-SPMC ratio. With ϵ , SPMC increases towards the coast and mineral fraction becomes dominant compared to organic particles at low ϵ (Eisma and Kalf, 1987). Hence, density and other physico-chemical properties of primary particles such as size and adhesion can change accordingly and thus influence w_s along a cross-shore transect. Our conceptual model illustrates the effect of varying particle properties on w_s independent of the observations used in our data analysis. However, the course of the parameters implicitly presumes certain environmental conditions, such as particle size being, among others, a function of SPMC and shear. Single parameter changes in the conceptual model result in mostly linear responses in w_s , while considering all parameters lead to a maximum in w_s on the conceptual cross-shore transect, in alignment with our findings in the data analysis. The assumption of linearly changing parameters is a first order approximation. Different and non-linear changes of parameters would lead to a transformation of the maximum sinking velocity zone such as a shifting and / or stretching (not shown). Former model studies show that for a given shear regime, varying particle properties such as particle adhesion can lead to a change in particle size and thus sinking velocity (e. g. Maerz et al., 2011). Hence, higher particle adhesion would translate to larger particles and

stronger fragmentation resistance, which would allow particles to sink out of the water column closer to the coast under higher turbulence intensity. This implies that spatial changes in SPMC and composition, and accordingly physico-chemical properties of particles, potentially affect the location of the transition zone. As discussed below this probably happens on a seasonal time scale due to the modulation of the transition zone by phytoplankton exudations which requires further investigations. While the cross-shore distribution of w_s is predominantly controlled by prevailing shear, the conceptual model indicates a potential additional relevance of physico-chemical properties for the formation of the high sinking velocity zone.

4.3 Implications of a cross-coastal maximum of sinking velocity for biogeochemical cycling in the coastal zone

The region of $\log_{10}(\epsilon) \approx -5.5$ with highest $\langle w_s \rangle$ coincides with a zone located off the coast at depth of about 15 m to 20 m. It is accompanied by a strong gradient in SPMC and SPM composition that is depicted by the F/SPMC ratio. This suggests that the region can be considered as a transition zone, hindering mineral particles to escape further off-shore. To simplify, consider the course from the coast to open waters. Once formed, relatively dense fast sinking flocs, whose properties are adapted to the transition zone turbulence level, would easily settle out of the water column when transported offshore. This is because turbulence becomes too weak to break those aggregates apart and to retain them in suspension. Thus, only loose, organic-rich particles are kept suspended in the water column while mineral-rich particles tend to settle out of the water column. Such trapping is previously described conceptually for other regions, e. g. by Mari et al. (2012) and Ayukai and Wolanski (1997) for river plumes. High $\langle w_s \rangle$ in the transition zone thus implies an enhanced link between pelagic and near-bottom processes. Among them, different transport mechanisms are discussed that potentially lead, on average, to a net transport of fine sediments shore-wards into the tidal back barriers. *Settling lag* is caused by the different time durations for particles in different water depths to sediment after bottom shear stress for erosion falls below the critical one (Postma, 1961). Another process is *scour lag* that potentially arises from lower bottom shear stresses needed to keep particles in suspension than for their erosion in combination with an asymmetry in the flow (Van Straaten and Kuenen, 1957). These effects are investigated under different topographical settings (van Maren and Winterwerp, 2013) and the authors concluded that w_s in the range of 0.5 mm s^{-1} and 1 mm s^{-1} lead to the highest deposition rates, and underlined the necessity of fine mineral particle flocculation for the build up of tidal flat systems. Recently, another potential major driver, the estuarine circulation, is suggested for sediment accumulation in the Wadden Sea. Residual currents of the estuarine circulation are driven by density gradients that can be caused by i) horizontal temperature gradients, ii) differential precipitation, and iii) river discharges, and iv) all possible superpositions. Residual currents of the estuarine circulation point on-shore in near-bottom, and off-shore in surface waters (Burchard et al., 2008; Flöser et al., 2011). This probably causes a net SPM flux into the Wadden Sea (Burchard et al., 2008, 2013). The latter is hypothesized to act as bio-reactor, where organic components of SPM become re-mineralized and are exported in dissolved form (Postma, 1984; Ebenhöf et al., 2004; Grunwald et al., 2010) by the off-shore-directed component of the estuarine circulation. Along this pathway, dissolved nutrients are assimilated by phytoplankton and thus transferred to bio-particulates. This is likely to happen particularly within the region of prevailing $\epsilon \approx 10^{-5.5} \text{ m}^2 \text{ s}^{-3}$, where phytoplankton experience favorable growth conditions as found by Pingree et al. (1978) for different places on the North-European continental shelf. Algae exude extracellular polymeric substances (EPS) like TEP that can undergo and mediate flocculation by bridging and embedding min-

erals. The latter ballasts the organic matrix and leads to enhanced sinking velocities, as formerly described for the deep ocean (Passow, 2004; De La Rocha and Passow, 2007), which again links pelagic processes to the residual near-bottom currents of the estuarine circulation. In summary, this would imply that the transition zone characterized by its high $\langle w_s \rangle$ accompanied with strong SPMC gradients acts as an important off-coastal feature for closing the near-shore nutrient and mineral cycle. Such a biologically mediated trapping mechanism would retain nutrients and minerals in the coastal zone and inside the Wadden Sea back barriers.

