#### **Theoretical insights from upscaling Michaelis-Menten microbial** dynamics in biogeochemical models: a dimensionless approach

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19 Abstract. Leading an effective response to the accelerating crisis of anthropogenic climate 20 change will require improved understanding of global carbon cycling. A critical source of 21 uncertainty in Earth Systems Models (ESMs) is the role of microbes in mediating both the 22 formation and decomposition of soil organic matter, and hence in determining patterns of CO<sub>2</sub> 23 efflux. Traditionally, ESMs model carbon turnover as a first order process impacted primarily by 24 abiotic factors, whereas contemporary biogeochemical models often explicitly represent the 25 microbial biomass and enzyme pools as the active agents of decomposition. However, the 26 combination of non-linear microbial kinetics and ecological heterogeneity across space and time 27 guarantees that upscaled dynamics will violate mean-field assumptions via Jensen's Inequality. 28 Violations of mean-field assumptions mean that parameter estimates from models fit to upscaled 29 data (e.g., eddy covariance towers) are likely systematically biased. Likewise, predictions of CO<sub>2</sub> 30 efflux from models conditioned on mean-field values will also be biased. Here we present a generic mathematical analysis of upscaling Michaelis-Menten kinetics under heterogeneity and 31 provide solutions in dimensionless form. We illustrate how our dimensionless form facilitates 32 33 qualitative insight into the significance of this scale transition and argue that it will facilitate 34 cross site intercomparisons of flux data. We also identify the critical terms that need to be

35 constrained in order to unbias parameter estimates.

# 36 1 Introduction

37 The current crisis of anthropogenic climate change is expected to accelerate during the 21st

- 38 century. Despite considerable effort to better constrain global biogeochemical models,
- 39 considerable uncertainty remains about how best to represent emerging mechanistic
- 40 understanding of soil element cycling into process-based models (Wieder et al., 2015; Todd-
- 41 Brown et al., 2018). This is a critical gap in knowledge because variations among models predict
- 42 hugely varying responses to global change drivers such as temperature, soil moisture, and  $CO_2$
- 43 enrichment. For example, a traditional first-order linear model forecasts no change or even slight
- 44 enhancement of soil organic carbon (SOC) pools by 2100 whereas one microbial-explicit model
- 45 forecasts a loss of ~70Pg of carbon (C), depending on whether microbial physiology acclimates
- 46 to higher temperatures (Wieder et al., 2013). In general, our understanding of how carbon (and
- 47 other elements) cycles in soil is undergoing significant revision toward a more microbial-centric
- paradigm. In contrast to traditional first-order linear models (e.g. CENTURY, Parton et al.,
  1987), microbial explicit models feature non-linear dynamics in which microbial biomass (or,
- similarly, microbially-driven enzyme pools) are responsible for decomposition, in addition to
- 51 providing substrate for synthesis of potentially long-term SOC (Blankinship and Schimel, 2018;
- 52 Blankinship et al., 2018). While indisputably a better representation of our scientific knowledge,
- 53 non-linear microbial models face several well-known challenges, including less analytical
- 54 tractability, greater computational challenges, and uncertainty about structural formulation and
- 55 dynamics (Georgiou et al., 2017; Sihi et al., 2016; Wang et al., 2014). However, one critical
- 56 consequence of non-linear microbial models that is only recently gaining attention is their
- 57 implications for addressing the upscaling challenge.
- 58 While the fields of population and community ecology have long confronted the challenges
- 59 posed by non-linearity and heterogeneity in spatiotemporal scaling of ecological dynamics
- 60 (Chesson, 2009; Levin, 1992), ecosystem ecology and biogeochemistry have tended to approach
- 61 the challenge of scale either by 1) utilizing mean-field assumptions, or 2) addressing the

62 challenge of scaling via grid-based computational/numeric methods. While there is nothing

63 wrong inherently with either approach, they unfortunately cannot yield theoretical insight into

64 the consequences of non-linearity and heterogeneity for scaling. Briefly, the combination of non-

(1)

65 linearity and heterogeneity means that aggregated behavior differs systematically from mean-

66 field predictions, a special case of Jensen's Inequality. In mathematical notation:

 $\mathbf{E}[f(x)] \neq f(\mathbf{E}[x])$ 

Although Jensen's Inequality is well-known from basic probability theory (Ross, 2002) it's

69 implications for ecological dynamics under heterogeneity were not well-appreciated until the

pioneering work of Peter Chesson in the 1990s (Chesson, 1998). In the case of carbon cycle
 science, there are a few immediate and critical applications. For instance, most trace gas

remission processes are well-known to be non-linear functions of underlying drivers such as

73 temperature and soil moisture. For example, ecosystem respiration is an exponential function of

temperature (usually expressed in  $Q_{10}$ ), and a unimodal function of soil moisture. Thus, when

matching observations of  $CO_2$  efflux ("F") to ecosystems, variations in soil temperature and

76 moisture could imply that *F* differs systematically from a mean-field prediction. Likewise,

variations in biotic interactions between microbes likely play a key role in biogeochemical

cycling (Buchkowski et al. 2017). In addition to missing critical analytical insight, not
 accounting for this behavior might have severe consequences for inverse modeling and

estimation of the parameters governing process-based models (Bradford et al., 2021). Moreover,

81 a significant advance in recent research has focused not only on microbial-explicit formulations,

82 but the role of microbe-substrate colocation in the complex and heterogeneous soil environment

in both the synthesis and decomposition of organic matter (Schimel and Schaeffer 2012,

84 Lehmann et al. 2020). This spatial colocation itself has very important implications for scale

transitions in soil systems, and thus requires specific theoretical attention from this perspective.

