



On physical mechanisms controlling air–sea CO₂ exchange

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Abstract.

Reducing uncertainties in the air–sea CO₂ flux calculations is one of the major challenges when addressing the oceanic contribution in the global carbon balance. In traditional models, the air–sea CO₂ flux is estimated using expressions of the gas transfer velocity as a function of wind speed. However, other mechanisms affecting the variability in the flux at local and regional scales are still poorly understood. The uncertainties associated with the flux estimates become particularly large in heterogeneous environments such as coastal and marginal seas. Here, we investigated the air–sea CO₂ exchange at a coastal site in the central Baltic Sea using nine years of eddy covariance measurements. To the best of our knowledge, this is the longest record of direct observations of CO₂ fluxes and the corresponding gas transfer velocities (k) in a marine environment. Based on these observations we were able to capture the temporal variability of the air–sea CO₂ flux and other parameters relevant for the gas exchange. The analysis of water-side and atmospheric control mechanisms showed that during wind speeds above 8 m s⁻¹ the conditions on both sides of the air–water interface were relevant for the gas exchange. Our findings further suggest that at such relatively high wind speeds, sea spray was an efficient mechanisms for air–sea CO₂ exchange. In contrast to high wind-speed conditions, during low wind speeds (<6 m s⁻¹) only water-side processes were found to be relevant control mechanisms, in particular, water-side convective processes. Altogether, our results show that currently existing wind-based parametrizations of k might be good approximations as long-term averages for environments with coastal characteristics. However, in order to reduce the uncertainty associated to these averages and produce reliable short-term k estimates, additional physical processes must be considered.

1 Introduction

Air–sea CO₂ exchange is an essential aspect of the global carbon cycle, having great implications for the Earth's climate. The global oceans are estimated to be net sinks of CO₂ taking up on average ca 25 % of the CO₂ emitted every year to the atmosphere due to anthropogenic activities (Friedlingstein et al., 2021). The current global ocean uptake is estimated to be between -2.0 and -3.1 GtC yr⁻¹ (Takahashi et al., 2009; Ciais et al., 2013; Friedlingstein et al., 2021). However, large uncertainties are still associated with the air–sea CO₂ flux estimates, mainly due to the incomplete understanding of the spatio-temporal variability in the controlling mechanisms.



25 The exchange of CO₂ across the air–sea interface can be described using the following bulk formula:

$$FCO_2 = kK_0\Delta pCO_2 = kK_0(pCO_2^w - pCO_2^a), \quad (1)$$

where the air–sea CO₂ flux (FCO_2) is a function of the gas transfer velocity (k), the difference in the partial pressure of CO₂ (ΔpCO_2) between the atmosphere and the seawater (superscripts a and w , respectively), and the salinity- and temperature-dependent solubility constant (K_0). The direction of the flux is determined by the sign of ΔpCO_2 and, by convention, positive
30 (upward) FCO_2 represents transport from the ocean to the atmosphere (i.e. positive ΔpCO_2).

The gas transfer velocity, k , represents the efficiency of the transfer processes across the air–sea interface. For CO₂ and other slightly soluble gases, such efficiency is particularly associated to the turbulent processes occurring in the oceanic boundary layer, which ultimately control the air–sea gas exchange. The wind can be associated, directly or indirectly, with most of the turbulent process near the ocean–atmosphere interface. Thus, the traditional approach suggests that k can be represented as a
35 function of the the wind speed since it is the largest source of kinetic energy to the upper ocean. With wind-speed data being a widely available resource globally, wind-based parametrizations of k have often been used to obtain global estimates of FCO_2 (e.g. Takahashi et al., 2009). However, large uncertainties in the estimated FCO_2 have been associated to the uncertainties in k (Woolf et al., 2019). At regional and local scales, the magnitude of these uncertainties becomes especially problematic, particularly in coastal environments where adequate representation of the physical and biogeochemical processes, and their
40 interactions, is necessary in order to avoid large biases in the flux estimates.

In addition to the wind speed, other water-side control mechanisms are well known to play a significant role in the gas transfer processes of slightly soluble gases, such as CO₂. Furthermore, the effect of atmospheric controls, and their impact on the upper layer of the ocean, are potentially relevant (e.g. Erickson III, 1993) but seldom considered. The relative importance of these forcing mechanisms on the gas exchange is highly dependent on the characteristics near the sea surface, which in
45 turn, can be categorized based on wind-speed regimes (e.g. Soloviev and Lukas, 2013). At moderately high wind speeds, above 8–10 m s⁻¹, the upper layer of the ocean is generally well mixed. Under these conditions, breaking waves (Zhao et al., 2003; Blomquist et al., 2017; Brumer et al., 2017), bubbles (Woolf, 1993, 1997; Bell et al., 2017), and sea spray (Andreas et al., 2016) have a significant role on air–sea interaction processes. At wind speeds lower than 4–5 m s⁻¹, when the effects of breaking waves and wind-induced mixing are limited, convective processes in the atmosphere (Erickson III, 1993) and the sea
50 (Rutgersson and Smedman, 2010) become relevant. Other processes such as surface films (Pereira et al., 2018; Ribas-Ribas et al., 2018), rain (Ashton et al., 2016), Langmuir circulation (Thorpe et al., 2003), and micro-scale wave breaking (Jessup et al., 1997) might be relevant over a wider range of wind speeds, including intermediate wind velocities. In order to explain the variability and reduce the uncertainty in FCO_2 estimates, it is necessary to understand the effect of the control mechanisms on gas exchange, particularly at higher wind speeds but even relevant at low and moderate wind speeds.

55 Coastal oceans and marginal seas are active and heterogeneous environments in terms of both physical and biogeochemical processes. These regions have been found to be net sinks of CO₂ at global scales (Borges et al., 2005; Laruelle et al., 2010; Chen et al., 2013) with disproportionately large contributions to the global carbon system when compared to the open ocean



(Laruelle et al., 2014). The complexity and heterogeneity of the coastal regions causes large spatio-temporal variability in the air–sea CO₂ exchange (Roobaert et al., 2019), variability that is rarely accounted for in global estimates.

60 The Baltic Sea is a semi-enclosed sea located at relatively high latitudes, stretching from 54° N to 66° N. The basin is largely affected by terrestrial inputs from surrounding watersheds and has relatively limited water exchange with the open ocean. This leads to a dynamic carbon system with significant spatio-temporal variability. In terms of the air–sea CO₂ exchange, several approaches have been used to estimate the regional fluxes, however, no consensus has been reached on the role of the Baltic Sea as a net source or sink of atmospheric CO₂ (Thomas et al., 2010; Kuliński and Pempkowiak, 2011; Norman et al., 2013b; 65 Parard et al., 2017). In order to resolve some of the key elements associated to the air–sea CO₂ exchange, previous studies have focused on the diurnal (Honkanen et al., 2021) and seasonal (Thomas and Schneider, 1999; Rutgersson et al., 2008; Schneider et al., 2014) variability in the partial pressure of CO₂ across the Baltic Sea, and on the spatial and temporal variability in the atmospheric CO₂ concentrations (Rutgersson et al., 2009). Furthermore, water-side convection (Rutgersson and Smedman, 2010; Norman et al., 2013b), upwelling events (Norman et al., 2013a; Jacobs et al., 2021), and ice coverage (Löffler et al., 70 2012) have all been recognized as important regional controls on the gas exchange. Despite these efforts, the effect of the different mechanisms modulating the air–sea gas exchange, and its variability, is still poorly understood in the Baltic Sea, as it is in many other coastal regions. Limited data availability is the main reason hindering our ability to resolve processes at relevant spatial and temporal scales. In this context, continuous and long-term monitoring of the air–sea CO₂ exchange in coastal areas is essential to improve our understanding of the gas transfer mechanisms.

75 In this study, we present and evaluate data collected during a nine-year period at the land-based station Östergarnsholm located on an island in the central Baltic Sea. This is, to the best of our knowledge, the longest record of air–sea CO₂ flux based on eddy covariance measurements. Using atmospheric and water-side data we evaluated different control mechanisms modulating the gas transfer velocity, *k*, covering a wide range of wind speed conditions.

2 Site and Data

80 2.1 The Östergarnsholm site

We used data collected between 2013 and 2021 from the Swedish marine Integrated Carbon Observation System (ICOS) station, Östergarnsholm. The station (57° 27' N, 18° 59' E) is located on a small and flat island located 4 km east of the bigger island of Gotland in the central Baltic Sea (Fig. 1). Measurements are performed in a 30 m land-based tower located on the southern tip of the island, with the base of the tower at 1.4 m above the mean sea level (Sjöblom and Smedman, 2002). The 85 tower has been used to monitor and study the marine atmospheric boundary layer and air–sea interaction processes since 1995 (e.g. Smedman et al., 1999; Rutgersson et al., 2001; Högström et al., 2008; Rutgersson and Smedman, 2010; Rutgersson et al., 2020).