4.4 Spatial biogeochemical implications of a coastal transition zone

Typically, spatial information on biochemical variables are needed to better understand system-wide cycling and fate of matter. Unfortunately, averaging of w_s over bins of ϵ means a possible loss of direct spatial reference. However, our approach made it possible to identify a general pattern of $\langle w_s \rangle$ with a defined maximum along prevailing energy dissipation rates. Even though we neither resolve tidal cycles nor consider potential seasonal variation of $\langle w_s \rangle$ along $\langle \epsilon \rangle$, a map of $\langle w_s \rangle$ allows for insights into the spatial distribution and variability of $\langle w_s \rangle$. It is re-constructed by mapping the found bin-wise relationship between $\langle \epsilon \rangle$ and $\langle w_s \rangle$ (Fig. 4C) onto the respective spatial $\langle \epsilon \rangle$ (Fig. 3). It must be highlighted that local, short term *in situ* measurements would most likely deviate from the obtained averaged picture (Fig. 7) that would be challenging to derive by *in situ* measurements. Nevertheless, strong spatial variability is visible with a pronounced cross-shore gradient of $\langle w_s \rangle$ featuring a maximum of sinking velocities along the coastline as indicated before. However, this maximum varies locally and is particularly less pronounced at the northern German and southern Danish coast, where $\langle w_s \rangle$ declines to values that are in the range of $\langle w_s \rangle$ in deeper waters in the southern German Bight. Applying the aforementioned concept of the transition zone as an off-coastal closing mechanism for nutrient cycling to the two contrasting German Wadden sea regions, East Frisian Wadden Sea and North Frisian Wadden Sea, in the latter particularly the Sylt-Rømø Bight, may help to better understand regional differences of nutrient concentrations. Lower nutrient concentrations and thus lower eutrophication levels in the Sylt-Rømø bight compared to the East Frisian Wadden sea are regularly observed (van Beusekom et al., 2009). In this case, the maximal $\langle w_s \rangle$ calculated in front of the Sylt-Rømø tidal inlet amounts only to $4 \cdot 10^{-4} \text{ m s}^{-1}$, about half the magnitude found in front of the other Wadden Sea regions. Hence, the ability for nutrient retention is diminished and would contribute to the lower nutrient concentrations compared to other Wadden Sea regions. By contrast, the potentially pronounced retention capacity of the transition zone, as predicted for the Dutch coast and German East-Frisian islands may contribute to the phenomenologically described *line of no return* (Postma, 1984), expected further off the coast and characterized by low SPMC. The spatial distribution of $\langle w_s \rangle$ is probably also reflected in the SPMC and their cross-coastal gradients (Fig. 8). Flocs with high w_s are likely mineral-dominated and would rapidly sink out of the water column and thus lead to strong cross-coastal diminishing of SPMC. Generally, dilution of SPMC occurs due to cross-shore wise increasing water depth and local currents parallel to the coast potentially confining horizontal SPMC distribution to the near-coast region (Staneva et al., 2009). Nevertheless, on average, lower ϵ accompanied by lower SPMC and smaller horizontal gradients are found along the North-Frisian than the Dutch and East-Frisian coast (Fig. 3 and 8). In combination, these conditions, in particular along the Sylt-Rømø islands, suppress the establishing of a defined transition zone. Hence, lower w_s should be found as implied by Fig. 7. As a result, retention efficiency is likely reduced and

potentially contributes to the observed lower nutrients concentrations compared to the East-Frisian Wadden Sea (van Beusekom et al., 2009).

Other concepts are presented to explain the observed spatial heterogeneity of eutrophication levels along the Wadden Sea coast. These explanations are categorized by Van Beusekom et al. (2012) in two main groups: either due to i) regional differences in organic matter import, or ii) differences in the size of the tidal basins. Potential differences of the import amount were attributed to different i) regional offshore primary production, ii) orientation of the coastline with respect to dominant wave and current directions and iii) intensity of shore-ward bottom currents. The eutrophication status can be related to the tidal basin width, where narrow tidal basins feature higher eutrophication status than wider ones (Van Beusekom et al., 2012). While the general build-up of nutrients gradients towards the coast is attributed to the above discussed processes of *settling and scour lag* in combination with the estuarine circulation, no clear mechanistic explanation is drawn as to why the size of the tidal basins lead to the found relation. Van Beusekom et al. (2012) suggested that POM imported into the tidal basins would be either distributed over a larger area for wide basins which would lead to gentle nutrient gradients compared to steeper nutrient gradients in narrower basins. However, other explanations cannot be ruled out, e. g. linked to differences in water exchange times (Van Beusekom et al., 2012). It is likely that a multitude of processes interact and detailed modeling studies are required to disentangle their relative contributions to the observed spatial heterogeneity of the Wadden Sea eutrophication status. Noticeably, the Sylt basin exhibits exceptionally low eutrophication status in the study of Van Beusekom et al. (2012), which would fit into the picture that the import of POM is reduced due to the rather weakly established transition zone in this area.