86 Overall, the basic consequences of Jensen's Inequality for estimation of trace gas emission ( $CH_4$ 

87 and  $N_2O$ ) were first discussed by Van Oijen et al. (2017), but have not been picked up on

elsewhere, until the present work and by Chakrawal et al. (2020)

89 Chakrawal et al. (2020) provide a detailed and compelling first-pass application of scale

90 transition theory to biogeochemical modeling. Our contribution here complements their laudable

91 effort by providing a more generic mathematical analysis of the scale transition, equally

applicable to both forward and reverse Michaelis-Menten microbial kinetics. As in Chakrawal et

al. (2020), we address the consequences of heterogeneity in both substrate/microbes

94 ("biochemical heterogeneity") as well as in the kinetic parameters ("ecological heterogeneity").

95 However, we diverge from their approach in that, rather than explore detailed simulation models,

96 we derive a completely non-dimensionalized expression for aggregating non-linear microbial

97 kinetics over both types of heterogeneity simultaneously. We illustrate the clarity this brings in

98 several special cases of our full analysis. Altogether, our approach provides new insight into the

properties of the scale transition and enables clear conclusions to be drawn across systems in

terms of the role of spatial variances and covariances in shaping ecosystem carbon efflux. Our

101 work provides a simplified, yet systematic framework around which to base subsequent

102 empirical and simulation-based studies.

### 103 **2. Carbon Efflux and the Scale Transition**

104 A variety of microbial-explicit process-based models have been proposed in the literature,

starting with the classic enzyme pool model of Schimel and Weintraub (2003). In order to

106 elucidate universal properties of the scale transition, we focus here on the  $CO_2$  efflux following

decomposition of a single substrate by a single microbial pool obeying Michaelis-Menten (MM)dynamics:

$$F = -f(C, MB, \theta) \tag{2}$$

110

109

111 where F is the  $CO_2$  flux, C is the carbon substrate, MB the live microbial biomass,  $\theta$  is a vector

112 of parameters, specifically  $V_{max}$  (the maximum reaction rate given saturation of either C, in

113 forward MM, or microbial biomass (MB), in reverse MM),  $k_h$  (the half-saturation constant), and 114 carbon-use efficiency,  $\varepsilon$ .

115 Our specific model for *F* is:

$$F = (1 - \varepsilon) \times C \times \frac{V_{max}MB}{k_h + MB}$$
(3)

117

116

118 Following the terminology of Chesson (1998,2012), the above is our "patch" model and our goal

119 is to understand how spatial variances and covariances impact the integrated flux, which

120 represents the spatial expectation or E[F] (hereafter denoted  $\overline{F}$ ), which represents

121 
$$\overline{F} = -C \times (1 - \varepsilon) \times \frac{V_{max}MB}{k_h + MB}$$
(4)

122 Where the bar over the expression represents the mean. The incorrect approach to solving for

123 E[F] is to simply plug-in the mean-field solution:

124 
$$\overline{F} = \overline{-C} \times \overline{(1-\varepsilon)} \times \frac{\overline{V_{max}} \overline{MB}}{\overline{k_h} + \overline{MB}}$$
(5)

125

126 Analytically, an exact solution would require specification of a joint distribution for C, MB and 127 parameters,  $\pi(MB, C, \theta)$ , and solution of the convolution integral:

128 
$$\iiint -f(MB,C,\theta)\pi(MB,C,\theta)dMBdCd\theta$$
(6)

129

130 However, following Chesson (2012) and Chakrawal et al. (2019) we are free to approximate the

131 solution for arbitrary distributions using a Taylor Series approximation expanded to the 2nd

moment. Specifically, we take the expectation over a multivariable Taylor Series expansion,

133 centered around the mean-field values of all parameters  $\theta$  (for simplicity, the variables *MB* and 134 *C* are included in the parameter vector  $\theta$ ):

135 
$$\overline{F} \approx \mathbb{E}\left[f(\overline{\theta}) + \frac{1}{2}\theta_{\theta-\overline{\theta}}^{T}H_{\overline{\theta}}[f(\theta)]\theta_{\theta-\overline{\theta}}\right]$$
(7)

136

137 where  $H[f(\theta)]$  represents the Hessian matrix of the function that determines the  $CO_2$  efflux *F* 138 (in this case Michaelis-Menten),  $\theta_{\theta-\overline{\theta}}$  represents the deviation from the mean at each instance 139 and for each of the parameters. It can easily be seen that  $\theta_{\theta-\overline{\theta}}^T \times \theta_{\theta-\overline{\theta}}$  is the variance-covariance 140 matrix, and that the first moment of the Taylor expansion cancels because the first derivative of 141  $\theta_{\theta-\overline{\theta}}$  is zero.

#### 142 2.1 Non-dimensionalization

- 143 Expanding equation 7 out, we have 5 terms involving the variances of C, MB,  $1 \varepsilon$ ,  $V_{max}$ , and
- 144  $k_h$ , and 10 terms involving covariances among the parameters. We can redistribute the
- 145 expectation operator over this approximation to see that we are dealing with the contributions
- 146 from the variance-covariance terms, weighted by the second partial derivatives evaluated at the
- 147 mean for each parameter. However, the resulting expression **does not readily yield insight into**
- 148 **the impact of scale transition upon the dynamics**, since second partial derivatives and cross
- 149 partial derivatives do not have easy intuition. Moreover, variances and covariances depend
- arbitrarily upon the scale of units and measurements involved, hindering both intuition and cross-
- 151 site comparisons. Therefore, we non-dimensionalize equation 7 for  $\overline{F}$  as follows:
- 152 1) We define a dimensionless quantity  $\lambda$  as  $\frac{\overline{MB}}{\overline{k_h}}$ .  $\lambda$  thus represents a multiplicative factor
- expressing the ratio of the mean microbial biomass over it's half-saturation value, indicatingthe microbial saturation for the decomposition.
- We divide all of the terms in 6 by their mean-field value, and represent the whole equation
  as a product:

157 
$$\overline{F} \approx f(\overline{\theta}) + f(\overline{\theta})((\theta - \overline{\theta})^T \frac{\frac{\partial^2 f}{\partial \theta^2}_{\theta = \overline{\theta}}}{f(\overline{\theta})}(\theta - \overline{\theta})) = f(\overline{\theta})(1 + \left((\theta - \overline{\theta})^T \frac{\frac{\partial^2 f}{\partial \theta^2}_{\theta = \overline{\theta}}}{f(\overline{\theta})}(\theta - \overline{\theta})\right)$$

- 158 (8)
- 159 3) We calculate the resulting expression for  $\overline{F}$
- 160 4) We notice that  $\frac{Var(\theta)}{\overline{\theta}^2}$  can be re-expressed as  $(\frac{SD(\theta)}{\overline{\theta}})^2$  which in turn is the square of the 161 dimensionless coefficient of variation  $(CV(\theta))^2$ . This enables us to reformulate the 162 variance terms in (7).
- 163 5) Similarly, since the covariance terms can be rewritten as  $COV(X, Y) = \rho_{X,Y}SD(X)SD(Y)$ , 164 we have the following equality:

165 
$$\frac{COV(X,Y)}{\overline{XY}} = \rho_{X,Y}CV(X)CV(Y)$$
(9)

Applying steps 1-5 to all the terms in the equation, we end up with a fully dimensionless 167 168 equation:

169  

$$\overline{F} \approx f(\overline{\theta})(1 + \frac{\lambda}{(1+\lambda)^2} [\rho_{k_h,MB} CV(k_h) CV(MB) - CV(MB)^2] + \frac{1}{(1+\lambda)^2} [CV(k_h)^2 - \rho_{k_h,MB} CV(h) CV(MB)] + \frac{1}{(1+\lambda)^2} [\rho_{C,MB} CV(C) CV(MB) + \rho_{V_m,MB} CV(V_m) CV(MB) + \rho_{V_m,MB} CV(V_m) CV(MB)] + \frac{1}{(1+\lambda)^2} [\rho_{C,MB} CV(C) CV(MB) + \rho_{V_m,MB} CV(V_m) CV(MB)] + \frac{1}{(1+\lambda)^2} [\rho_{C,MB} CV(C) CV(MB) + \rho_{V_m,MB} CV(V_m) CV(MB)] + \frac{1}{(1+\lambda)^2} [\rho_{C,MB} CV(C) CV(MB) + \rho_{V_m,MB} CV(V_m) CV(MB)] + \frac{1}{(1+\lambda)^2} [\rho_{C,MB} CV(C) CV(MB) + \rho_{V_m,MB} CV(V_m) CV(MB)] + \frac{1}{(1+\lambda)^2} [\rho_{C,MB} CV(C) CV(MB) + \rho_{V_m,MB} CV(V_m) CV(MB)] + \frac{1}{(1+\lambda)^2} [\rho_{C,MB} CV(C) CV(MB) + \rho_{V_m,MB} CV(V_m) CV(MB)] + \frac{1}{(1+\lambda)^2} [\rho_{C,MB} CV(C) CV(MB) + \rho_{V_m,MB} CV(V_m) CV(MB)] + \frac{1}{(1+\lambda)^2} [\rho_{C,MB} CV(C) CV(MB) + \rho_{V_m,MB} CV(V_m) CV(MB)] + \frac{1}{(1+\lambda)^2} [\rho_{C,MB} CV(C) CV(MB) + \rho_{V_m,MB} CV(V_m) CV(MB)] + \frac{1}{(1+\lambda)^2} [\rho_{C,MB} CV(C) CV(MB) + \rho_{V_m,MB} CV(V_m) CV(MB)] + \frac{1}{(1+\lambda)^2} [\rho_{C,MB} CV(C) CV(MB) + \frac{1}{(1+\lambda)^2} [\rho_{C,MB} CV(C) CV(MB)] + \frac{1}{(1+\lambda)^2} [\rho_{C,MB} CV(C) CV(MB)$$

171 
$$\rho_{\epsilon,MB}CV(\epsilon)CV(MB) + \rho_{C,k_h}CV(C)CV(k_h) - \rho_{V_m,k_h}CV(V_m)CV(k_h) - \rho_{k_h,\epsilon}CV(k_h)CV(\epsilon)] +$$

(10)

172 
$$\rho_{V_m,C}CV(V_m)CV(C) + \rho_{C,\epsilon}CV(C)CV(\epsilon) + \rho_{\epsilon,V_m}CV(\epsilon)CV(V_m)$$

- 173
- 174

Note that by symmetry, we have also solved for the case of the forward Michaelis-Menten 175

176 kinetics. This can be expressed simply by interchanging C and MB, and by correspondingly

177 altering  $\lambda$  to represent the ratio of substrate availability over half-saturation.