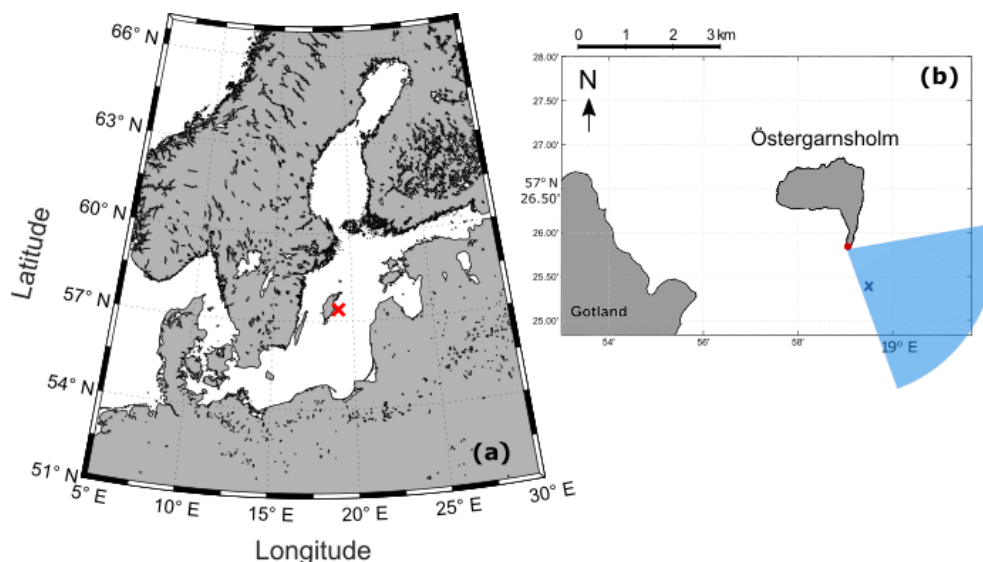


Figure 1. (a) Map of the Baltic Sea; the red cross in the central Baltic Sea indicates the location of the Östergarnsholm station. (b) Map of the Östergarnsholm station ca 4 km off from the Gotland Island; the red dot indicates the location of the tower, the blue cross is the location of the mooring with water-side instrumentation (Sect. 2.1.2), and the shaded blue area is the so-called "open-sea" sector (see text for details).

2.1.1 Atmospheric data

The tower at Östergarnsholm was instrumented with high-frequency sensors for continuous turbulence measurements. CO₂ fluxes were calculated using the fluctuations of the vertical wind component measured by a CSAT3-3D sonic anemometer (Campbell Scientific, Inc., Logan, UT, USA) and the fluctuations of the atmospheric CO₂ molar densities measured with a LI-7500 open-path gas analyzer (LI-COR, Inc., Lincoln, NE, USA). For details about the flux calculations see Sect. 2.2. Both instruments were located at a height of 9 m from the tower base and sampled at a rate of 20 Hz. In addition to the high-frequency data, profile measurements of wind speed, wind direction, and temperature at 7, 12, 14, 20, and 29 m height were carried out at 1 Hz and averaged over 30 min periods. Relative humidity, atmospheric pressure, incoming solar radiation, and precipitation were also measured at the site.

The measurements at the Östergarnsholm site have been found to be representative of open-sea or coastal conditions depending on the wind direction (Rutgersson et al., 2020). Only data with wind directions from southeast ($80^\circ < \text{WD} < 160^\circ$) representing open-sea conditions were included in the analysis of the current study (Fig. 1b). For this wind sector, it was considered that no disturbances occurred in the tower measurements due to flow distortion, and that the wave field was not affected by the shallowing of the seafloor (Högström et al., 2008). Furthermore, the biogeochemical water properties were assumed to be homogeneous for this sector (Rutgersson et al., 2008). A more detailed description of the Östergarnsholm site can be found in Rutgersson et al. (2020).



105 According to the flux footprint estimates (Kljun et al., 2015), the source/sink area of the fluxes from the open-sea sector was located a few hundred meters upwind from the tower (Fig. 2). For unstable and neutral conditions, the main flux source/sink was found as a localized area near the tower. While for stable conditions, the source/sink contributions per unit area were smaller closer to the tower and rather spread over a larger region.

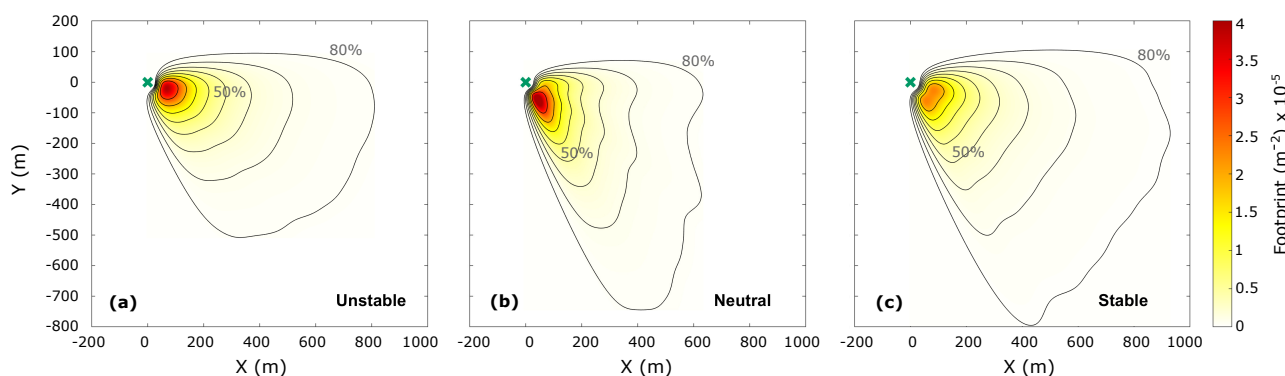


Figure 2. Average footprint distribution for (a) unstable, (b) neutral, and (c) stable atmospheric conditions. The green cross indicates the position of the tower and the contours represent the percentage of source area from 10–80%. The flux footprint was calculated using the model developed by Kljun et al. (2015) using all data available for the open-sea sector between mid-2013 and 2020.

The atmospheric stability was represented as z/L , where z is the measurement height and L is the Monin-Obukhov length. The latter is given by $L = -u_*^3 \Theta / g \kappa \overline{w'T'_s}$, where u_* is the friction velocity, Θ is the potential temperature, $g = 9.81 \text{ m s}^{-2}$ is the acceleration due to gravity, $\kappa = 0.4$ is the Von Karman constant, and $\overline{w'T'_s}$ is the buoyancy flux. The sonic temperature, T_s , was considered to be almost equal to the virtual temperature, T_v , which is often used for buoyancy flux calculations (Aubinet et al., 2012). Following Sjöblom and Smedman (2002), we use $z = 10.4 \text{ m}$ as the mean measurement height with respect to the mean sea level. Here, unstable conditions were defined as $z/L \leq -0.05$, near-neutral conditions $-0.05 < z/L < 0.05$, and stable conditions $z/L \geq 0.05$.

115 2.1.2 Water-side measurements

Water-side measurements were carried out continuously at a mooring located 1 km southeast of the tower (see Fig.1b). At the mooring, seawater temperature and partial pressure of CO_2 were measured every 30 min using a SAMI- CO_2 sensor (Sunburst Sensors, LLC, MT, USA) at a depth of 4 m. Additionally, continuous wave measurements were made with a Directional Waverider buoy located 4 km southeast of the tower (outside the domain in Fig. 1b). The Waverider buoy is operated and maintained by the Finnish Meteorological Institute (FMI). In this study, the characteristics of the sea state were represented by the significant wave height (H_s), the wave steepness calculated as H_s/L_p where L_p is the peak wavelength, and the wave age, C_p/U_{10N} , where C_p is the phase velocity of the waves and U_{10N} is the neutral equivalent wind speed at 10 m height.



Daily data of the mixed layer depth (MLD) from the Baltic Sea Reanalysis product provided by the Copernicus Marine Environment Monitoring System (CMEMS) (Von Schuckmann et al., 2016) were used as an indicator of the vertical mixing
125 in the water column, and for the water-side convection calculations (Rutgersson and Smedman, 2010; Norman et al., 2013b). These data are freely available from CMEMS website at <http://marine.copernicus.eu/services-portfolio/access-to-products/>.