Since the driving mechanisms behind the transition zone are probably ubiquitous for coastal marine systems with sufficiently strong cross-shore water density gradients generated by freshwater run-off, precipitation-evaporation balance or heat fluxes, local hydrodynamics will determine its eventual formation. This implies that the concept of the transition zone with its retention capacity for nutrients may be applicable to other coastal marine and estuarine systems in general, and will help to better understand different nutrient cycling behavior among those systems.

4.5 Biological modulation of the transition zone?

There is increasing evidence that sticky algae excretions such as transparent exopolymer particles mediate aggregation and increase shear resistance of particles against fragmentation, potentially enabling algae to clear the water column from mineral load (Fettweis et al., 2014). As shown in the conceptual model (Sec. 2.4), high w_s occur due to the interplay of shear-driven flocculation and primary particle density and size. Similarly, Hamm (2002) found a non-linear effect of mineral-organic weight/weight-content on sinking velocity showing first an increasing w_s for increasing mineral load, but a stagnation or even converse effect when exceeding a certain threshold. To hypothesize, a high concentration of SPM with an optimal ratio between dense mineral and sticky bio-particles might exist in the transition zone under favorable turbulent conditions to form larger flocs compared to near-coast regions, and denser flocs compared to open waters. In sum, this could lead to the found high $\langle w_s \rangle$. This raises the important question as to what extent algae can affect the transition zone and its seasonal location to sustain the effective coastal nutrient cycle that is driven by the estuarine circulation. By sustaining the nutrient cycle, algae potentially support the development of the pronounced nutrient gradients towards the Wadden Sea. As a consequence, 3-D bio-

geochemical models require a representation of the tight coupling between sediment dynamics and biogeochemical, potentially even adaptive phytoplankton physiological processes to improve models' capability to estimate particulate matter fluxes.

5 Conclusions

The present study provides a strong indication of a maximum of sinking velocities along a cross-shore transect that is defined by a gradient of prevailing energy dissipation rates. Towards the off-shore, the enhanced sinking velocities are accompanied by a strong decline of SPMC, and an increasing F/SPMC ratio. The interplay of processes leading to the observed features defines the region as coastal transition zone. In turn, processes feed back on SPMC and composition.

The transition zone is probably an important feature on the course from the coast to the continental shelf. Predominantly driven by turbulent shear generated by energy dissipation, the transition zone with highest sinking velocities potentially acts as a retention zone for dissolved nutrients leaking from near-shore waters by providing favorable conditions for algae growth. Phytoplankton take up nutrients and excrete extracellular polymeric substances (EPS) that mediate flocculation processes. Embedding minerals into the organic matrix leads to enhanced sinking velocities. Algae thus seem to possess the ability to clear the water column (Fettweis et al., 2014), and link the off-shore dissolved nutrient fluxes to residual landward bottom fluxes. Consequently, these interlinked processes would have the potential to retain fine sediments and nutrients in coastal areas. The strength of the links eventually affects the eutrophication state of the coastal region.

It is scarcely understood what relevance individual processes have in forming the transition zone. Besides favoring energy dissipation rates of about $\epsilon \approx 10^{-5.5} \text{ m}^2 \text{ s}^{-3}$, algae and their extracellular polymeric excretions are potentially important in forming of the transition zone as EPS are known to enhance collision efficiency and resistance against fragmentation. Since primary production in the studied area features a pronounced seasonal cycle, further studies are thus needed to investigate the temporal and spatial extension of the transition zone. Long term or repeated spatial *in situ* measurements, including SPMC, ϵ profiles and particle properties such as size, w_s and (volume-) composition, are required to underpin the indirectly found enhanced sinking velocities as a feature of the transition zone. Additionally, modeling approaches are required to gain a deeper process-based understanding, and to disentangle the closely linked processes among other eutrophication state affecting environmental factors. This also implies the necessity to incorporate biological-mineral interactions in models, especially when system-wide studies are carried out.