#### 178 **3.** Discussion

#### 179 Having fully non-dimensionalized equation 7, we are in a much better position to gain analytical

- 180 insight into the scale transition. To begin, we note the pivotal role played by the quantity  $\lambda$
- throughout this equation.  $\lambda$  scales the contributions of the parameter variation and correlation 181
- terms to the deviation from mean field behavior according to the ratios  $\frac{\lambda}{(1+\lambda)^2}$ ,  $\frac{1}{(1+\lambda)^2}$ , and  $\frac{1}{1+\lambda}$ . All 182
- of the parameter variance terms (which have become  $CV(\theta)^2$  upon non-dimensionalization), are 183
- 184 scaled by one of these three  $\lambda$  ratios, alongside 7 out of 10 of the covariance terms. Overall, low
- 185  $\lambda$  (here  $\lambda < \approx 1$ ) keeps all the spatial correction terms in play, while increasing  $\lambda$  tends to
- simplify matters. As noted by others (Sihi et al., 2016; Buchkowski et al., 2017), as  $MB \rightarrow \infty$ 186
- 187 (equivalent to MB >>  $k_h$  or  $\lambda \to \infty$ ), reverse Michaelis-Menten kinetics converge to first order,
- 188 leaving:

189

$$\overline{F} = -C \times (1 - \epsilon) \times V_{max} \tag{11}$$

190 Accordingly, in our setup, the multiplicative factor for the scale transition correction approaches a simplified expression, as  $\lambda \to \infty$ : 191

192 
$$\overline{F} \to f(\overline{\theta}) \left( 1 + \rho_{V_m, \mathcal{C}} \mathcal{C}V(V_m) \mathcal{C}V(\mathcal{C}) + \rho_{\mathcal{C}, \epsilon} \mathcal{C}V(\mathcal{C}) \mathcal{C}V(\epsilon) + \rho_{\epsilon, V_m} \mathcal{C}V(\epsilon) \mathcal{C}V(V_m) \right)$$
(12)

193 This is quite remarkable. Despite invoking the situation where microbial biomass (and its 194 enzyme supply) is effectively infinite - thus linearizing the underlying patch models - we cannot 195 eliminate the possibility of a potentially substantial deviation from mean-field when scaling decomposition kinetics. We note that in this resulting expression, we have reduced the situation 196 197 to a set of three critical correlations involving two microbial physiological parameters ( $\epsilon$ , and  $V_m$ ), and substrate availability (C). Regardless of their respective variabilities (CV terms), if 198

these correlations are close to zero, then the whole expression converges to mean field. 199

166

200 Returning to the situation where  $\lambda$  is not large, if we ignore the correlation terms (temporarily

201 setting to zero), we see that there are direct contributions to the scale transition from the

202 variability in *MB* and  $k_h$  that may, to some extent, balance each other:

203 
$$\overline{F} = f(\overline{\theta})(1 + \frac{1}{((1+\lambda)^2)} [CV(k_h)^2] - \frac{\lambda}{(1+\lambda)^2} [CV(MB)^2])$$
(13)

204 Focusing on the offsetting correction terms, we can re-write as:

205 
$$\frac{\lambda}{(1+\lambda)^2} \left[ \frac{CV(k_h)^2}{\lambda} - CV(MB)^2 \right]$$
(14)

206 and for the case of  $\lambda = 1$ , this becomes:

207 
$$\frac{1}{4} [CV(k_h)^2 - CV(MB)^2]$$
(15)

Thus, variability in the factors of soil protection that impact upon  $k_h$  in practice, can offset the impact of variability in microbial biomass itself.

210 More generally, starting with our dimensionless equation 10 puts modelers and empiricists in a

211 better position to assess the quantitative significance of the scale transition correction across

212 systems compared to expressions with opaque second partial derivatives and cross derivatives,

and arbitrarily scaled variance terms. By re-expressing  $\overline{F}$  in terms of dimensionless coefficients

214 of variation, correlation coefficients and  $\lambda$ , we can plug-in realistic values for variability in any

215 relevant parameter and assess the % effect on  $\overline{F}$  in terms of deviation from mean field behavior.

216 We argue that this formulation possesses significant advantages not only in understanding how

- 217 to scale flux estimates  $(\overline{F})$  within a site, but going forward will help facilitate intercomparison
- 218 *among* sites in terms of their scale-free variability. In particular, we explore variation in
- 219 dominant environmental drivers of inter-site variation (temperature and soil moisture) below. But
- 220 first, we analyze how the scale transition sheds new light on microbe-substrate colocation.

# 221 **3.1 Spatial Colocation of Microbes and Substrate**

222 To illustrate these advantages in interpretability, we first take the special case of a model where

we treat all parameters as constant (and known) except substrate and microbial biomass. This

224 corresponds to setting the other CV and  $\rho$  terms to 0. In this case, we are isolating the impact of

the spatial colocation of substrate and decomposers. Our equation becomes:

226 
$$\overline{F} \approx f(\overline{\theta})(1 - \frac{\lambda}{(1+\lambda)^2}CV(MB)^2 + \frac{1}{(1+\lambda)}(\rho_{C,MB}CV(C)CV(MB)))$$
(16)

227 In the case of this formulation, there is a very clear dual convergence as  $\lambda$  increases:

- 228 1. deviation from mean-field behavior declines, and
- 229 2. first order kinetics are approached

230 Indeed, our equation 16 reveals the exact speed of this convergence in terms of dimensionless  $\lambda$ 

and a balance of CV(MB), CV(C) and their correlation.

We illustrate the scale transition solutions to equation 16 as a function of  $\lambda$  for various choices of CV(C), CV(MB) and  $\rho$  in Fig. 1:



234

235 Figure 1: Scale transition correction for models given spatial colocation between microbes and

substrate across a gradient of  $\lambda$  values, and for a variety of correlation  $\rho$  values (0-1), with

237 CV(SOC) held constant at 0.5, a) CV(MB) = 0.25, b) CV(MB) = 0.5, and c) CV(MB) = 1.Note

238 that in panel c) the system appears to converge on a value lower than 1. However, as  $\lambda$ 

239 *increases, convergence to 1 does occur, albeit slowly, as it must according to equation 16.* 

240 In the case of pure spatial colocation, with no variation in the kinetic parameters, the scale

transition correction factor varies from a maximum of 1.5 to a minimum around 0.75, and in all cases indeed converges to 1 as  $\lambda$  increases. The variability assumed for C and MB impacts only

242 cases indeed converges to 1 as  $\lambda$  increases. The variability assumed for C and MB impacts only 243 the scale of the correction factor, not the qualitative behavior as  $\lambda$  and  $\rho$  vary. One benefit of

having a simplified, generic dimensionless equation of this sort is that it enables us to think in a

245 unit-free/scale-free manner about the *plausible* range of the scale transition correction given

transparent assumptions about variability and correlations.