2.2 Data processing

High frequency data obtained at Östergarnsholm site (Sect. 2.1.1) were used for FCO_2 calculations using the eddy covariance method (Baldocchi et al., 1988; Aubinet et al., 2012) and the subsequent gas transfer velocity analysis. The turbulent fluctu-
130 ations for the flux calculation (described below) were obtained following a Reynolds decomposition using an average period of 10 min. These averages were calculated as block averages from the linearly detrended time series of the 20 Hz data. The turbulent fluctuations were used to calculate variances and covariances. Other statistical moments were also calculated from the turbulent fluctuations and used as part of the quality control.

The raw wind-speed components were transformed to earth-system coordinates and corrected using a double rotation
135 (Kaimal and Finnigan, 1994). Wind speed and direction were calculated from the corrected components to avoid effects caused by the tilting of the sonic anemometer. Following the convention, measured wind speeds were adjusted to a neutral equivalent wind speed at 10 m height (U_{10N}). Only wind directions representing open-sea conditions were used in the analysis; see Sect. 2.1.1 for details. Only periods with the three consecutive 10 min averages from $80^\circ < WD < 160^\circ$ (i.e. open-sea sector) were included. Data with wind speeds lower than 2 m s^{-1} was excluded for the FCO_2 calculations.

The performance of the gas analyzer was evaluated based on the relative signal strength indicator (RSSI). Following Nilsson
140 et al. (2018), we used the variance of the RSSI (σ_{RSSI}^2) to remove low quality data. The variance was calculated over the 10 min periods, and only data with $\sigma_{RSSI}^2 < 0.001$ were considered in the analysis. Additionally, thresholds of different statistical parameters were used to ensure the homogeneity of the data and avoid outliers. Data was excluded if the absolute value of the fourth order moment of the CO_2 signal was higher than 100 ppm^4 to filter out outliers, and the variance of the
145 vertical wind speed was $\sigma_w^2 < 1e^{-6} \text{ m}^2 \text{ s}^{-2}$ to exclude unrealistically low values of the vertical wind variance.

According to the eddy covariance method, the FCO_2 were calculated from:

$$FCO_2 = \rho_a \overline{w'c'}, \quad (2)$$

where, ρ_a is the mean density of dry air, and the term $\overline{w'c'}$ represents the time-averaged covariance between the turbulent
150 fluctuations of the vertical wind component (w) and the dry mole fraction of the gas (c). The FCO_2 were calculated over 30 min periods by averaging three consecutive 10 min periods fulfilling all the quality-control steps. The flux was directly calculated from the CO_2 dry mole fraction (i.e. mole fraction of CO_2 relative to dry air) obtained from the measured molar densities of CO_2 relative to the ambient air. By using this direct conversion method (Sahlée et al., 2008), the corrections for dilution effects (Webb et al., 1980) were avoided. A description of the direct conversion method and detailed discussion can be found in Sahlée et al. (2008). Fluxes with magnitudes below a minimum detection limit of $\pm 0.05 \text{ } \mu\text{mol m}^{-2} \text{ s}^{-1}$ were removed.



155 This limit was empirically defined to avoid data with low signal-to-noise ratio. In addition to FCO_2 , enthalpy fluxes were also estimated from the turbulent measurements as the sum of the sensible ($\rho C_p \overline{w'T'}$) and latent ($\rho \lambda \overline{w'q'}$) heat fluxes. Furthermore, cases with high relative humidity conditions ($RH > 95\%$) were excluded to avoid data possibly affected by condensation on the instruments.

The gas transfer velocity was calculated from Eq. 1 using the calculated FCO_2 (Eq. 2). The solubility constant (K_0) was determined from the relationship suggested by Weiss (1974) using a constant salinity value of 7 PSU and in situ water temperature from the SAMI sensor. Changes in the salinity, which oscillates between 6.5 and 7.5 PSU in these region of the Baltic, are not expected to have significant effects on the solubility in comparison to the effect of the water temperature variability. The ΔpCO_2 was obtained from pCO_2^w measured with the SAMI sensor, and pCO_2^a calculated from the molar densities obtained with the gas analyzer. A normalized gas transfer velocity (k_{660}) was calculated as $k_{660} = k(660/Sc)^{-1/2}$, where the reference, $Sc = 660$, corresponded to the Schmidt number of CO_2 at $20^\circ C$ for seawater ($S = 35\text{‰}$), and Sc was the Schmidt number calculated with the corresponding in situ water temperature (T_w) for each data point. Periods with $\Delta pCO_2 < \pm 50 \mu\text{atm}$ were excluded from the analysis. Furthermore, during conditions of strong water-side stratification, pCO_2^w measurements carried out at 4 m depth might not be representative of the air–sea CO_2 fluxes measured at the tower. Therefore, all the data occurring during strongly stratified conditions according to the in situ observations were not considered in the analysis. Data were removed when the water-side temperature gradient (ΔT_w) was larger than $1^\circ C$. The ΔT_w was defined as the difference between T_w measured at 4 m depth and the near-surface water temperature (T_{ns}) measured at 0.35 m depth with the Waverider buoy. Finally, scattered negative k_{660} were expected due to the stochastic behaviour of turbulence. However, more negative k_{660} values than what can be explained by this inherent property of eddy covariance data were observed. There was no methodological reason to exclude such data from the data set as they fulfilled all previous quality control steps. Nonetheless, they were not used for the more detail analysis presented in Sect. 3.2.1 to Sect. 3.2.3, due to the lack of a viable physical explanation.

The calculated k_{660} were used to study the effect of water-side and atmospheric control mechanisms on air–sea CO_2 exchange. In order to remove the wind-speed dependency from k_{660} , a "residual" gas transfer velocity was calculated (k_r). This residual term was calculated for each half-hourly value as the difference between k_{660} and the wind-based parametrization from Wanninkhof (2014) (hereafter referred to as k_{W14}). For the analysis, the data were divided into three different wind-speed regimes from light breeze to moderate gale. The thresholds for these regimes were chosen depending on the expected conditions at the sea surface according to the Beaufort Scale (Barua, 2005). Low wind speeds were defined as $U_{10N} < 6 \text{ m s}^{-1}$, covering conditions of light to moderate breeze; only ripples and small waves causing little disturbance on the surface were expected under these conditions. Intermediate conditions representing moderate breezes included wind speeds of $6 > U_{10N} > 8 \text{ m s}^{-1}$. Finally, relatively high wind-speed conditions were defined as $U_{10N} > 8 \text{ m s}^{-1}$ when moderate to long waves were expected, whitecaps and sea spray were likely to be observed; these wind speeds correspond to fresh breeze to moderate gale.



3 Results

3.1 Oceanic and meteorological conditions: the annual cycle

The annual cycle, obtained using the nine years of data (2013–2021), showed a seasonal pattern in both pCO_2^w and pCO_2^a (Fig. 3a). The pCO_2^a variability was small, at least when compared to the variability in pCO_2^w . The lowest pCO_2^a were observed during the late summer and autumn with values often below $380 \mu\text{atm}$. Higher values occurred during winter, reaching $440 \mu\text{atm}$. While the monthly means of pCO_2^a oscillated around $410 \mu\text{atm}$. An increasing trend in pCO_2^a was observed during the study period (not shown); a linear regression using the monthly averages suggested an increase of 0.2 ppm per month which corresponds to a total increase of approximately 20 ppm during the nine-year period. Both the trend and the seasonal variability in pCO_2^a were masked by the variability in pCO_2^w (Fig.3a). A strong seasonal pattern was observed for pCO_2^w with values higher than those in the atmosphere during the winter, and lower during summer. The lowest pCO_2^w reached values below $100 \mu\text{atm}$ during the summer of 2018. The highest values of pCO_2^w occurred during the winters of 2018–2019 and 2019–2020 with observed values higher than $800 \mu\text{atm}$. Furthermore, lower summer pCO_2^w values were observed in the last three years in comparison to previous years. In this way, the inter-annual variability of pCO_2^w was mostly noticeable in terms of an increasing amplitude of the seasonal cycle during the last years.

The monthly means of FCO_2 showed a seasonal cycle with positive fluxes during the winter and negative during the summer (Fig. 3b). This seasonal pattern in the flux was consistent with the thermodynamic forcing (i.e. $\Delta pCO_2 = pCO_2^w - pCO_2^a$) which suggested an upward transport during the winter and downward during the summer. However, a high variability in the half-hourly data was observed year-round.