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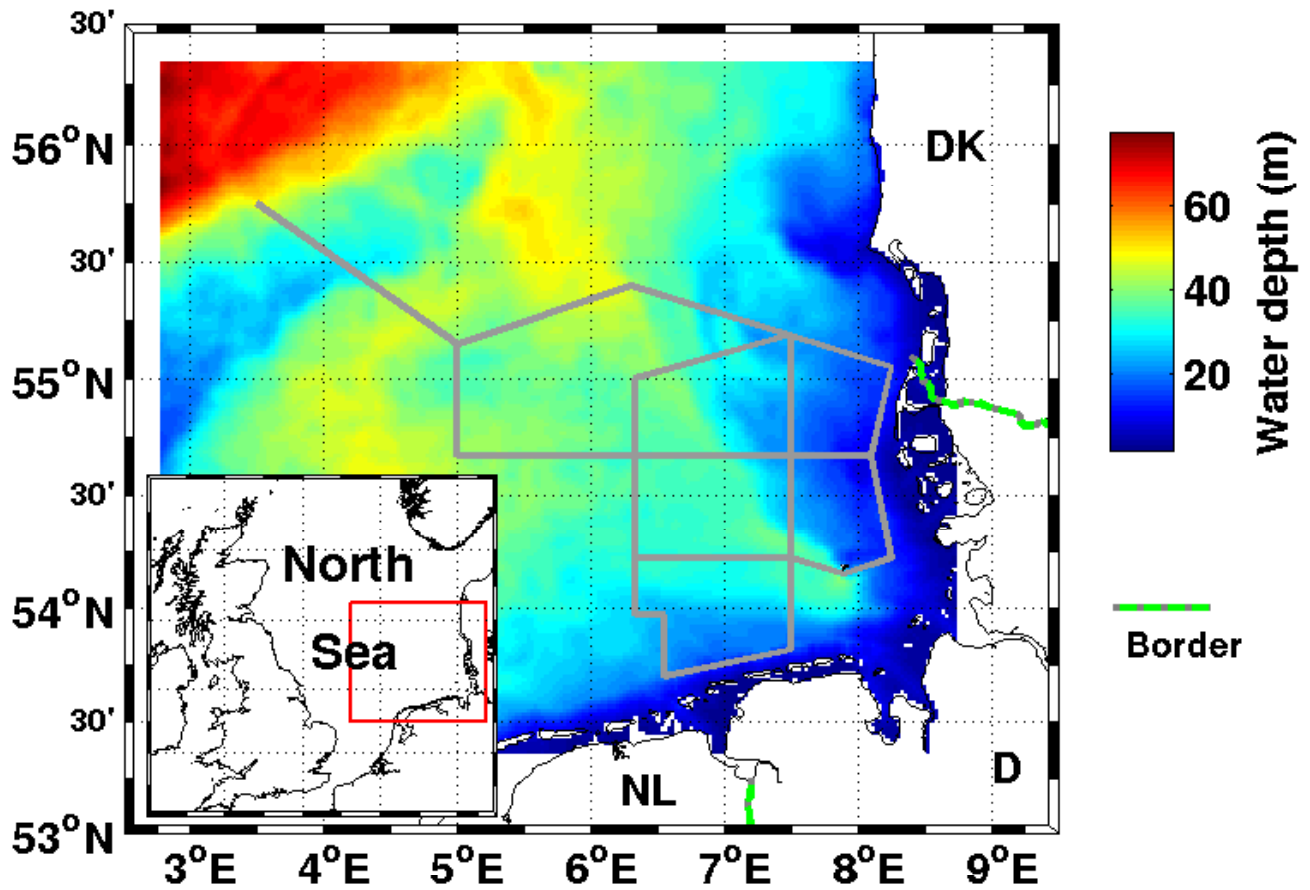


Figure 1. Map of water depth in the German Bight. Gray lines indicate the ScanFish sampling transects. The inset shows the North Sea and in red the drawn region. For East- and North-Frisian sub-regions see Fig. 7.

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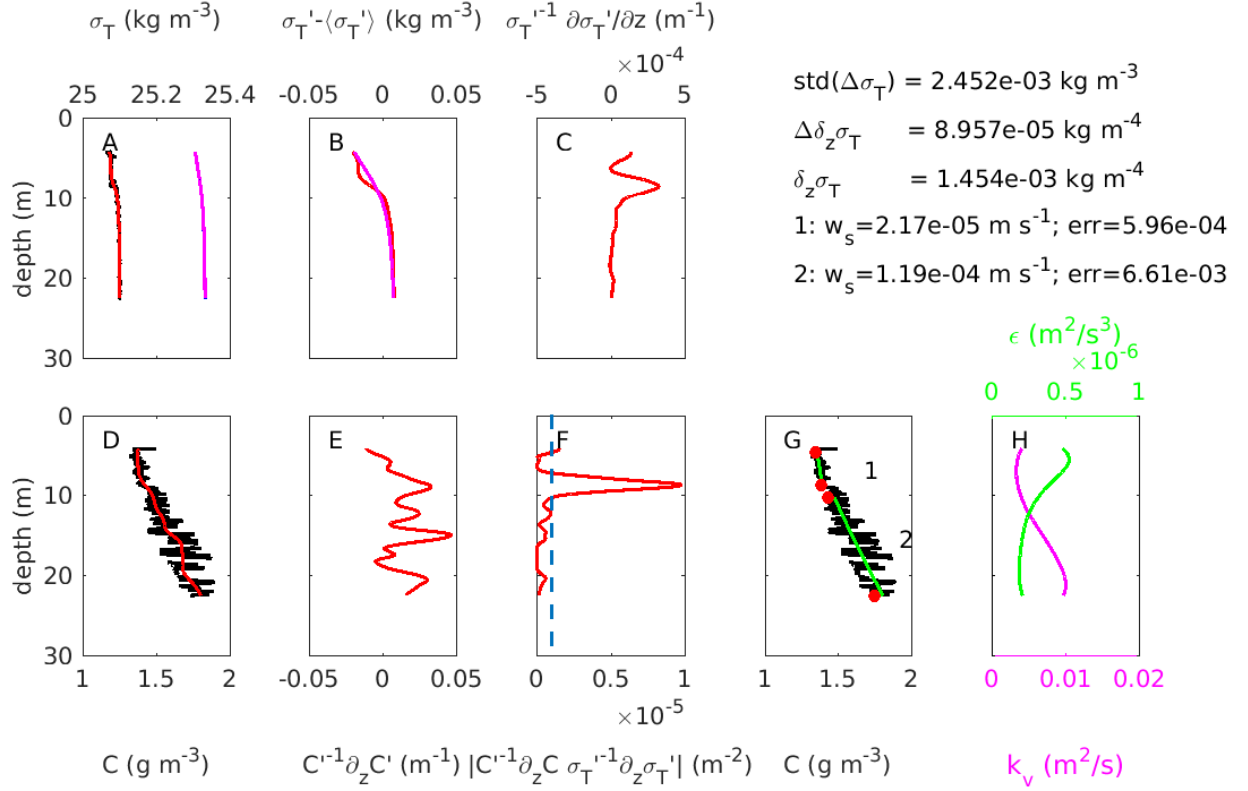