- 247 Another benefit is that it is mathematically tractable to see how the variance and covariance
- terms can balance each other, and to solve for where they are equal. If we introduce a new term
- 249  $\lambda_2$  representing the relationship between CV(MB) and CV(C) as follows  $CV(MB) = \lambda_2 CV(C)$ ,

250 we can re-express the the deviation of the mean-field correction from 1 as:

251 
$$CV(MB)^{2}\left[\frac{1}{1+\lambda}\left(\rho\lambda_{2}-\frac{\lambda}{(1+\lambda)}\right)\right]$$
(17)

Thus, whether the correction is positive or negative depends crucially on the product of the colocation correlation coefficient  $\rho$  and the extent of variability in substrate relative to variability in microbes.

255 If we fix  $\lambda$  to unity, as done in our Fig.1, our mean-field deviation simplifies to:

256 
$$CV(MB)^2[\frac{1}{2}(\rho\lambda_2 - \frac{1}{2})]$$
 (18)

- 257 In general, the scale transition correction is larger to the extent that microbial variability exceeds
- substrate variability under reverse Michaelis-Menten kinetics (the opposite relation holds for
- 259 forward Michaelis-Menten by symmetry). Thus, variability in microbial biomass is not only
- 260 important by itself in driving Jensen's Inequality, but also with respect to variability in substrate 261 supply. **Our analysis thus highlights another route of convergence back to the mean field**
- 201 supply. Our analysis thus inglinghts another route of convergence back to the mean field 262 beyond the simple increase of  $\lambda$ : variability in substrate increasing to match variability in
- 263 **microbes in the presence of positive spatial colocation factor**. We also note that the
- magnitude of the scale transition correction scales as the *square* of the coefficient of variation of
- 265 microbial biomass. Quadratic scaling means that at low to moderate levels of variability, the
- 266 deviation from mean field behavior is likely to be minimal, but at moderately high to high levels
- 267 of variability, severe deviations can be expected. Finally, we note that throughout, our
- development of these kinetics assume proportionality to microbial biomass, but it is really the
- 269 live/active fraction that matters. Since the active fraction vary considerably with environmental 270 conditions (e.g. soil temperature and moisture explored below), we believe it is reasonable to
- conditions (e.g. soil temperature and moisture explored below), we believe it is reason
   expect large coefficients of variation overall in most real-world ecosystems.

# 272 **3.2 Environmental Heterogeneity**

- 273 So far, we have analyzed in depth the role of variability in microbes and their substrate, but not
- in the ecological drivers underlying maximal reaction rates (i.e.  $V_{max}$ ) or half-saturation (i.e.  $k_h$ ).
- 275 We start with the observation that both linear first order and non-linear microbial models will

show characteristic scale transitions given heterogeneity in temperature and soil moisture.

277 Consider the asymptotic convergence of the reverse MM to first order

$$\frac{dC}{dt} = -V_{max}\overline{C} \tag{19}$$

279 This is mathematically equivalent to the more standard way of writing these models down as

$$\frac{dc}{dt} = -kC \tag{20}$$

In the analysis that follows, we will consider both temperature and soil moisture as factors that could drive variations in  $V_{max}$  over space or time.

283

278

280

# 284 **3.2.1 Scale Transition over Temperature Heterogeneity**

- To make matters clear, we re-express the rate limiting maximal reaction velocity  $V_{max}$  first as a function of temperature (assuming all else constant):
- $V_{max} = e^{aT} \tag{21}$
- 288 In this case, our integrated flux equation will be:

289 
$$\frac{\overline{dc}}{dt} = \overline{-e^{aT} \times (1 - \epsilon) \times C}$$
(22)

290 Allowing for variability in *T*, this integrated equation will show characteristic scale transitions

291 given the convex (exponential) relationship with *T*.

292 Using the Taylor expansion again to second order we have:

293 
$$\overline{V_{max}} \approx e^{a\overline{T}} (1 + \frac{1}{2}a^2 Var(T))$$
(23)

294 The critical scale transition correction term here is again multiplicative, and we re-express it into

a function of a dimensionless coefficient of variation parameter more suited to ready

interpretation. First, the exponential dependence of respiration on temperature is canonically

297 codified in terms of  $Q_{10}$  scaling. We substitute  $a = \frac{log(Q_{10})}{10}$ , and end up with:

298 
$$1 + \frac{1}{2}a^{2}Var(T) = 1 + \frac{1}{2}\left(\frac{\log(Q_{10})}{10}\right)^{2}Var(T) = 1 + \frac{1}{200}\left(\log(Q_{10})\right)^{2}(SD(T))^{2} =$$

299 
$$1 + \frac{1}{200} (\log(Q_{10}))^2 (\overline{T}CV(T))^2$$
(24)

For a "typical"  $Q_{10}$  of 2.5, and a  $\overline{T}$  of 25, we see the multiplicative scale transition correction in figure 2:



302

Figure 2: Scale transition correction for Q10 temperature response scaling given coefficient of
 variation CV(Temp) from 0 to 0.5

305 As is clear in Fig. 2, the scale transition for temperature is extremely convex. Integration of

306 fluxes over ecosystems with significant heterogeneity in temperature invokes substantial

307 deviation from a mean-field model. For instance, at a CV of 0.2, the scale transition correction is

308 1.10, but by a CV of 0.5 it is 1.66. Obviously, the significance of this depends on the scale and

309 heterogeneity over which an accurate flux model is desired. For a smaller footprint eddy

310 covariance tower (e.g. Gomez-Casanovas et al., 2018) over a uniform habitat type, soil (and near

surface) temperatures probably do not vary by much more than 20%. Regardless, our general

312 mathematical analysis quantifies and clarifies exactly how the scale of variation influences the

313 degree of the scale transition correction.