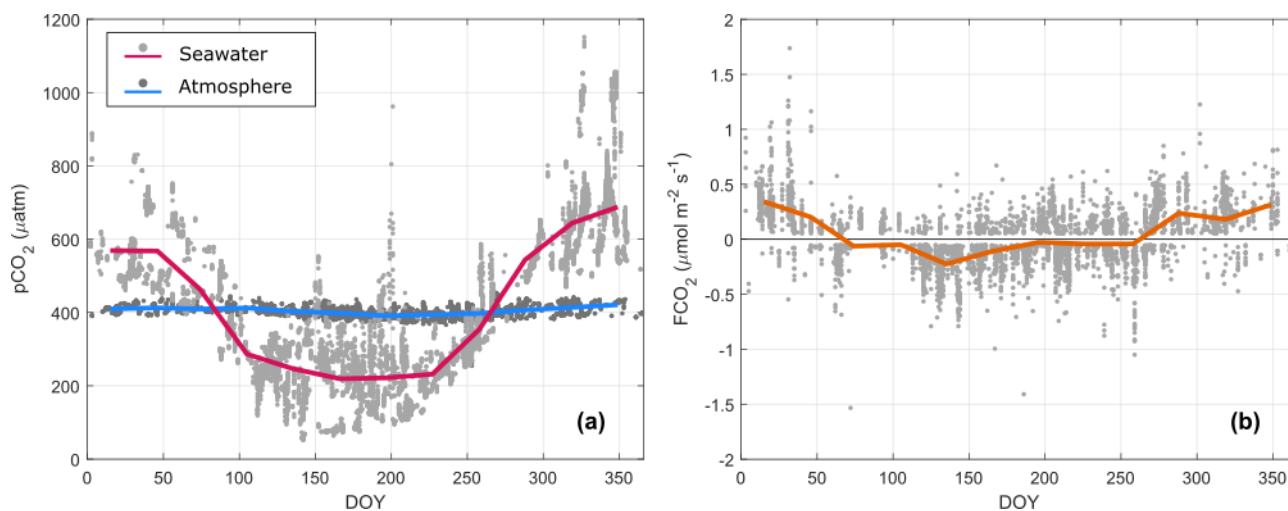


Figure 3. Annual cycle of (a) CO_2 partial pressure ($p\text{CO}_2$) in the seawater and in the atmosphere, and (b) air–sea CO_2 fluxes. The dots represent the half-hourly values while the solid lines show the monthly averages.



The atmospheric and water-side variables describing the physical characteristics showed clear seasonal cycles (Fig. 4). However, a large scatter was observed from the individual half-hourly values, thus, highlighting the large variability and heterogeneity of the environment. The monthly means of wind speed and significant wave height (H_s) were higher during the autumn and winter, in comparison to the summer (Fig. 4a and 4b). However, short-term events with high winds ($U_{10N} > 10 \text{ m s}^{-1}$) and waves ($H_s \sim 3 \text{ m}$) were observed during all seasons. The temperature gradient ($\Delta T = T_w - T_a$, Fig. 4c) showed that from September to February (DOY 250–50) the ocean was, on average, $1.5 \text{ }^\circ\text{C}$ warmer than the overlying air (i.e. positive gradient). During late spring, the atmospheric temperature was warmer than the seawater with an average ΔT of $-1 \text{ }^\circ\text{C}$. The enthalpy flux showed mean monthly values of $0\text{--}75 \text{ W m}^{-2}$ (Fig. 4d). The monthly means of relative humidity (RH, Fig. 4e), ranged between 60 and 85 % throughout the year, but a large scatter was observed. Particularly during autumn and winter (DOY 250–50), when moderately low values of RH ($< 40 \%$) often occurred.

3.2 The gas transfer velocity

The normalized gas transfer velocity, k_{660} , followed a clear increasing relationship with wind speed (Fig. 5). At low and moderate wind speeds (up to ca 10 m s^{-1}), the bin-averaged values of k_{660} followed commonly-used parametrizations. However, at wind speeds higher than 10 m s^{-1} , the wind-speed dependency was more pronounced than the quadratic (Wanninkhof, 2014) and cubic (McGillis et al., 2001) relationships. Although the bin-averaged k_{660} values were generally in agreement with published parametrizations, we observed a large scatter in the half-hourly values along the entire wind-speed range. To further assess this variability in k_{660} , we evaluated atmospheric and water-side controls involved in the gas exchange under relatively high (Sect. 3.2.1), low (Sect. 3.2.2), and intermediate (Sect. 3.2.3) wind-speed conditions. For the description of the wind-speed regimes used in the following analysis see Sect. 2.2.

3.2.1 Controls on k_{660} at high wind speed conditions

At high wind speed conditions, the characteristics of the wave field (Fig. 6) were found to be associated to the gas exchange. The residual gas transfer velocity, k_r , showed a clear relationship with the significant wave height (H_s) as the highest values of k_r occurred during high-wave conditions of $H_s > 1.5 \text{ m}$ (Fig. 6a). Low values of wave steepness (H_s/L_p) were observed even at the highest wind speeds with maximum values of 0.06 (Fig. 6b), much lower than the theoretical wave breaking threshold (Stokes, 1880). While the wave age (C_p/U_{10N} , Fig. 6c) suggests that most of the waves were locally generated (i.e. wind sea) at these wind speed conditions.

In addition to the wave field, we found ΔpCO_2 and MLD to be relevant water-side parameters which were associated to the behaviour of the gas transfer velocity (Fig. 7). Note that ΔpCO_2 was considered here as a water-side control due to the importance of pCO_2^w in modulating the variability of the gradient, in comparison to the rather constant values of pCO_2^a (see Sect. 3.1). Under conditions of strong gradients, where the absolute magnitude of ΔpCO_2 was higher than $200 \text{ } \mu\text{atm}$, k_{660} followed the wind-based parametrization ($k_r \approx 0$). Nonetheless, the largest values of k_{660} were not related to such strong gradients, instead, they were related to moderate positive ΔpCO_2 values in the order of $50\text{--}100 \text{ } \mu\text{atm}$ (Fig. 7a). Large negative values of ΔpCO_2 were, in most cases, linked to negative values of k_r , thus, lower-than-expected from the wind-

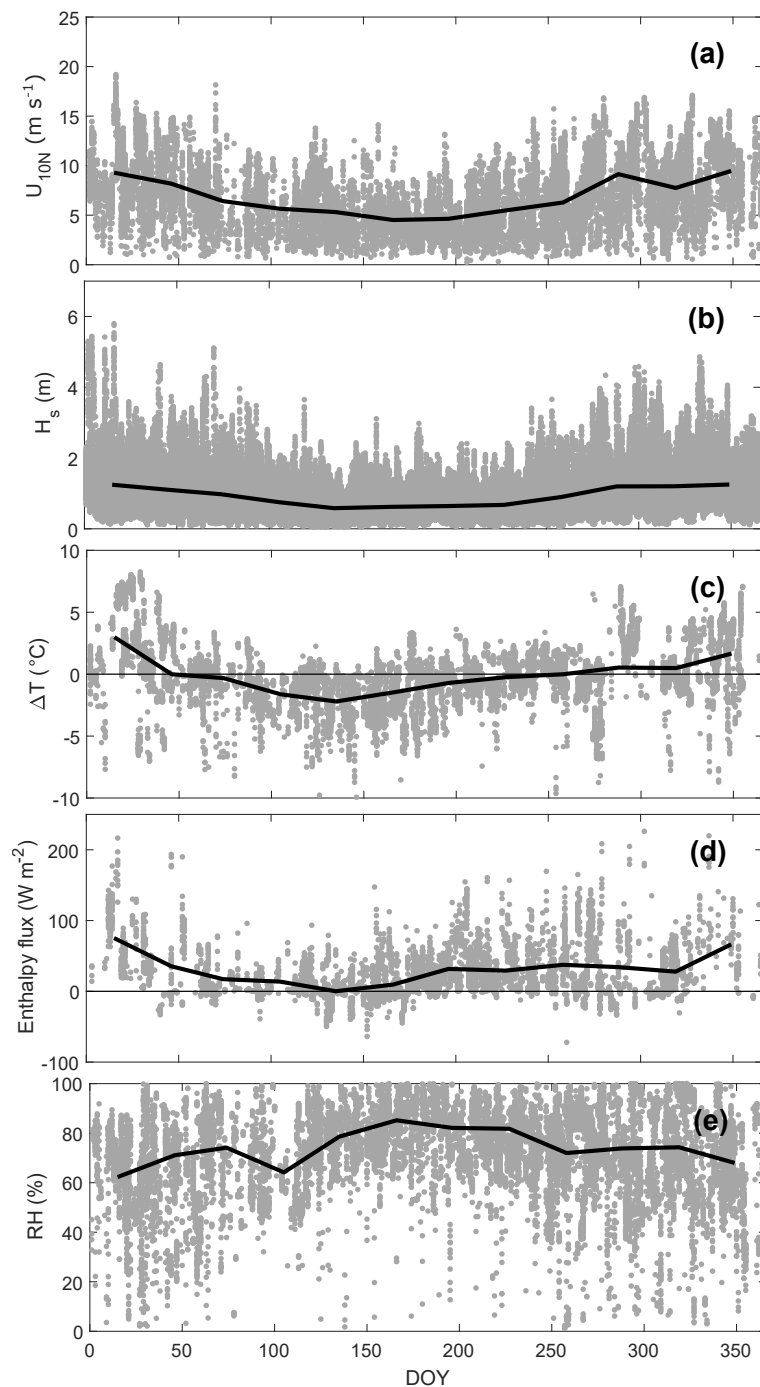


Figure 4. Annual cycle of (a) 10 m neutral wind speed, (b) significant wave height, (c) temperature gradient ($\Delta T = T_w - T_a$), (d) enthalpy flux, and (e) relative humidity. The dots are the half-hourly values while the solid black lines represent the monthly average of each parameter.