Figure 2. Example for the A: spline-fitted profiles, B: comparison between observed and modeled spline-fitted σ_T' , C-G the splitting and subsequent fitting of the analytical model to vertical SPMC subprofiles; derived w_s and analytical model errors (Eq. (6)) are given by the according numbers. H: turbulence model-calculated ϵ (green) and k_v (magenta) profile. For A-G: lines in black: raw observations; red: spline-fits; magenta: hydrodynamical spline-fits; blue dashed line in F: critical depth averaged value (right hand side of Eq. (5)); green: analytical model fit; red dots: derived split points from Eq. (5).

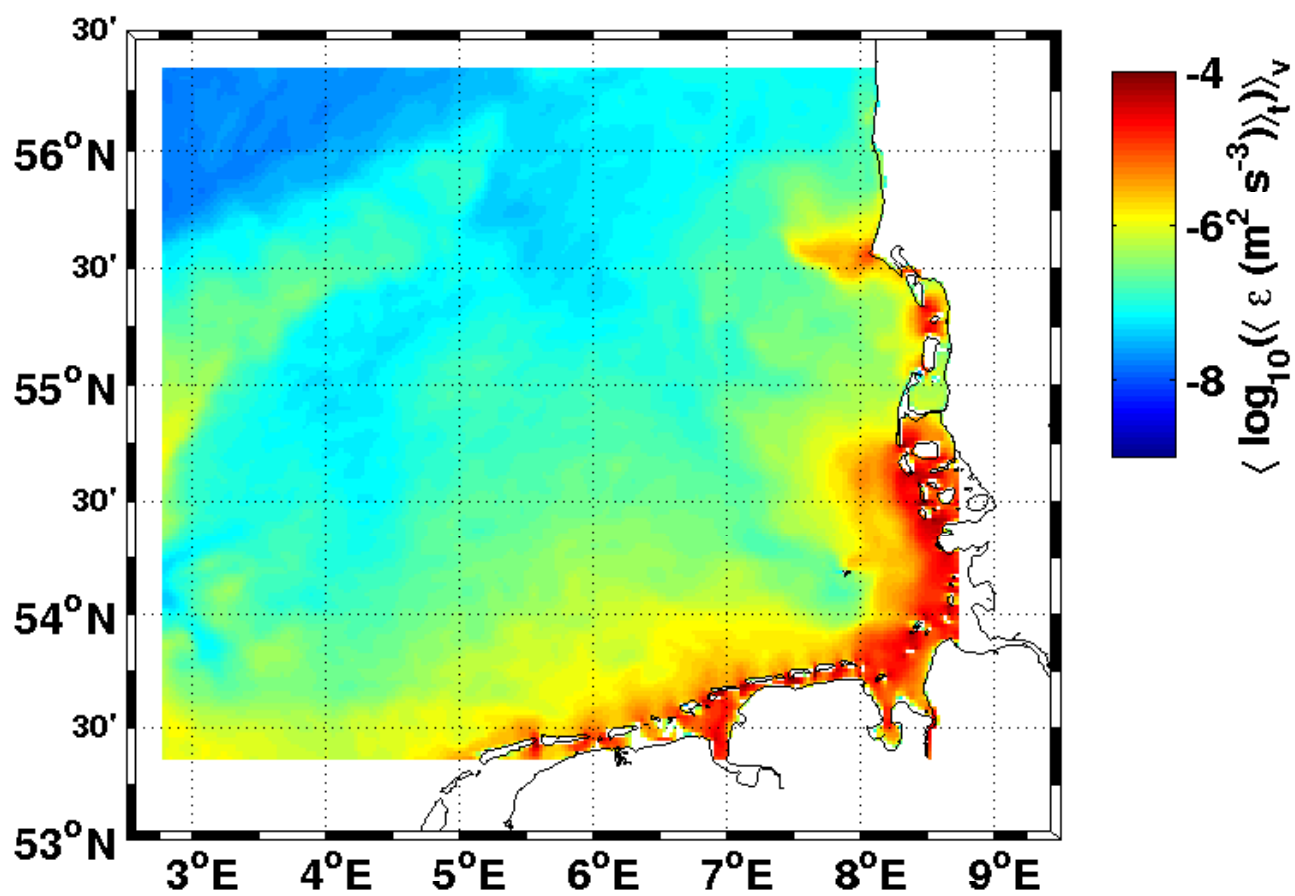


Figure 3. Map of hydrodynamic model-calculated time- and depth-averaged energy dissipation rate ϵ for the times of the cruises.