314 Notably, the only difference between the scale transition correction for first order and for reverse

315 Michaelis-Menten kinetics is that in the latter there would be additional correlation terms to

316 consider, e.g. the correlation between temperature and  $V_{max}$ , temperature and  $k_h$ , as well as

temperature and *C* and *MB*.

### 318 **3.2.2 Scale Transition over Soil Moisture Heterogeneity**

319 Unlike soil temperature, we expect heterotrophic respiration to respond in a unimodal fashion to 320 soil moisture. At low levels of soil moisture, microbes are moisture limited, and at high levels 321 they are oxygen limited, with some optimum range of values in the middle. Although a 322 considerable amount of work has gone into developing soil moisture functions, including both 323 empirical and theoretical derivations (Yan et al., 2018; Tang and Riley, 2019), there is no clear 324 consensus on an optimal representation. Moreover, many of the candidate functions complicate 325 analysis considerably by virtue of stepwise formulation (Linn and Doran, 1984). Therefore, to 326 study the implications of the scale transition we proceed via a powerful simplifying abstraction, 327 and simply represent the soil moisture response as a quadratic of the form:

$$V_{max} = \beta \phi - \beta \phi^2 \tag{25}$$

329 where  $\phi$  represents the soil moisture content. We normalize our function in two senses: first in

330 output space we assume that it has a maximum of 1 (i.e. represents heterotrophic respiration

relative to a maximum value of 1), and second that the soil moisture content  $\phi$  is itself bounded

between 0 and 1, with a peak in the middle at 0.5. Thus, our function captures the unimodal

abstraction in a symmetric form. Given these conditions, there is a unique solution at  $\beta = 4$ .

334 We seek the scale transition:

$$\overline{V_{max}} = \overline{4\phi - 4\phi^2} \tag{26}$$

As before, we can approximate this as a mean-field plus a correction to the mean field, which after some re-writing becomes:

338 
$$\overline{V_{max}} \cong 4\overline{(\phi} - \overline{\phi}^2 - Var(\phi))$$
(27)

339 We then substitute:  $Var(\phi) = \overline{\phi}^2 C V(\phi)^2$ , and re-express:

340 
$$\overline{V_{max}} \cong 4\overline{(\phi} - \overline{\phi}^2 (1 + CV(\phi)^2))$$
(28)

341



342 Figure 3: **a.** heterotrophic respiration as a function of soil moisture given solution in equation

343 25. Note that soil moisture is normalized with a maximum response at 0.5, where 1 represents

344 complete waterlogging, and respiration is normalized to a maximum of 1. **b.** Scale transition

approximate solution for respiration as a function of the dimensionless coefficient of variation of
 soil moisture where mean field soil moisture is either 0.5 (top curve) or 0.25 (bottom curve,

347 equivalent to 0.75 by symmetry).

348 As shown in Fig. 3, where mean field soil moisture is close to the optimal value, scale transition 349 effects are expected to be quite large. For instance, by the time the coefficient of variation is 0.5, 350 efflux would only be 75%, and declines rapidly to 0 as the coefficient of variation approaches 1. 351 Clearly, this latter outcome is not necessarily biologically realistic, and a more detailed 352 numerical experiment should be done to explore scenarios with that much variation. However, 353 our abstractions yield the simple insight that mean-field solutions invariably *overestimate* the 354 real flux, and this overestimation can be considerable. In our experience, soil moisture varies 355 tremendously both over space, especially given contrasts in topography, relief and underlying 356 soils, but perhaps even more so over time, including within a small area, due to day-day and even 357 hour-by-hour variations in precipitation, evapotranspiration and drainage. To the extent that 358 ecosystems deviate from a stable, consistent soil moisture regime, we should expect strong scale

359 transition effects.

360 Our results on soil moisture relate to the argument by Tang and Riley (2019) that heterotrophic

361 respiration arises from a two-step process whereby substrate must diffuse into the vicinity of

362 microbes, and then be taken up – the latter by a Michaelis-Menten kinetic. However, microscale

363 variations in soil moisture mediate and regulate the first step of the process, so that the "effective

substrate affinity" (the  $k_h$  term in the Michaelis-Menten model) deviates from the base substrate affinity of the second step. Tang and Riley (2019) point out that the effective substrate affinity

therefore reflects microscale heterogeneity, and they argue that experimentalists should account

367 for this when fitting efflux data to models. But what about scaling up in the field from small

368 plots to fields to larger ecosystem units? Fortunately, our analytical framework can be readily

369 queried to account for heterogeneity in substrate affinity  $(k_h)$ .