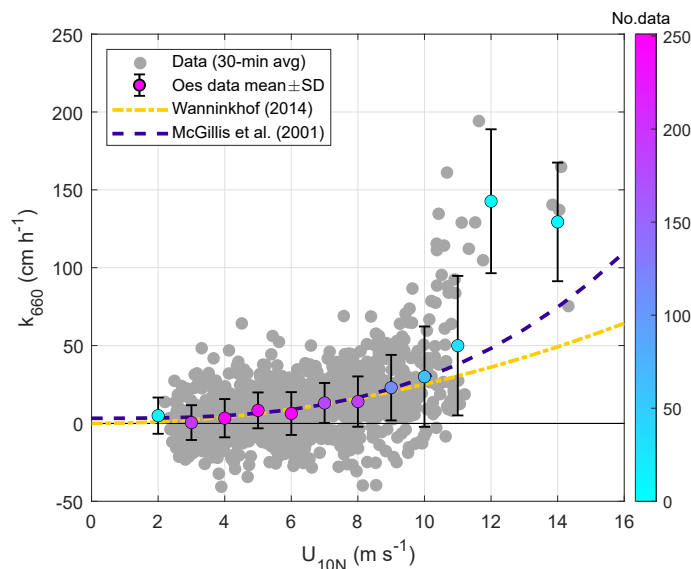


Figure 5. Gas transfer velocity for CO_2 (normalized to a Schmidt number of 660) as a function of the 10 m neutral wind speed. For reference, a quadratic (Wanninkhof, 2014) and cubic (McGillis et al., 2001) wind-based parametrizations were included. The grey dots represent the half-hourly values, the colored circles with bars represent the averages over 1 m s^{-1} bins and $\pm 1\sigma$. The colorbar indicates the amount of data (No. data) available per bin. The calculations of k_{660} were based on the data retrieved during the nine-year period from 2013 to 2021.

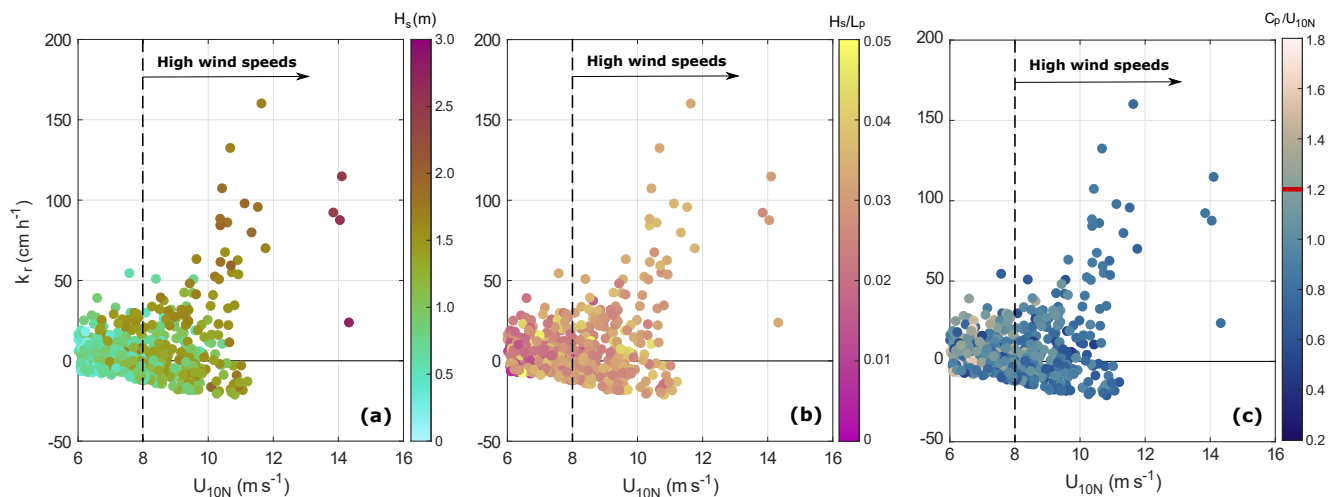


Figure 6. Residual gas transfer velocity ($k_r = k_{660} - k_{W14}$) as a function of the 10 m neutral wind speed under high wind speed conditions. The color represents a) significant wave height, b) wave steepness (H_s/L_p), and c) wave age (C_p/U_{10N}). The vertical dashed lines mark the threshold between intermediate ($6 < U_{10N} < 8 \text{ m s}^{-1}$) and high ($U_{10N} > 8 \text{ m s}^{-1}$) wind speeds. The red mark in c) indicates the theoretical threshold between wind sea ($C_p/U_{10N} < 1.2$) and swell ($C_p/U_{10N} > 1.2$).



240 based parametrization. While, moderate negative values of ΔpCO_2 occurred for both positive and negative k_r values. At these relatively high wind-speed conditions ($U_{10N} > 8 \text{ m s}^{-1}$), the water column was well mixed and, while some scatter is observed, the largest values of k_r were associated with a deep mixed layer (Fig. 7b). Relatively shallow MLD values remained closer to the $k_r = 0$ line and below it. The combined effect of ΔpCO_2 and MLD, was strongly modulated by seasonal patterns. During the winter, strong and persistent vertical mixing occurred along with positive ΔpCO_2 (Fig. 3a). During the summer, shallower MLD and negative ΔpCO_2 were observed.

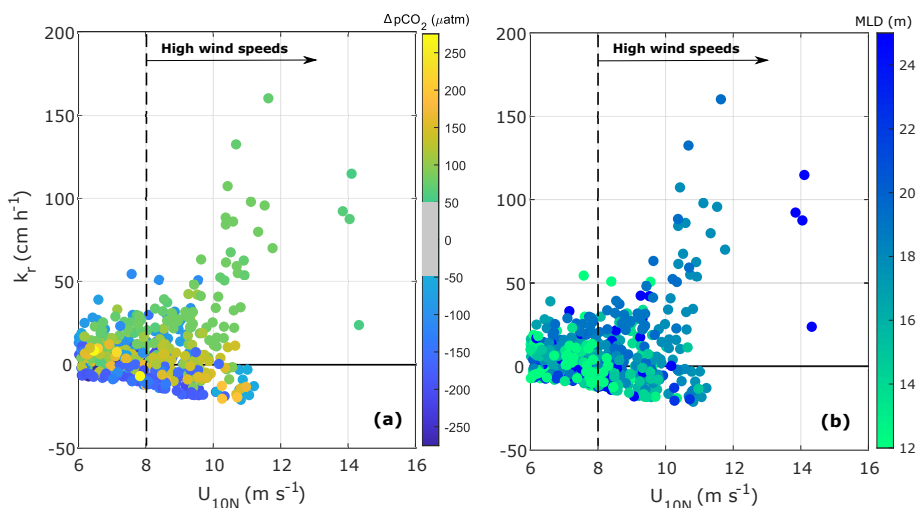


Figure 7. Residual gas transfer velocity ($k_r = k_{660} - k_{W14}$) as a function of the 10 m neutral wind speed under high wind speed conditions. The color represents a) the ΔpCO_2 , and b) the mixed layer depth (MLD). The vertical dashed lines mark the threshold between intermediate ($6 < U_{10N} < 8 \text{ m s}^{-1}$) and high ($U_{10N} > 8 \text{ m s}^{-1}$) wind speeds.