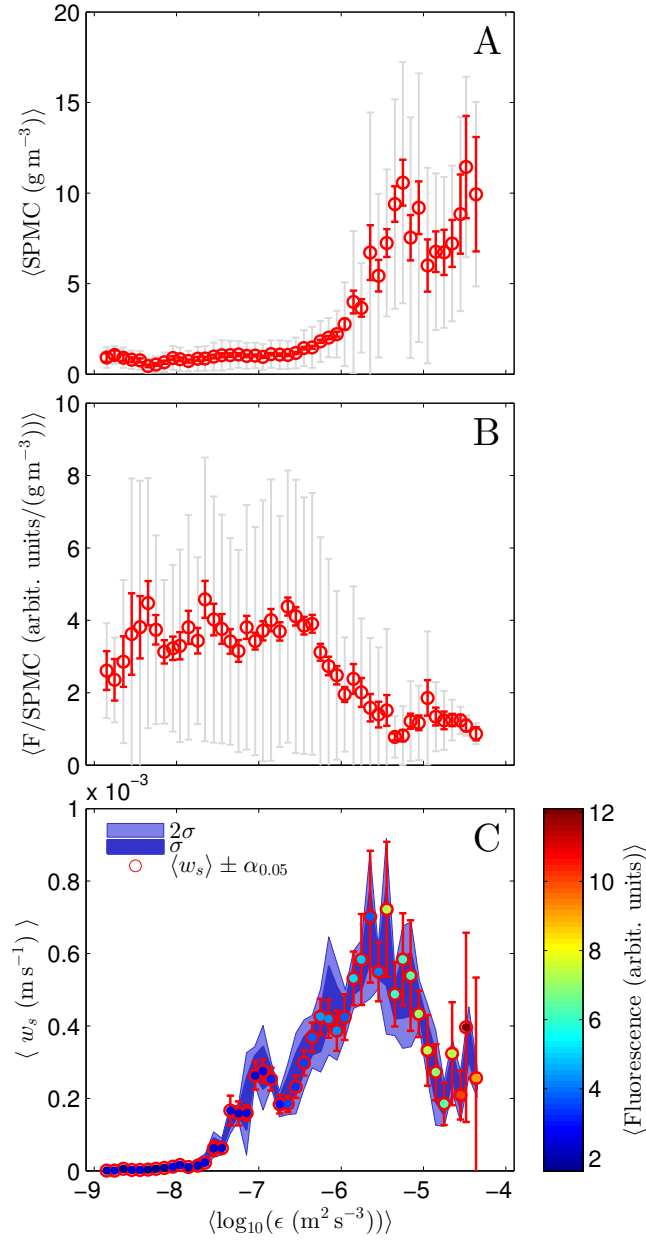


Figure 4. A: $\langle \text{SPMC} \rangle$, B: $\langle \text{F/SPMC} \rangle$ ratio and C: $\langle w_s \rangle$ versus $\langle \log_{10}(\epsilon \text{ (m}^2 \text{ s}^{-3})) \rangle$. Gray error bars represent the standard deviation and red error bars represent the confidence intervals (α -quantile=0.05) for the averages in each bin. Average values for the bins were used for the color coding. In C: σ and 2σ represent the confidence levels of 68% and 95%, respectively, for the Monte-Carlo-type parameter variation.