## 370 3.2.2.1 Heterogeneity in Substrate Affinity

371 We proceed by first by holding all terms constant except allowing the half saturation constant  $k_h$ 

to vary, reflecting variations in soil moisture, or frankly any other factor regulating microbial

- access to substrate (e.g. soil aggregation, organomineral complexes, etc. Schmidt et al., 2011,
- 374 Lehmann et al., 2020). As usual, we seek the scale transition over  $\frac{V_{max}MB}{k_h+MB} = f(k_h)$ . We can

375 recover this transition quickly from equation (10) by extracting only the term with  $(k_h)$ , or

376 rederive from scratch holding everything else as constant. The result is that the dimensionless377 scale transition correction term is:

378 
$$1 + \frac{1}{(1+\lambda)^2} CV(k_h)^2$$
(29)

379 Intriguingly, this result shows that heterogeneity in substrate affinity per se results in a convex

380 correction term, implying that mean-field models will under-estimate rather than overestimate

381 the resulting fluxes. Given that the correction is proportional to the inverse of the square of  $\lambda$ ,

- this correction converges rapidly to 1 (no scale transition) as mean microbial biomass increases
- (Figure 4). Nevertheless, where heterogeneity is high and  $\lambda$  is around 1 or lower, the correction
- could be substantial.

385 More broadly, our analysis highlights that, under non-linear Michaelis-Menten kinetics for

386 representing carbon processing, the impact of environmental heterogeneity acting on the

387 substrate affinity parameter is opposite of when it acts on the  $V_{max}$  parameter. Thus, if we

388 represent soil moisture as a modifier to the  $V_{max}$  in the numerator, heterogeneity in soil moisture

389 should result in lower carbon efflux than mean field, whereas if we represent soil moisture

390 heterogeneity by way of substrate affinity it is the opposite. At first glance, this finding appears

inconsistent with Tang and Riley (2019). However, we note that their full kinetic formulations

include soil moisture acting in both roles ultimately, and therefore resulting in the familiar

unimodal soil moisture-respiration relationship. For instance, their application of 'Dual Monod'
 kinetics include soil moisture driving effective substrate affinity terms for both carbon and water,

395 as well as an effective fraction of active microbes (which modifies the numerator).

396 Thus, for the analysis of upscaling fluxes in the presence of soil moisture heterogeneity, we

expect the concave corrections of Figure 3b to hold, regardless of the fine-scale details of the soil moisture function used.



399

400 Figure 4: Scale transition factor for variations in the substrate affinity ("k") parameter in

401 Michaelis-Menten kinetics as a function of dimensionless  $\lambda$  (the ratio of mean-field microbial

402 biomass over mean-field k), for two scenarios of variability in affinity, one where coefficient of

403 variation = 1 (top curve), the other where coefficient of variation = 0.5 (bottom curve).

### 404 **3.3 Lessons for Scientific Inference**

405 We close our discussion by considering the implications of the scale transition for advancing the

406 state of biogeochemical modeling. Critically, the representation of non-linear (microbial driven)

407 kinetics is a crucial modeling choice with large implications for long-term SOC forecasts.
 408 Traditional first-order process-based models dodge explicit representation of these kinetics, but

408 Traditional first-order process-based models dodge explicit representation of these kinetics, but 409 nonetheless have worked well in practice. This state of affairs persists because both non-linear

409 nonemeters have worked wen in practice. This state of analys persists because both non-inteal 410 and linear kinetics are capable of representing coarse-scaled biogeochemistry reasonably well, at

410 least in certain respects. Since first order kinetics are known to be a crude approximation, the

412 crucial question for practice is not whether they are "true", but rather whether there is significant,

412 systematic information loss inherent to their use. Fortunately, the scale transition offers a clear,

414 clean path to discriminate between these alternative model formulations.

415 As noted throughout, the dimensionless term  $\lambda$  plays a critical role in linking the non-linear

416 (Michaelis-Menten) kinetics to the first order kinetics. As  $\lambda$  increases, the non-linear kinetics

417 converge to first order. Thus, in seeking to infer where the non-linear kinetic models provide

418 substantial advantages, ensuring that  $\lambda$  is not too large (>>1) is the first priority.

419 Previous work Sihi et al. (2016) has approached this question theoretically, from first principles.

420 Here, we point out that demonstrating substantial deviation from mean-field model when fitting

- 421 non-linear kinetics to data is both a necessary and sufficient condition for inferring that  $\lambda$  is not
- 422 too large. Thus, we recommend that time series of flux data be fit to both a first order and a non-
- 423 linear kinetic model, where crucial covariates including substrate (SOC), microbial biomass, and
- 424 possibly environmental parameters such as temperature, have been measured sufficiently well to
- 425 quantify the relevant variances and covariances. Where predictive performance and forecasting
- 426 are the primary goals, we recommend careful consideration of model parameterizations
- 427 (i.e. based on leave-one-out cross validation), and model combination via "stacking" where it is
- 428 difficult to infer a decisive "winner" (Yao et al., 2018) acknowledging that carrying this out is a
- 429 significant enterprise.
- 430 In addition to the role of  $\lambda$ , our analysis also cleanly shows the contribution of other terms to the
- 431 scale transition, and thus alternative metrics to assess. First and foremost, accounting for the
- 432 spatial colocation of microbial biomass and substrate (according to equation 16 above) or the
- 433 various correlation terms between microbial biomass and kinetic/environmental factors in
- 434 equation 10. Moreover, recent theoretical developments offer quantitative insights into the
- 435 interpretation of the half saturation constant (or the substrate affinity parameter) and thus
- 436  $\lambda$  (Tang and Riley, 2019). Tang and Riley (2019) decompose microbial access to substrate into a
- 437 two-step process, which is often strongly modified by soil moisture. Moreover, conceptual
- 438 advances suggest that colocation is a potentially important factor in organic matter
- 439 decomposition vs. stabilization (Schimel and Schaeffer 2012, Lehmann et al. 2020). Here, we
- 440 show that both affinity and colocation are co-dependent in their effects on scale transition.
- 441 In addition to fitting fully parameterized flux models (as above), simpler statistical models could
- 442 be fit examining the role of variations in microbial biomass, or colocation of microbial biomass
- 443 and SOC, in explaining across-site variations in ecosystem respiratory fluxes (F). A substantial
- 444 role for either correlation of MB and C, or their variability, would constitute ipso facto
- 445 evidence of the preferability of well-formulated non-linear kinetic models. On the other
- hand, small roles for colocation, or evidence of large values of  $\lambda$  in practice would suggest
- 447 minimal advantage to abandoning first order models in favor of more complex microbial models.
- 448 A meta-analytical approach across sites will benefit greatly from our formulation in terms of
- 449 dimensionless quantities like  $\lambda$  and the various coefficients of variation.
- 450 We further note that the scale transition presented here is closely related to global sensitivity
- 451 analysis (GSA, Saltelli et al. 2010). In its fundamental setup, a GSA tests effects of variability in
- 452 parameters. While GSA has been typically used towards characterizing the uncertainty of
- 453 parameters, it is directly applicable to spatial and temporal variability. For example, the first
- 454 order results of a GSA (or the result of a one at a time parameter substitution), provides the
- 455 contribution of that parameter to the scale transition. Similarly, the 'all but one' perturbation
- 456 offers insights into how the net effect of all parameters (and variables) violates the mean field
- 457 approximation. Therefore, a computationally expensive GSA can be leveraged to garner further
- 458 insights on top of sensitivity effects, allowing for the characterization of the scale transition.
- Indeed, a computationally intensive approach to simulating scale transitions was utilized byChakrawal et al. (2020) to good effect. However, we suggest future computational studies build
- 461 off of the dimensionless approach studied here, including those extended to multiple microbial
- 462 populations which would result in multiple dimensionless lambdas and corresponding
- 463 multiplicative contributions to the scale transition. Obviously, the parameter space needs to be
- 464 properly chosen (or subsetted) to reflect appropriate means, variabilities, and perhaps most