245 Atmospheric conditions such as atmospheric stability, relative humidity, and the total enthalpy flux seemed to be linked to the gas exchange efficiency as well (Fig. 8). In terms of the atmospheric stability, the differences between unstable, neutral, and stable conditions became evident at the highest wind speeds (Fig. 8a). An increased wind-speed dependency of k_{660} (and therefore of k_r) was observed under unstable conditions for wind speeds higher than 8 m s^{-1} . While k_r remained close to zero up to 10 m s^{-1} during stable and neutral conditions. At wind speeds higher than 10 m s^{-1} , k_r tended to decrease with wind speed during stable conditions, and to a lesser extent under neutral conditions. Furthermore, high k_r values were observed when high enthalpy fluxes occurred; meanwhile, the lowest values of the enthalpy flux were associated to negative k_r values
250 (Fig. 8b). The relative humidity seemed to also play a role, with high RH values associated to negative k_r for wind speeds higher than 10 m s^{-1} . On the contrary, low RH were observed mostly during positive k_r (Fig. 8c). Some low RH values linked to negative k_r were observed at wind speeds higher than 10 m s^{-1} , however, these values were also related to close-to-zero values of the enthalpy flux. According to the results in Fig. 8, dry air during unstable atmospheric conditions was not only associated to large enthalpy fluxes, but it also triggered a significant increase in the efficiency of the gas exchange.

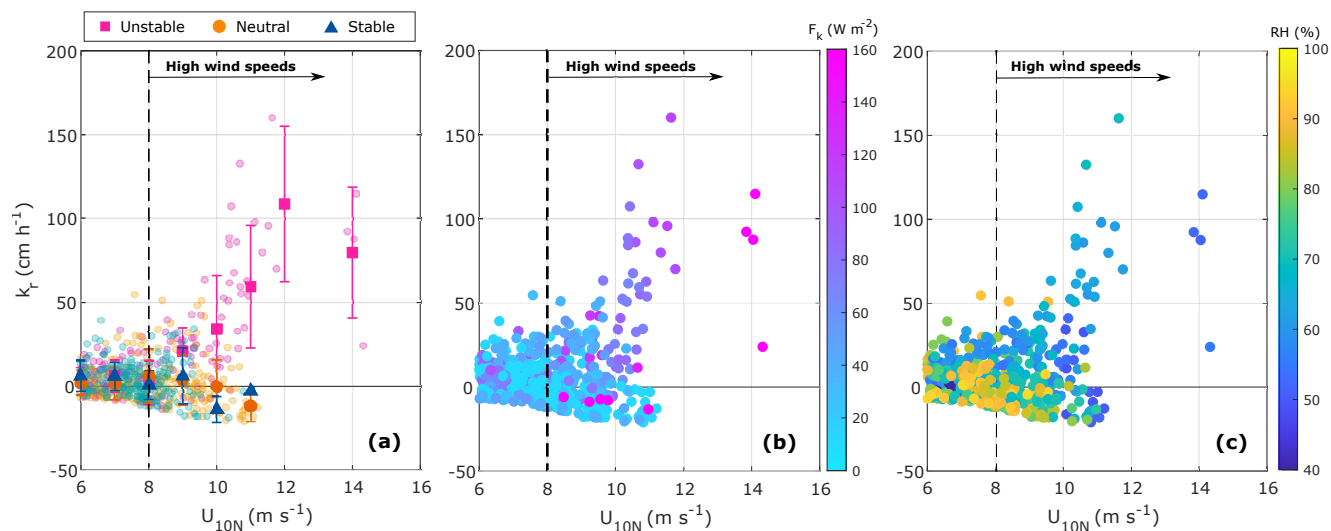


Figure 8. Residual gas transfer velocity ($k_r = k_{660} - k_{W14}$) as a function of the 10 m neutral wind speed under high wind speed conditions. The color represents a) atmospheric stability, b) enthalpy flux, and c) relative humidity. In a) the filled symbols represent the bin average under each atmospheric regime and the bars indicate $\pm 1\sigma$. The vertical dashed lines mark the threshold between intermediate ($6 < U_{10N} < 8 \text{ m s}^{-1}$) and high ($U_{10N} > 8 \text{ m s}^{-1}$) wind speeds.

255 3.2.2 Controls on k_{660} under low wind-speed conditions

Under low wind-speed conditions ($U_{10N} \leq 6 \text{ m s}^{-1}$) the wave field showed smaller ($H_s < 1 \text{ m}$) and less steep waves (Fig. 9a and b). Under these conditions, the waves tended to be older ($C_p/U_{10N} > 1.2$, Fig. 9c), indicating a larger proportion of swell waves in comparison to the locally generated waves observed at higher wind speeds. The characteristics of the wave field did not show a clear relation with the behaviour of k_r . Furthermore, the analysis of other water-side parameters such as ΔpCO_2 , MLD, and water-side convection helped explain part of the variability in the gas transfer velocity (Fig. 10). The largest ΔpCO_2 values, with absolute magnitudes higher than $200 \mu\text{atm}$, were observed within $\pm 10 \mu\text{atm}$ of the horizontal line, i.e. $k_r \approx 0$ (Fig. 10a). Positive ΔpCO_2 corresponded to positive k_r , while strong negative values of ΔpCO_2 were mostly associated with negative k_r for wind speeds below 6 m s^{-1} , but a large scatter is observed. Under these calm wind speed conditions, the values of MLD were generally low, with the lowest ($\text{MLD} < 15$) values showing a larger scatter (Fig. 10b). The pattern of ΔpCO_2 and MLD was associated with the seasonal cycles, suggesting that shallow MLD typical of the summer months may hinder the downward transport of CO_2 into the sea. In addition to ΔpCO_2 and the MLD, water-side convection was a relevant parameter under unstable atmospheric conditions (Fig. 10c). Low values of the water-side convective scale (w^*), as defined in Rutgerosson and Smedman (2010), were observed on and below the horizontal line (i.e. $k_r \approx 0$). Meanwhile, some of the scattered data could be associated with larger values of the water-side convective scale ($w^* > 0.01 \text{ m s}^{-1}$). Water-side convection seemed to be particularly relevant for positive fluxes, which occurred during the winter when $\Delta T > 0$ (Fig. 4c) indicating cooling of the surface. The rest of the scattered data which occurred during negative ΔpCO_2 and shallow MLD (Fig. 10a and 10b)



cannot be explained by water-side convection or any other parameter assessed in this study. Furthermore, the interpretation of these data should be taken with some caution as the strong stratification, relatively weak ΔpCO_2 , and the possibility of strong heterogeneity in terms of the biogeochemical properties might hinder our capacity to calculate k_{660} from pCO_2^w and FCO_2 . Under neutral and stable atmospheric conditions, water-side convection is not a relevant mechanisms (Rutgersson and Smedman, 2010; Rutgersson et al., 2011; Norman et al., 2013b). At low wind speed conditions, atmospheric controls (not shown) did not play any significant role on the gas exchange.

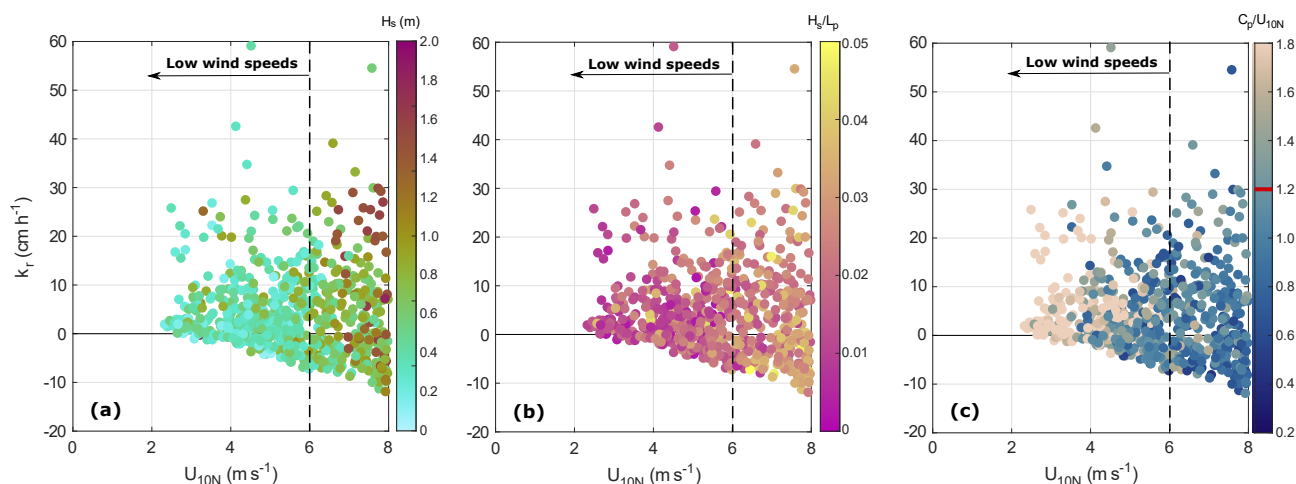


Figure 9. Residual gas transfer velocity ($k_r = k_{660} - k_{W14}$) as a function of the 10 m neutral wind speed under low wind speed conditions. The color represents a) significant wave height, b) wave steepness (H_s/L_p), and c) wave age (C_p/U_{10N}). The vertical dashed lines mark the threshold between low ($U_{10N} < 6 \text{ m s}^{-1}$) and intermediate ($6 < U_{10N} < 8 \text{ m s}^{-1}$) wind speeds. The red mark in c) indicates the theoretical threshold between wind sea ($C_p/U_{10N} < 1.2$) and swell ($C_p/U_{10N} > 1.2$).