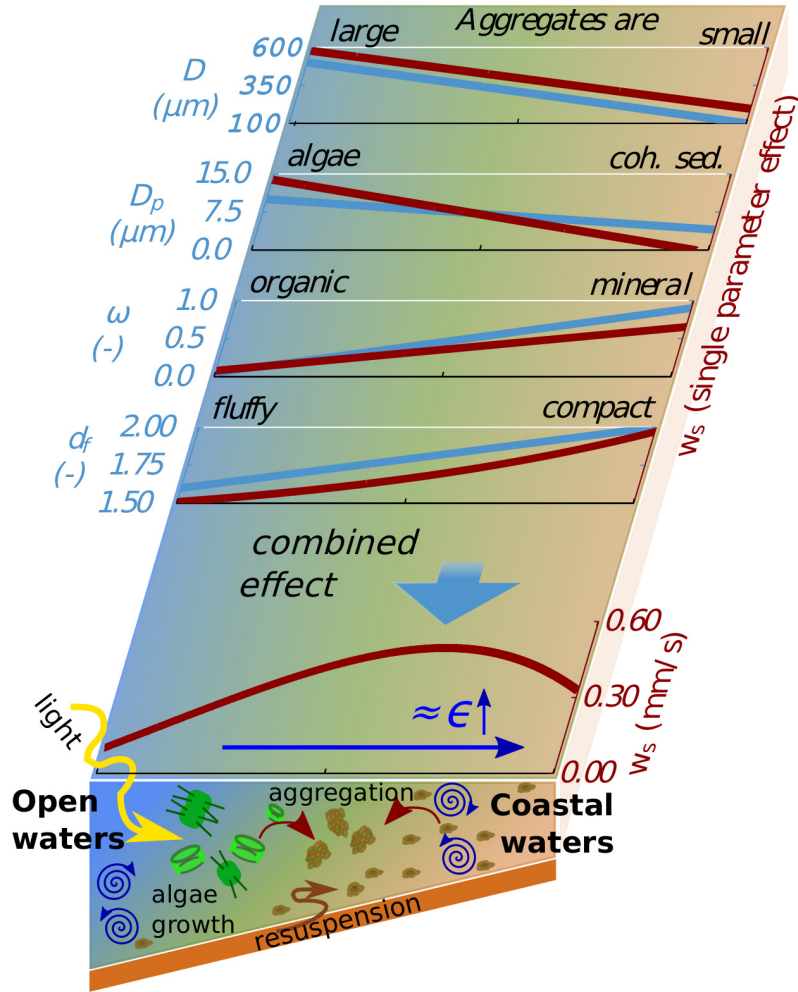


Figure 5. Schematic view of potential influences on sinking velocity (red lines) from open waters (left) to coastal waters (right) and its total impact on the resulting sinking velocity (lowest panel) calculated according to Eq. (10). The sinking velocities shown in the upper 4 panels are based on a particle of $D = 100\mu\text{m}$, $\omega = 0.5$, $D_p = 4\mu\text{m}$, and $d_f = 2$ while the respective parameter shown in the panel is varied (blue lines). (constant parameters are $\mu = 0.001\text{ kg m}^{-1}\text{ s}^{-1}$, $\rho = 1000\text{ kg m}^{-3}$, $\rho_o = 1100\text{ kg m}^{-3}$, and $\rho_s = 2650\text{ kg m}^{-3}$).

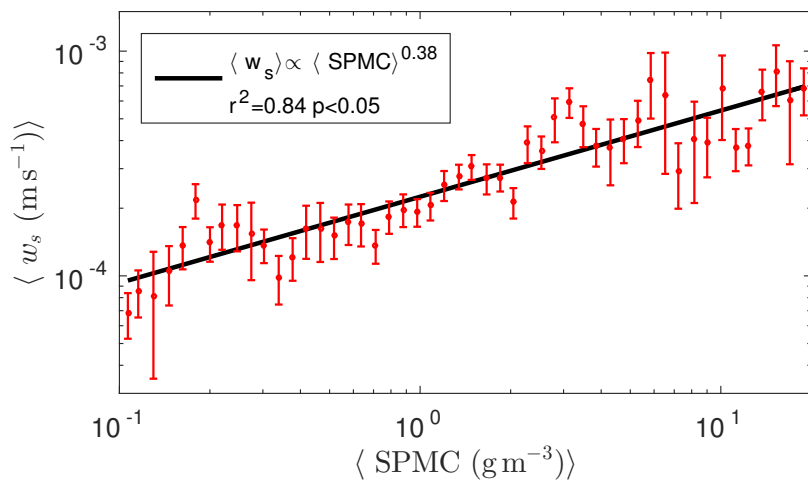


Figure 6. Scatterplot of $\langle \text{SPMC} \rangle$ and $\langle w_s \rangle$. Error bars represent the the confidence intervals (α -quantile=0.05) for the averages in each bin. The rather poor correlation can be attributed to the heterogeneity of the German Bight system. See text for discussion.