465 challenging - correlations. Equation 10 would then provide analytical, albeit approximate, insight

- into the scale transition effects, while the GSA would enable study of any shortcomings from
- approximation, and also allow for quantification of individual variable importance for those
- 468 parameters that enter into the dynamics in multiple places.
- 469 Finally, our analysis of environmental factors including temperature and soil moisture leads to
- 470 readily testable predictions. For temperature, the scale transition is convex and thus, ceteris
- 471 paribus, variation in soil temperature should lead to greater effluxes than mean field models
- 472 would predict. The implications of this for climate-feedback should be studied in greater detail.
- 473 For soil moisture, which varies considerably across both space and especially time, our analysis
- based on an idealized quadratic representation yields a concave scale transition correction, i.e.
   the mean-field soil moisture will over-estimate efflux. Likewise, when represented in both
- 476 substrate affinity and multiplicative active microbial biomass fraction terms, as in Tang and
- 477 Riley (2019), the scale transition remains concave. However, environmental factors that act only
- 478 through substrate affinity would result in a convex correction as in Figure 4. Once again, we
- 479 highlight that the nature of scaling corrections, wherever it is possible to be studied empirically,
- 480 can provide insight into the most productive representations of our models.

# 481 **4. Conclusions**

- 482 Here, we have illustrated how the spatial scale transition can be expressed in dimensionless form,
- 483 yielding insight into the systematic operation of Jensen's Inequality in upscaling microbial
- 484 decomposition kinetics. Our analysis has identified the central role of the dimensionless quantity
- 485  $\lambda$  representing the ratio of mean-field microbial biomass over its half-saturation value in
- 486 governing the extent of the scale transition correction, expressed here in multiplicative form best
- 487 facilitating comparison among systems. For somewhat simplified scenarios such as restricting
- 488 to spatial colocation of substrate and microbes as  $\lambda \to \infty$ , the mean-field correction goes to 0
- 489 and the model converges to first order.
- 490 This dual sense of convergence also provides opportunity to empirically test for the presence of
- 491 significant non-linear microbial dynamics in upscaled field data: to the extent that upscaled
- 492 fluxes deviate from the flux estimated at mean-field conditions, we have *ipso facto* evidence for
- the importance of formulating our biogeochemical models with these non-linear terms.
- 494 Conversely, where there is close agreement between mean-field and upscaled fluxes, there are
- 495 arguably stronger reasons for retaining first-order process model formulations.
- 496 In closing, we would like to point out how this mathematical analysis illustrates the challenge of
- 497 scaling quite nicely. In the context of non-linear models, for each parameter that is allowed to
- 498 vary in space, there is not only a new variance parameter, but a number of new covariance terms
- 499 are induced, growing as the factorial of the number of varying parameters  $\binom{5}{2}!$  (Fig. 3). Thus, in
- 500 the case of the 5 parameter function considered here, the full approximation has 5 mean field
- 501 terms, 5 coefficients of variation, 10 correlation coefficients, and the dimensionless quantity  $\lambda$ .



Figure 3: Model complexity grows exponentially with number of spatially varying parameters.
We argue to keep models as simple as possible for both analytical and computational

505 *tractability*.

502

506 Even with a maximally generic and simplified expression, fitting such non-linear time series

507 models to field data still represents quite a challenge, especially while adequately accounting for

and propagating uncertainty. Modelers and theoreticians should appreciate the complexity of the

509 task at hand. Fortunately, our analysis has identified a potentially robust route to limiting model

510 complexity: screen systematically for the importance of various correlations in explaining

511 variations in fluxes. Accordingly, we recommend that research focus first upon spatial colocation

of MB and C, which is readily measured, and then to thoughtfully and carefully expand models

513 with additional terms as needed.

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515 analysis, and wrote the manuscript. SG developed the concepts, contributed to the mathematical 516 analysis, and co-authored and edited the manuscript.

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