3.2.3 Controls on k_{660} in the transition range

Wind speeds between 6 and 8 m s^{-1} were considered here as the transition range between low and high wind speeds. For convenience, the results of this transition range were included in the figures of both high and low wind speed conditions (Sect. 3.2.1 and Sect. 3.2.2, respectively). During the transition range, the atmospheric controls did not seem to be relevant (similar to low wind-speed conditions). Furthermore, the water-side mechanisms showed a mixed behaviour. The wave field showed waves with an average H_s of 0.9 m, in comparison to the 1.4 and 0.6 m of the high and low wind speed regimes, respectively. However, wave steepness and wave age at intermediate winds showed average values of $H_s/L_p=0.03$ and $C_p/U_{10N}=1.1$ similar to the mean values of $H_s/L_p=0.03$ and $C_p/U_{10N}=1.0$ observed at higher wind speeds. Both the wave field and the wind speed seemed to cause stronger mixing as larger values of MLD reaching the entire water column (ca 25 m) were observed under the transition range in contrast to the more persistent stratification at lower wind speeds (Fig. 10b). Values of MLD during these intermediate wind speeds were still lower, on average, than those observed during the high wind speed regime. Meanwhile,

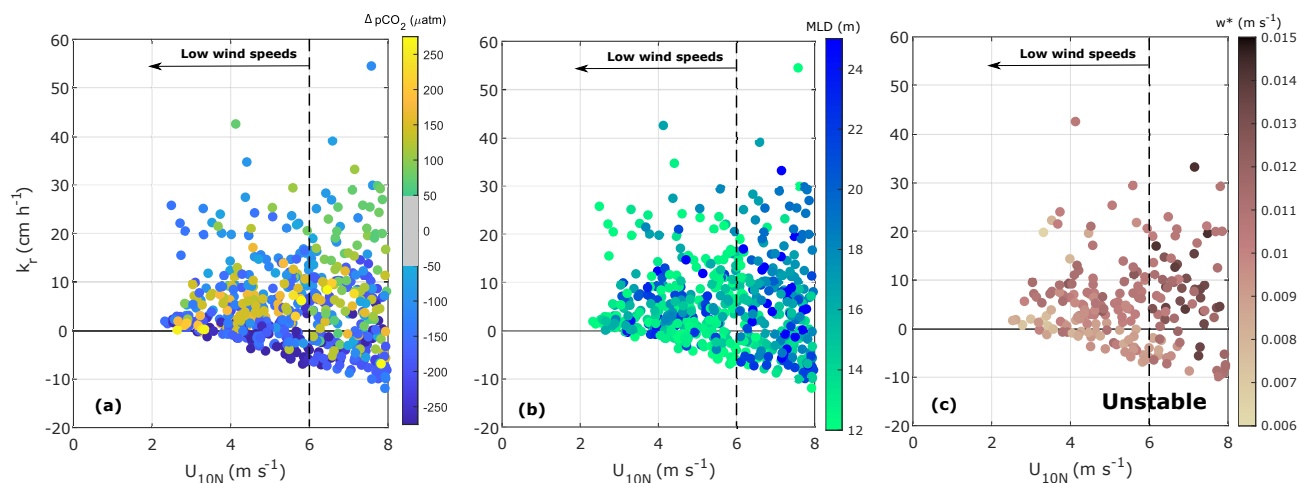


Figure 10. Residual gas transfer velocity ($k_r = k_{660} - k_{W14}$) as a function of the 10 m neutral wind speed under low wind speed conditions. The color represents a) the ΔpCO_2 , b) the mixed layer depth (MLD), and c) the water-side convective scale (w^*) as defined in Rutgersson and Smedman (2010), during unstable atmospheric conditions. The vertical dashed lines mark the threshold between low ($U_{10N} < 6 \text{ m s}^{-1}$) and intermediate ($6 < U_{10N} < 8 \text{ m s}^{-1}$) wind speeds.

large values of w^* suggest the mixing caused by water-side convective processes during unstable atmospheric conditions might be relevant at the intermediate wind speeds (Fig. 10c). In terms of the thermodynamic forcing, positive values of ΔpCO_2 were associated to positive values of k_r similar to the behaviour at lower wind speeds. However, the increased mixing observed at intermediate winds seemed to also enhance fluxes during moderately positive ΔpCO_2 , similar to the observed pattern during high wind speed conditions (Fig. 7a).

4 Discussion

We used nine years of eddy-covariance-based FCO_2 data to evaluate the effect of different control mechanisms on air–sea CO_2 gas exchange. By using this long record, we were able to capture the seasonal variability of the FCO_2 and other parameters relevant to the gas exchange (Sect. 3.1), as well as directly assess controls on k_{660} (Sect. 3.2). The empirically derived k_{660} values showed, in general, a similar wind-speed dependency as the two wind-based parametrizations, the quadratic relationship from Wanninkhof (2014), and the cubic relationship from McGillis et al. (2001). These parametrizations were used as references and we considered that using other commonly-used parametrizations (e.g. Wanninkhof and McGillis, 1999; Nightingale et al., 2000; Ho et al., 2006) would show similar results. However, the scatter of the data suggested that a large proportion of the variability in k_{660} was not explained by variations in the wind speed (Fig. 5). Mechanisms other than the wind have been recognized to be relevant controls on air–sea gas exchange processes, particularly in heterogeneous environments such as marginal and coastal seas (e.g. Upstill-Goddard, 2006; Gutiérrez-Loza et al., 2021). Based on the results presented here, we showed that wind-based parameterizations developed for larger basins and open-ocean conditions can be adequate approxima-



tions for long-term gas transfer velocity calculations, even in heterogeneous environments such as the studied area. However, we noticed conditions where the existing models failed to replicate the observations, and we further showed that in order to capture the short-term variability in k_{660} additional process need to be accounted for.

The characteristics of the wave field are one of the most important parameters involved on the air–sea gas transfer directly affecting the conditions of the interface. However, the effect of waves on gas exchange has been very seldom studied in coastal regions (e.g. Gutiérrez-Loza et al., 2018). We analyzed different parameters describing the wave field (significant wave height, wave steepness, and wave age) and their effect on the gas transfer velocity. At high wind speeds, the waves were found to be locally generated by the wind (i.e. wind sea), which made it hard to separate the effect of the waves from that of the wind. The wave climate in the Baltic Sea is characterized by predominant short wind sea and scarce long swell (Soomere, 2022). Furthermore, the effect of swell waves was not relevant at high wind speed conditions, and while at lower wind speeds some effect from swell could be expected, such effect was partly excluded here as wind speeds lower than 2 m s^{-1} were removed from the analysis. Values of wave steepness suggested that there was no systematic wave breaking. However, empirical knowledge of the study site suggests the presence of breaking, particularly of the smallest/youngest waves. These waves are not expected to generate significant bubble entrainment into the upper layer of the sea, but generation of sea spray is expected at the highest wind speeds.

The effect of sea spray in the air–sea gas exchange has been reported before (Andreas et al., 2016, 2017). However, the impact of a broken surface caused by wave breaking and whitecapping, has been shown to cause an asymmetric transport due to the injection of bubbles into the ocean enhancing the negative (downward) fluxes (Woolf, 1997; Leighton et al., 2018). This was not seen here, instead, we found that the increase in the k_{660} was associated to the ocean-to-atmosphere (upward) transport in what we suggest is the effect of sea spray. This effect was clearly affected with the seasonal variability as the increased k_{660} were observed during the winter months, when dry and cold air (colder than the ocean) are common (Fig.4c and e). During the winter months positive CO_2 gradients prevail and positive gas and heat fluxes are most often observed. Dedicated analysis of the effect of bubbles and sea spray in coastal regions are still necessary. Particularly at high latitudes, where strong seasonality may cause significant differences in the forcing mechanisms modulating the influx and efflux of gas into/from the ocean (see Sect. 3.1). A brief analysis of the effect of sea spray on air–sea fluxes made using a bulk algorithm for turbulent air–sea fluxes, freely available from www.nwra.com/resumes/andreas/software.php (Andreas et al., 2008). The results showed that the large enthalpy fluxes observed under unstable conditions at high wind speeds (Fig. 8b) can be attributed to sea spray. Figure 11 shows sea spray mediated fluxes of up to 30 W m^{-2} (up to 20 % of the total enthalpy flux) under these conditions. We suggest sea spray can enhance the transport of CO_2 across the interface, as it does with heat and moisture. Based on our results, this mechanism can be a relevant control of the gas exchange by enhancing upward CO_2 fluxes at wind speeds higher than $8\text{--}10 \text{ m s}^{-1}$ under unstable atmospheric conditions and low relative humidity. Further data under these particular conditions is necessary to establish a better understanding of the effect of sea spray mediated gas fluxes.