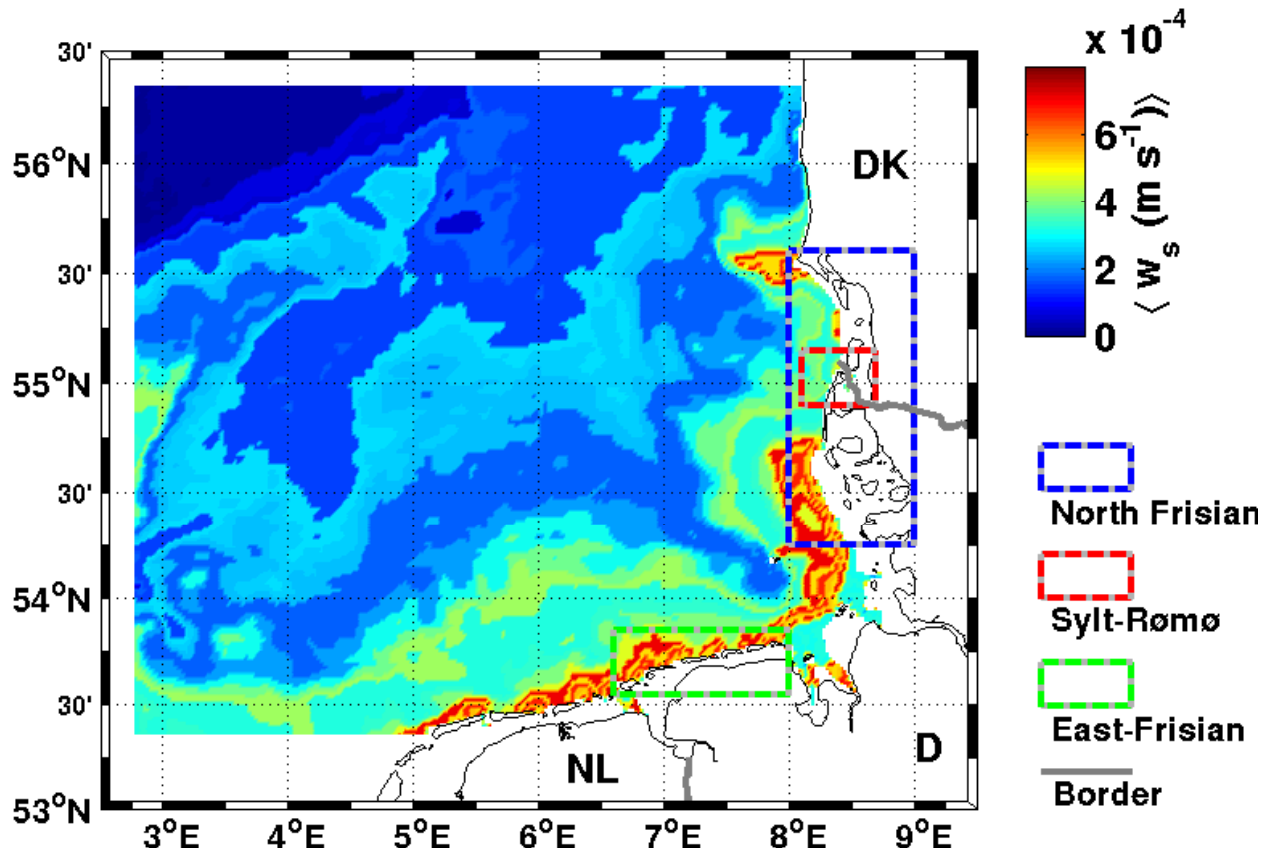


Figure 7. Map of spatial distribution of $\langle w_s \rangle$ in the German Bight. Results for $\langle w_s \rangle$ from Fig. 4 were mapped onto the modeled $\langle \epsilon \rangle$ shown in Fig. 3 to provide spatial information on potentially prevailing $\langle w_s \rangle$. Areas of water depth smaller than 5 m are not considered.

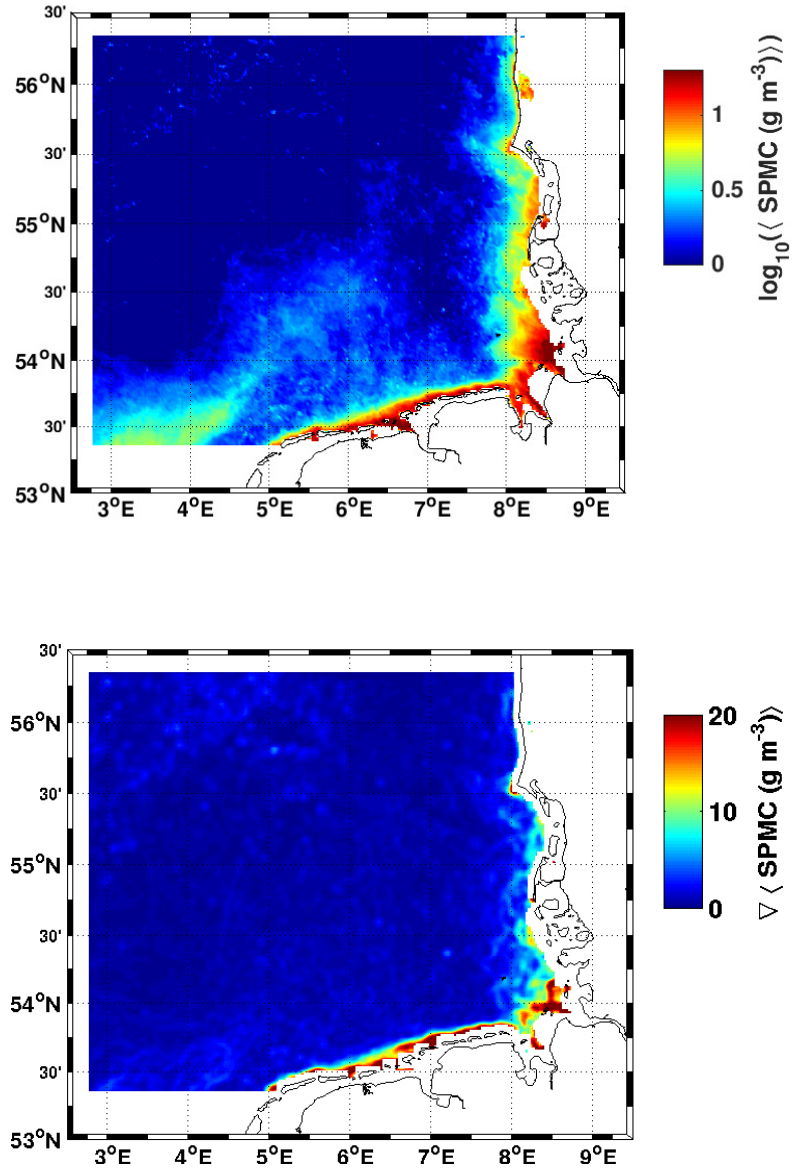


Figure 8. Top: MERIS-derived temporally-averaged SPMC in the German Bight. Averaging of available MERIS scenes was carried out for the times of the cruises. Regions of water depths smaller than 5 m were whited out. Bottom: Pixel-wise gradients of SPMC calculated by $((\partial_x \langle \text{SPMC} \rangle)^2 + (\partial_y \langle \text{SPMC} \rangle)^2)^{\frac{1}{2}}$, where x and y are east-west and south-north pixel directions, respectively. A Wiener filter of 3x3 pixel was subsequently applied. Notice the stronger gradients along the Dutch and East-Frisian coastline compared to the North-Frisian coast with the Sylt-Rømø-bight at 55 °N. For regions see also Fig. 7.