Under calm wind speed conditions, with $U_{10N} < 6 \text{ m s}^{-1}$, the effect of water-side convection was found to be relevant to explain the larger values of k_{660} during unstable atmospheric conditions at these wind speeds. Previous studies have shown the effect of convective processes on the gas transfer in the Baltic Sea (Rutgersson and Smedman, 2010; Norman et al., 2013b).

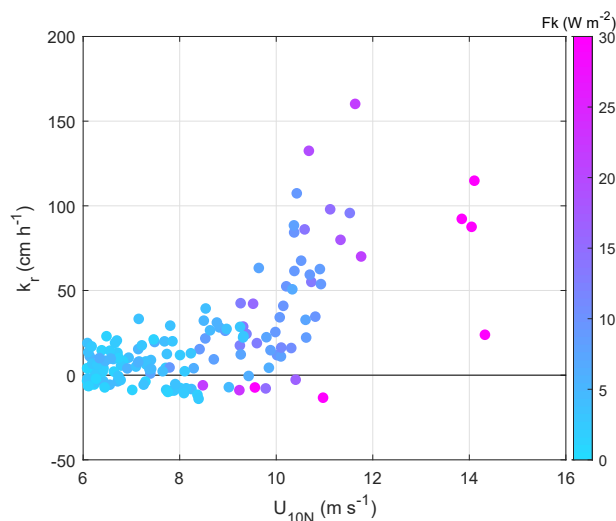


Figure 11. Residual gas transfer velocity ($k_r = k_{660} - k_{W14}$) as a function of the 10 m neutral wind speed under low wind speed conditions. The color represents the sea-spray-mediated enthalpy flux calculated using a bulk algorithm for air–sea turbulent fluxes (Andreas et al., 2008).

The same studies showed that the water-side convective processes are not relevant under stable and neutral conditions, as well as at wind speeds higher than 8 m s^{-1} .

The ability to correctly calculate the gas transfer velocity from eddy covariance FCO_2 and ΔpCO_2 measurements is a matter of discussion. For the data presented here, this limitation was evident at low wind speeds during the summer months, when a strongly stratified and highly heterogeneous surface layer might not be well represented by the pCO_2^w measurements at 4 m depth. In order to avoid this uncertainty in our k_{660} calculations, we removed data based on a ΔT_w criterion (see Sect. 2.2). However, we still observed scattered data at the lowest wind speeds corresponding to moderate ΔpCO_2 negative values and shallow MLD (Fig. 10a and b). Due to the uncertainty, we did not make further conclusions based on these data, and we encourage the reader to be cautious when interpreting these results.

The occurrence of upwelling events can be up to 25-30% in some regions of the Baltic Sea (Lehmann et al., 2012). Norman et al. (2013a) found that effect of these events on the annual mean FCO_2 can reach up to 25% for the Baltic Sea conditions. In this study we did not analyze the relative importance of upwelling on FCO_2 or k_{660} , however, we recognize that part of the observed inter-annual variability might be caused by upwelling events. An in-depth analysis of the data is necessary to evaluate individual upwelling events captured during the nine years of measurements. Furthermore, using the parameterization suggested by Pereira et al. (2018) based on the skin temperature, we calculated a rough estimate of the effect of surfactants on our long-term k_{660} . The results suggest an overall effect of merely -0.1 %, suggesting a very small reduction of k_{660} when surfactants are taken into account. A detail analysis of the effect of surface films is beyond the scope of this study. Nevertheless, we recognize that this process might be particularly relevant in coastal seas and other shallow bodies of water. The analysis of the effect of sea ice, precipitation, Langmuir circulation, and other processes in the Baltic Sea is still pending.



360 We presented what we believe is the longest record of directly measured FCO_2 and the corresponding gas transfer velocities
from in situ data. Regardless of the length of our data set, we still found three major limitations when addressing the effect of
forcing mechanisms on air–sea gas exchange: 1) the use of a land-based station provides significant advantages, but it hinders
the evaluation of the spatial variability of the exchange; 2) the intrinsic characteristics of the eddy covariance technique, in
combination to the exclusion of wind directions not representative of open-sea conditions, resulted in a patchy data set with
365 significant gaps throughout the time series; and 3) limited amount of data at high wind speeds ($U_{10N} > 10 \text{ m s}^{-1}$) was recorded,
which hindered the analysis and increases the uncertainties at those critical wind conditions. The results presented here provide
significant insights on the air–sea gas exchange and its variability. Nevertheless, the limitations are still large and a continued
effort from the scientific community is required.

Global coastal oceans and marginal seas are one of the most vulnerable environments subjected to the effect of climate
370 change and anthropogenic pressures (Pachauri et al., 2014). Understanding the role of these regions in the global carbon cycle
has become an essential aspect in order to address the challenges of the current and future climates. In this sense, the Baltic
Sea can be seen as a test basin which provides a wide variety of physical and biogeochemical conditions. At the same time,
the carbon system of the Baltic Sea has been relatively well-documented (e.g. Kuliński and Pempkowiak, 2011; Schneider and
Müller, 2018), and the region has been a relevant study area in terms of mitigation and environmental management (Reusch
375 et al., 2018). The results presented here are most probably relevant for other marginal seas and coastal areas.

5 Conclusions

We presented a large data set of directly measured air–sea CO_2 fluxes by eddy covariance from a land-based station in the
Baltic Sea. The forcing mechanisms acting on the surface of the ocean and their relative effect on the gas exchange can widely
vary depending on the wind-speed. Therefore, the air–sea gas exchange, controlled by such forcing mechanisms, can also be
380 expected to be affected in different ways depending on the wind-speed regime. We investigated the effect of the water-side and
atmospheric conditions on the gas transfer velocity under relatively high, intermediate, and low wind speed regimes.

At high wind speeds larger values of the residual gas transfer velocity were observed. These values occurred during events of
cold and dry air, under unstable atmospheric conditions. We suggest, based on these results, that sea spray might be one of the
most effective mechanisms enhancing air–sea CO_2 fluxes under these conditions, most probably due to large evaporation rates.
385 The effect of the wave field was particularly evident in terms of the significant wave height, with high gas transfer velocity
values occurring when large values of H_s were observed. However, the effect of the wave field was not completely decoupled
from the effect of the wind given that most of the waves were locally generated. Thus, theoretically removing the wind effect
might not be sufficient to individually evaluate the effect of waves on the gas exchange. Intermediate wind speeds showed a
mixed behaviour, thus, we defined these wind speeds as a transition range. Under these conditions, the mixing induced by the
390 wind and waves to the upper layer of the sea was found to be the most relevant parameter for the gas exchange. Particularly
when moderate ΔpCO_2 values occurred, similarly to high wind speed conditions. Under low wind speed conditions, water-side



convection was the only parameter explaining part of the variability on k_r , particularly during the winter. During the summer, strong stratification occurred, hindering the downward fluxes.

395 Currently existing wind-based parametrizations showed to adequately represent k_{660} in the long-term averages. However, further investigation of parameters affecting the seasonal and inter-annual variability of the fluxes is needed to improve our understanding of air–sea gas exchange and adequately represent k_{660} in shorter time scales. Similarly, a detail analysis of bubble- and sea-spray-mediated fluxes is needed to contribute to the understanding of CO₂ fluxes in the coastal regions.



Data availability. The raw data supporting the conclusions of this manuscript is available upon request.

Author contributions. LG-L, MW, ES and AR were involved in the conceptualization of the project. AR was in charge of the project administration and funding acquisition. The formal analysis was carried out by LG-L with the support of EN. LG-L, MW, EN and AR participated in the continuous maintenance of the Östergarnsholm station and data acquisition. LG-L wrote the original draft with contribution from all the co-authors.

Competing interests. The authors declare that they have no conflict of interest.

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